



Diffusion of Oxygen in Ni and Fe due to Substitutional Vacancy from DFT Calculations

Zi-Kui Liu, Huazhi Fang and ShunLi Shang

Department of Materials Science & Engineering

Pennsylvania State University, University Park, PA 16802

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First-Principles Calculation of Self-Diffusion Coefficients

M. Mantina, Y. Wang, R. Arroyave,^{*} L. Q. Chen, and Z. K. Liu

Department of Materials Science and Engineering, The Pennsylvania State University, University Park, Pennsylvania 16802, USA

C. Wolverton

Department of Materials Science and Engineering, Northwestern University, Evanston, Illinois 60208, USA

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We demonstrate a first-principles method to compute all factors entering the vacancy-mediated self-diffusion coefficient. Using density functional theory calculations of fcc Al as an illustrative case, we determine the energetic and entropic contributions to vacancy formation and atomic migration. These results yield a quantitative description of the migration energy and vibrational prefactor via transition state theory. The calculated diffusion parameters and coefficients show remarkably good agreement with experiments. We provide a simple physical picture for the positive entropic contributions.

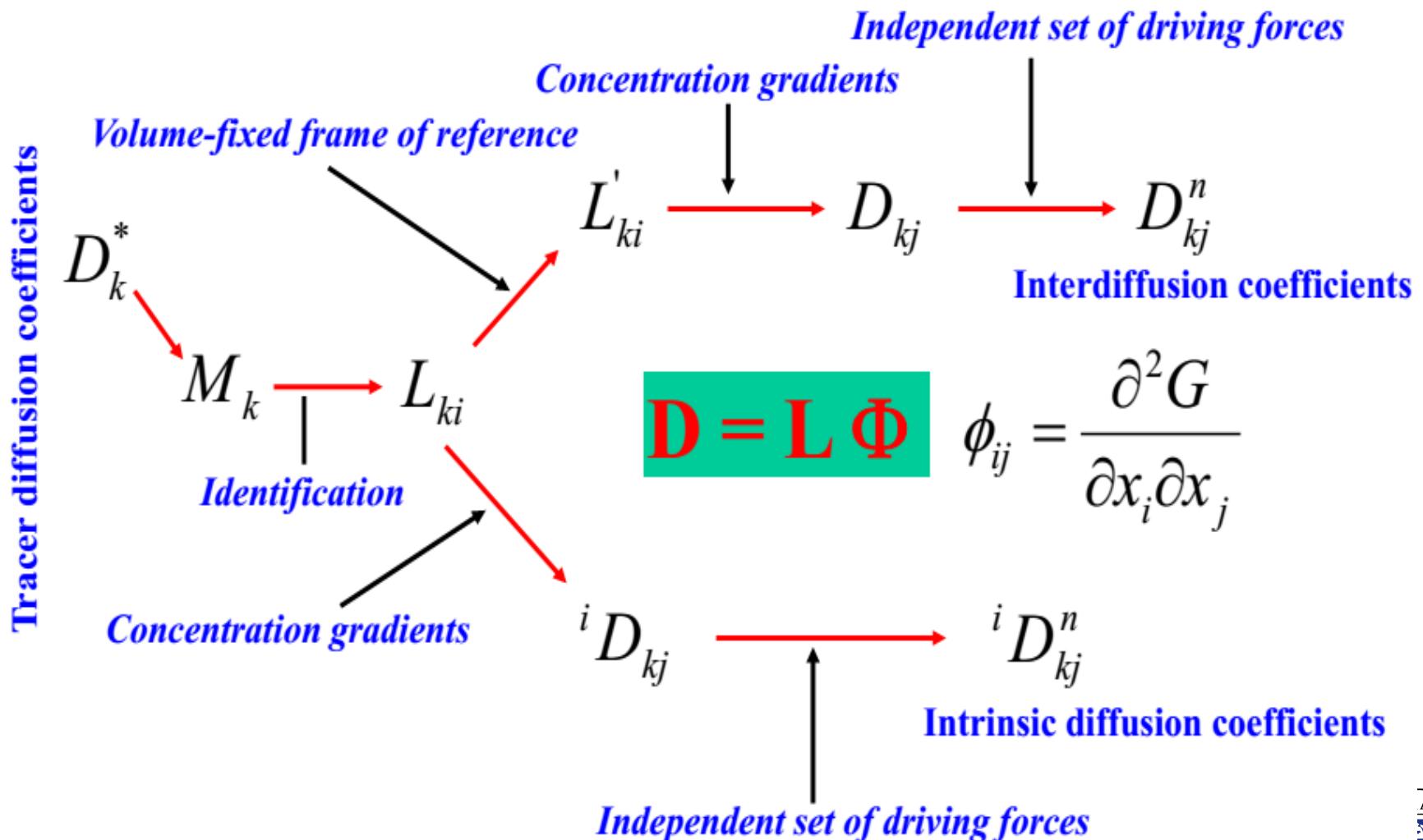
- Use density functional theory calculations to determine the energetic and entropic contributions to vacancy formation and atomic migration, i.e. all parameters used to evaluate vacancy-mediated diffusion coefficients
- Compare the calculated values with available experimental data



Outline

- **Interstitial diffusion from first-principles**
 - Theory and finite temperature thermodynamics
- **Oxygen diffusion in fcc nickel**
 - Site preference & diffusion pathway
 - Va-modified site preference & diffusion path
 - Oxygen diffusivity with and without Va
- **Oxygen diffusion in bcc iron**
 - Site preference with and without Va
 - Oxygen diffusivity with and without Va
- **Summary**

Diffusion coefficient **D** is a product of kinetic coefficient **L** and thermodynamic factor **Φ**



Interstitial diffusion equation

Following Wert & Zener, interstitial diffusivity can be described by:

$$D = n\beta d^2 \Gamma$$

Coordinate number Jump distance
 Probability factor Jump frequency

$$\Gamma = \frac{kT}{h} \frac{Z_{TS}}{Z_{IS}} e^{-\Delta E/kT} = \frac{kT}{h} \frac{\prod_{i=1}^{3N-3} \left[2 \sinh \left(\frac{h\nu_i^{IS}}{2kT} \right) \right]}{\prod_{i=1}^{3N-4} \left[2 \sinh \left(\frac{h\nu_i^{TS}}{2kT} \right) \right]} e^{-\Delta E/kT}$$

High temperatures ($h\nu/2kT \ll 1$) $\Gamma = \frac{\prod_{i=1}^{3N-3} \nu_i^{IS}}{\prod_{i=1}^{3N-4} \nu_i^{TS}} e^{-\Delta E/kT} = \nu^* e^{-\Delta E/kT}$

\Rightarrow **Vineyard** equation

Low temperatures ($h\nu/2kT \gg 1$) $\Gamma = \frac{kT}{h} \frac{\exp \sum_{i=1}^{3N-3} \frac{h\nu_i^{IS}}{2kT}}{\exp \sum_{i=1}^{3N-4} \frac{h\nu_i^{IS}}{2kT}} e^{-\Delta E/kT} = \frac{kT}{h} e^{-(\Delta E + \Delta G_{vib})/kT}$

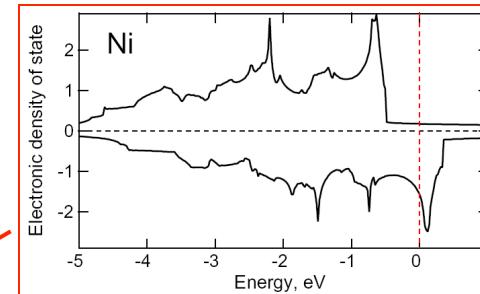
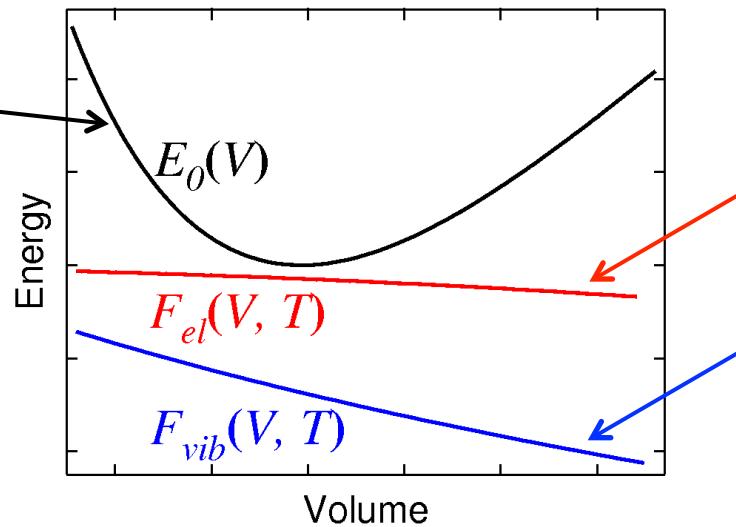
\Rightarrow **Eyring** equation

Finite temperature thermodynamics - The quasiharmonic approach

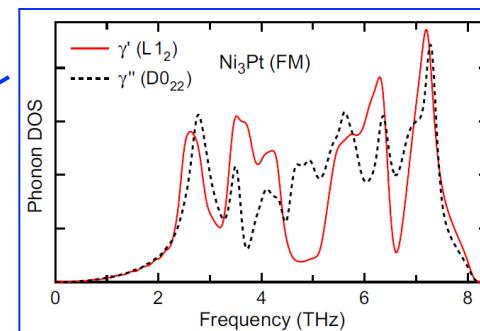
$$F(V, T) = E_0(V) + F_{vib}(V, T) + F_{el}(V, T)$$

- $E_0(V)$ Static energy at 0 K and volume V , i.e., EOS (by VASP)
- $F_{vib}(V, T)$ Vibrational contribution at V & T (Phonon or Debye model)
- $F_{el}(V, T)$ Thermal electronic contribution at V & T (by VASP)

DFT directly



E-DOS



P-DOS

Calculation methods – O in fcc Ni & O in bcc Fe

- DFT based first-principles calculations
 - VASP code
 - PAW method for electron-ion interaction
 - Exchange-correctional (X-C) functionals
 - LDA for O in fcc Ni
 - GGA-PBE for O in Fe
 - Also tested others: PW91, PBE+U, PBEsol, AM05
 - CI-NEB for migration barrier calculations
 - Phonon & Debye model for finite T thermodynamics



Case 1 of 2

- Oxygen diffusion in fcc Ni
with and without vacancy



Oxygen diffusion in fcc Ni

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First-principles studies on vacancy-modified interstitial diffusion mechanism of oxygen in nickel, associated with large-scale atomic simulation techniques

H. Z. Fang,^{1,2} S. L. Shang,^{1,2} Y. Wang,^{1,2} Z. K. Liu,^{1,2} D. Alfonso,^{1,3} D. E. Alman,^{1,3} Y. K. Shin,^{1,4} C. Y. Zou,^{1,4} A. C. T. van Duin,^{1,4} Y. K. Lei,^{1,5} and G. F. Wang^{1,5}

¹National Energy Technology Laboratory Regional University Alliance, U.S. Department of Energy, Pittsburgh, Pennsylvania 15236, USA

²Department of Materials Science and Engineering, The Pennsylvania State University, University Park, Pennsylvania 16802, USA

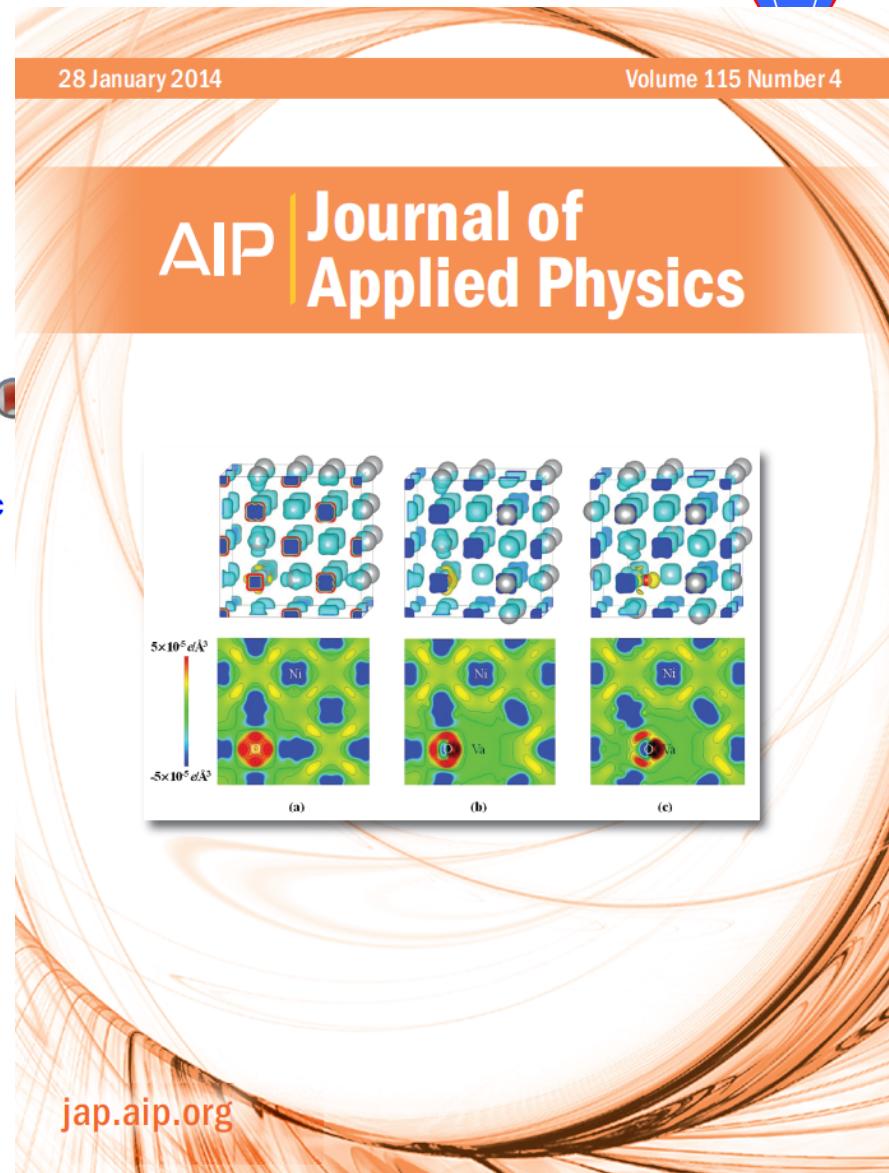
³National Energy Technology Laboratory, U.S. Department of Energy, Pittsburgh, Pennsylvania 15236, USA

⁴Department of Mechanical and Nuclear Engineering, The Pennsylvania State University, University Park, Pennsylvania 16802, USA

⁵Department of Mechanical Engineering and Materials Science, University of Pittsburgh, Pennsylvania 15261, USA

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This paper is concerned with the prediction of oxygen diffusivities in *fcc* nickel from first-principles calculations and large-scale atomic simulations. Considering only the interstitial octahedral to tetrahedral minimum energy pathway for oxygen diffusion in *fcc* lattice, greatly underestimates the migration barrier and overestimates the diffusivities by several orders of magnitude. The results indicate that vacancies in the Ni-lattice significantly impact the migration barrier of oxygen in nickel. Incorporation of the effect of vacancies results in predicted diffusivities consistent with available experimental data. First-principles calculations show that at high temperatures the vacancy concentration is comparable to the oxygen solubility, and there is a strong binding energy and a redistribution of charge density between the oxygen atom and vacancy. Consequently, there is a strong attraction between the oxygen and vacancy in the Ni lattice, which impacts diffusion. © 2014 AIP Publishing LLC. [<http://dx.doi.org/10.1063/1.4861380>]



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Oxygen diffusion in fcc Ni – site preference

Solution energy:

For interstitial solution:

$$E_O^{\text{sol}} = E_{n\text{Ni}+\text{O}} - E_{n\text{Ni}} - E_{\text{O}},$$

For substitutional solution: $E_O^{\text{sol}} = E_{(n-1)\text{Ni}+\text{O}} - \frac{n-1}{n} E_{n\text{Ni}} - E_{\text{O}}$

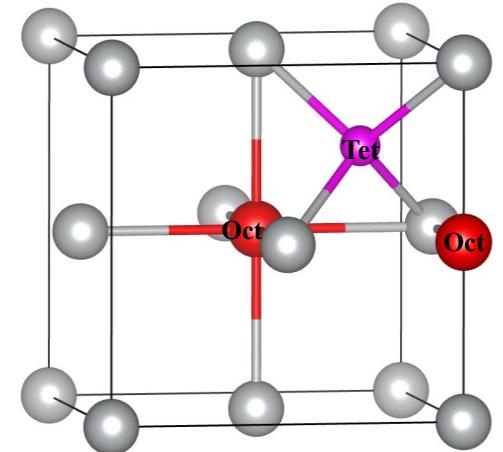
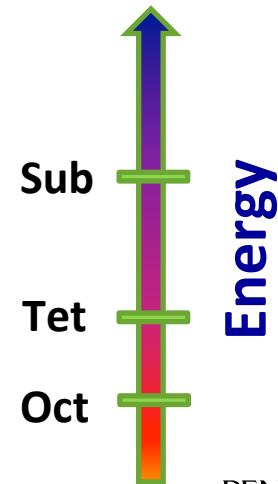


TABLE I. Solution energies (eV) of oxygen in pure nickel.

Substitutional	Interstitial		Method
	Octahedral	Tetrahedral	
E_O^{sol}	-2.31 —	-3.63 -2.65, ^a -3.39 ^b	-3.28 -2.38 ^a
			This work Other DFT

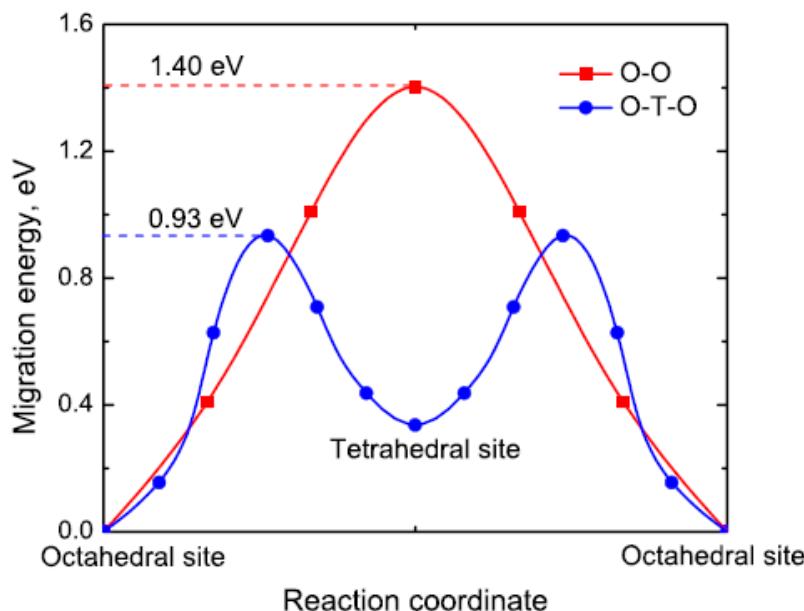
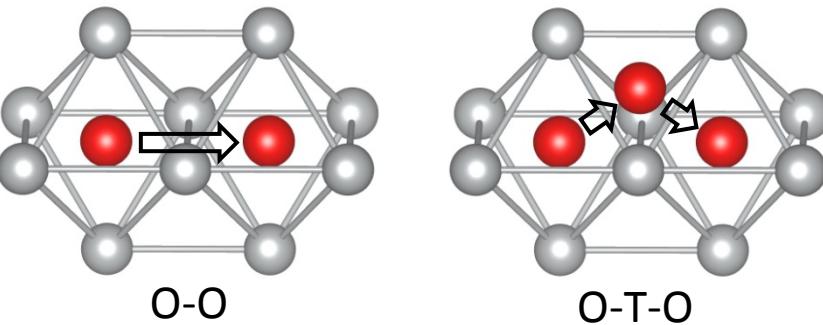
^a E. H. Megchiche, et al, JPCM 19 (2007) 296201

^b J. J. Kim, et al, APL 100 (2012) 131904



Oxygen diffusion in fcc Ni – diffusion pathway

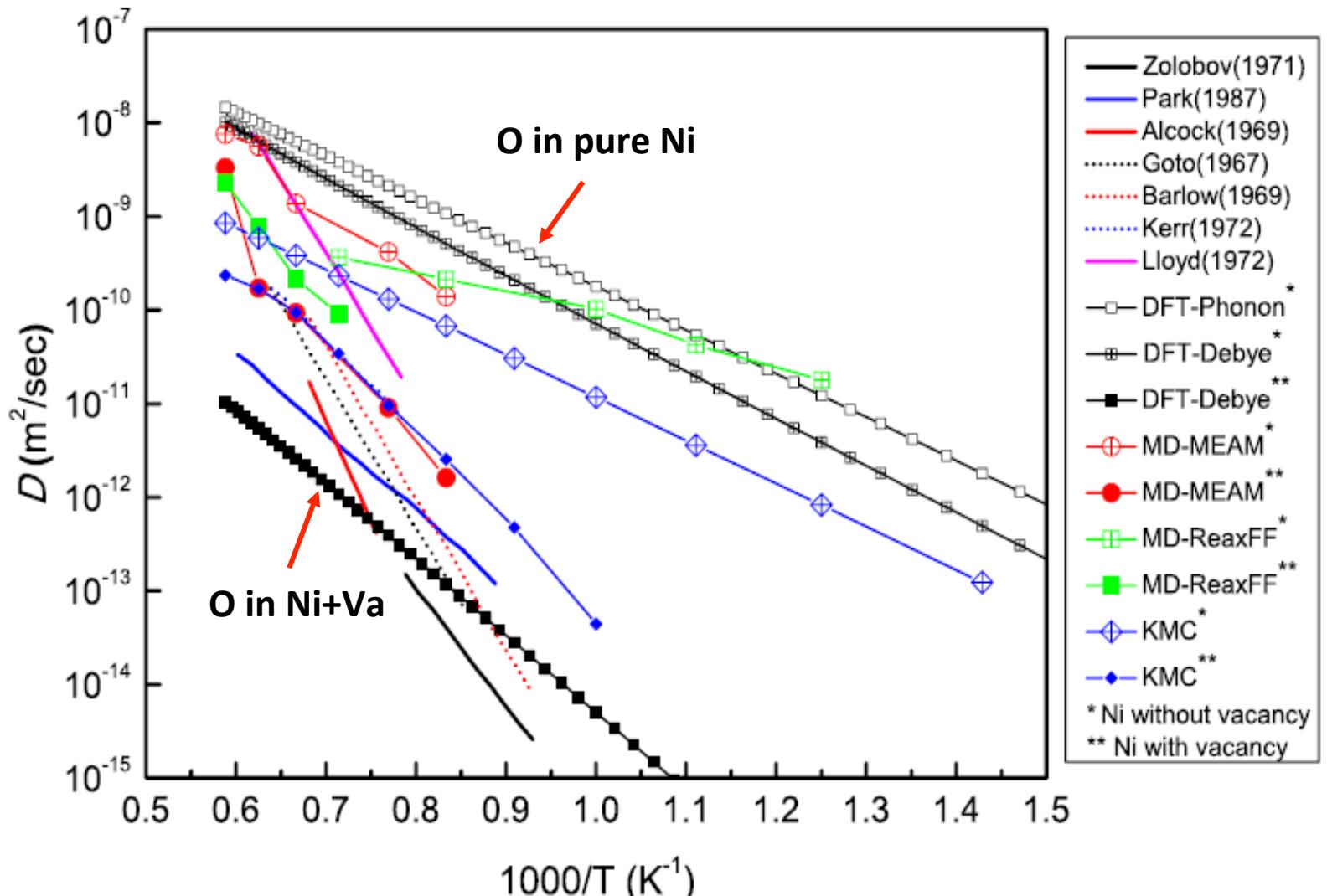
- Two migration pathways are considered



	O-O	O-T-O
This work	1.40	0.93
Other DFT	1.57 1.44	1.23 1.12
Experiment	1.70, 3.20, 3.12, 3.08, 4.28, 2.49, 1.89	

Predicted barriers much lower than measurements

Oxygen diffusion in fcc Ni – finite T diffusivity

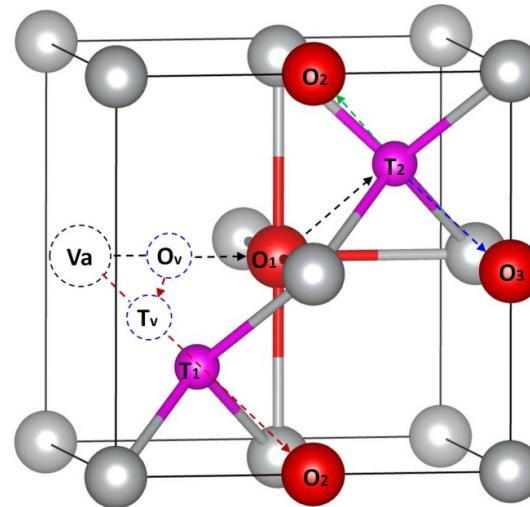
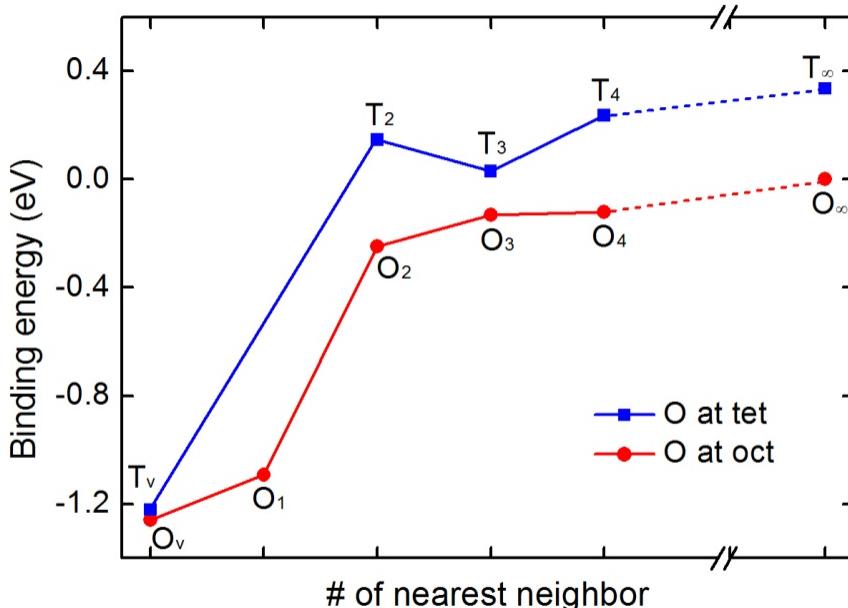




Oxygen diffusion in fcc Ni – vacancy modified site preference

- With Va, the neighboring Oct & Tet sites shift, *i.e.*:

$$\text{O}_1 \rightarrow \text{O}_v; \quad \text{T}_1 \rightarrow \text{T}_v$$

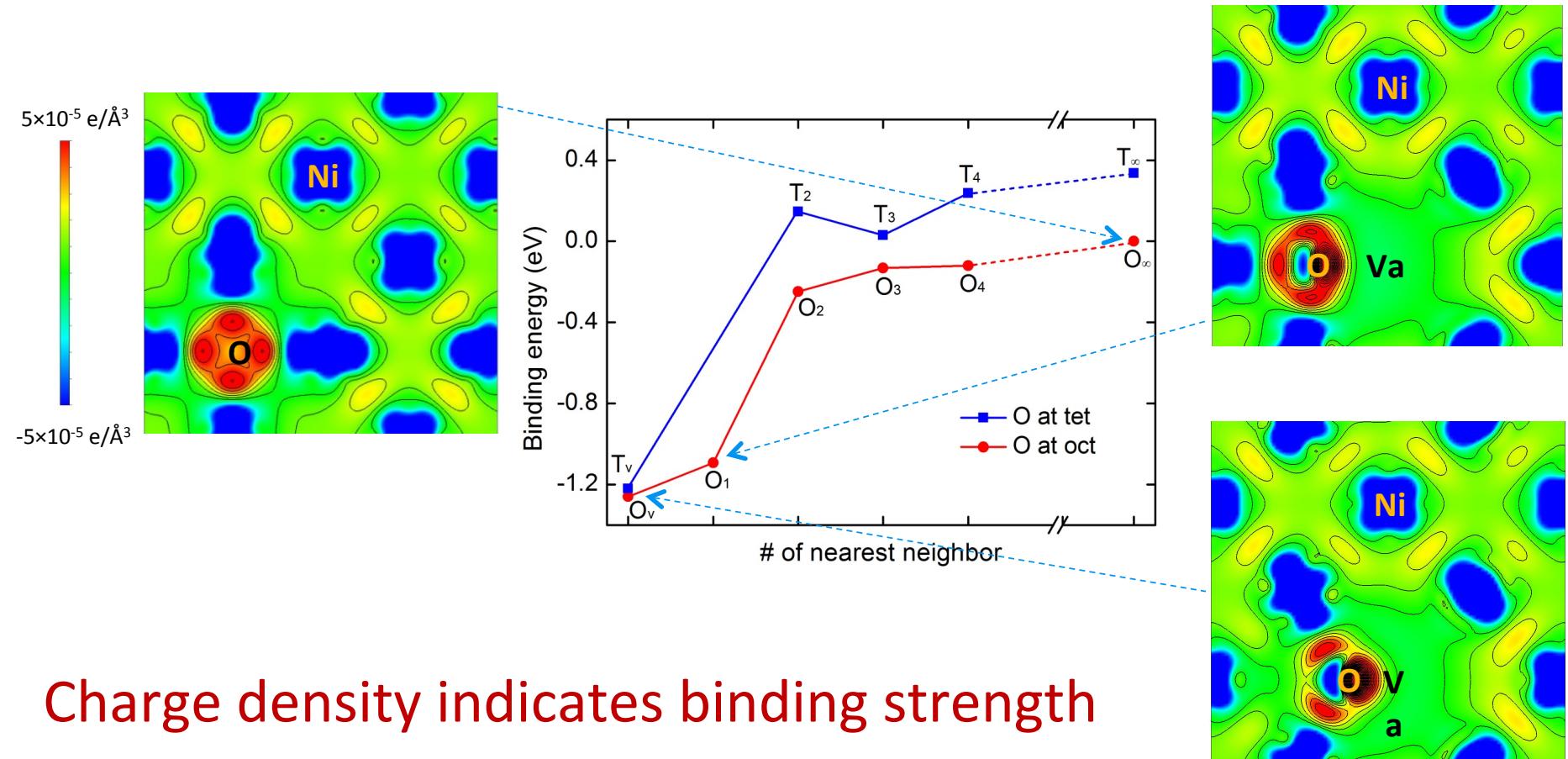


- O-Va binding energy ~ 1.2 eV.
- Effect up to the 3rd NN (O₃).





Oxygen diffusion in fcc Ni – binding energy of O and Va

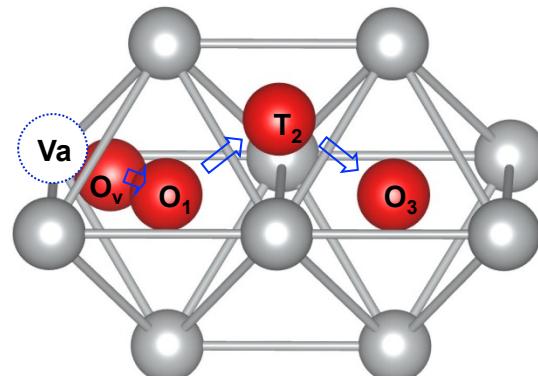


Charge density indicates binding strength

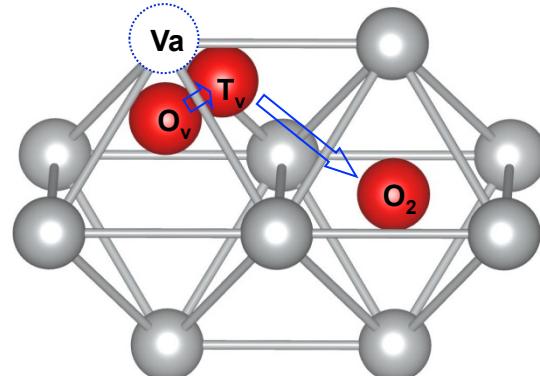


Oxygen diffusion in fcc Ni – Va modified diffusion pathway

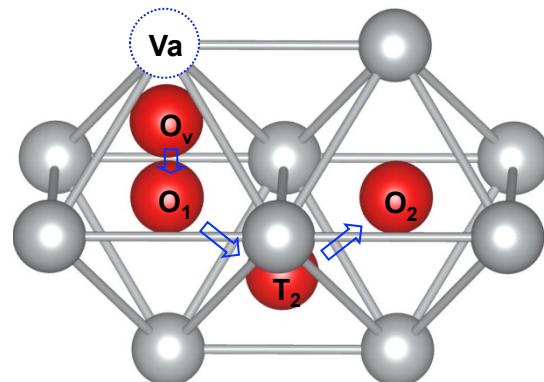
- Three possible paths for O diffusion from Va



(a) O_v - O_1 - T_2 - O_3



(b) O_v - T_v - O_2



(c) O_v - O_1 - T_2 - O_2

Oxygen diffusion in fcc Ni – Va modified diffusion barrier

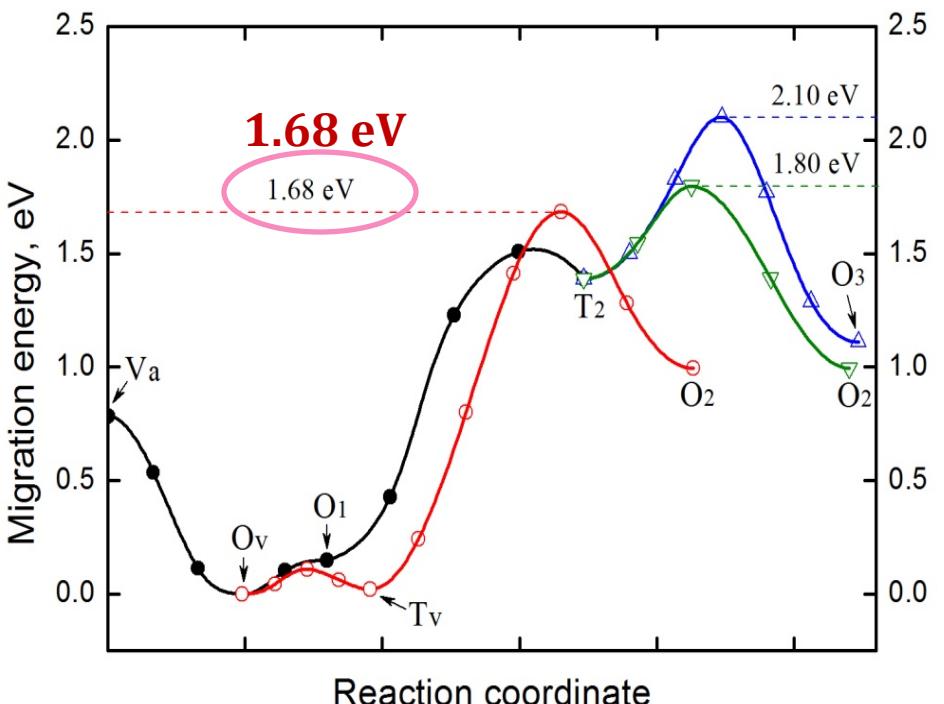


TABLE II. Migration energies (eV) of oxygen in nickel without and with vacancy effect.

	<i>O-O</i>	<i>O-T-O</i>	Method
E_O^{mig}	1.40 ^a 1.57, ^c 1.44 ^d	0.93, ^a 1.68, ^b 1.80, ^b 2.10 ^b 1.23, ^c 1.12 ^d 1.99 ^e 1.70, ^f 1.89, ^g 3.08, ^h 2.49, ⁱ 3.2, ^j 4.28, ^k 3.12 ^l	This work Other DFT Other MD Experiment

^aNi without vacancy.

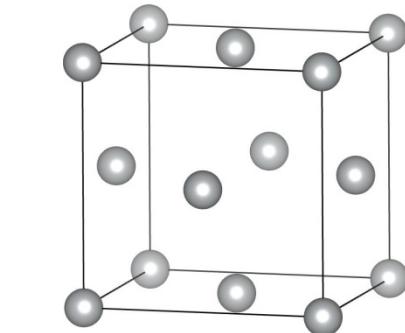
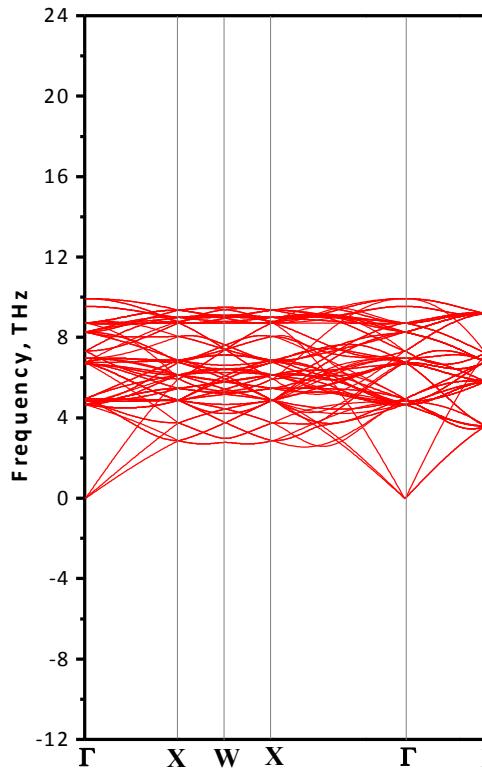
^bNi with vacancy.

Significant increase of diffusion barrier from 0.93 to 1.68 eV

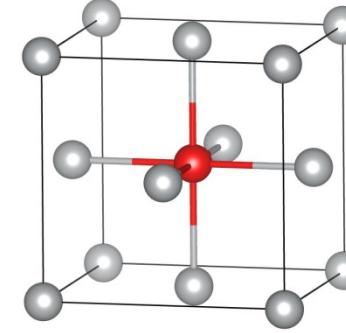
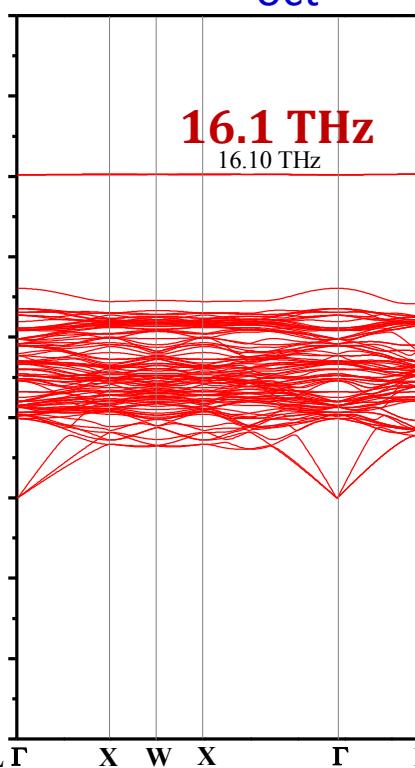


Oxygen diffusion in fcc Ni - phonon

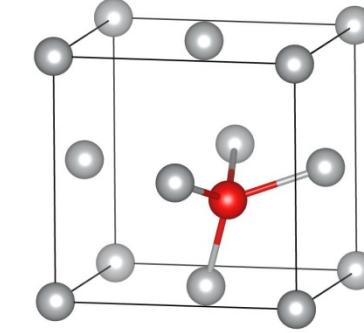
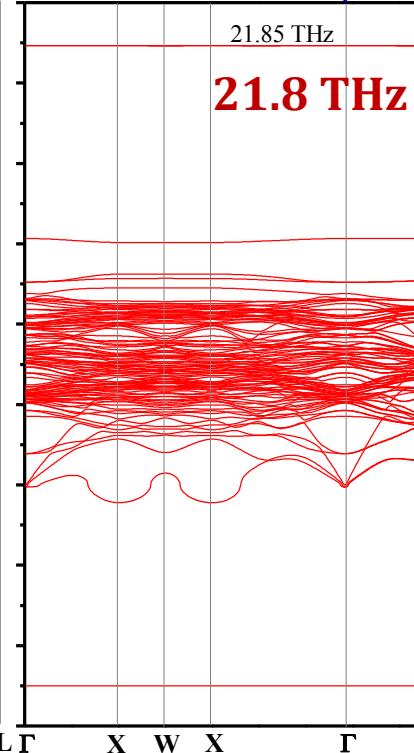
Pure Ni



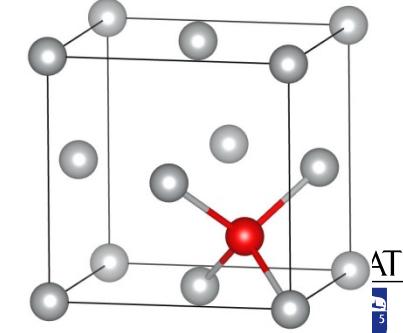
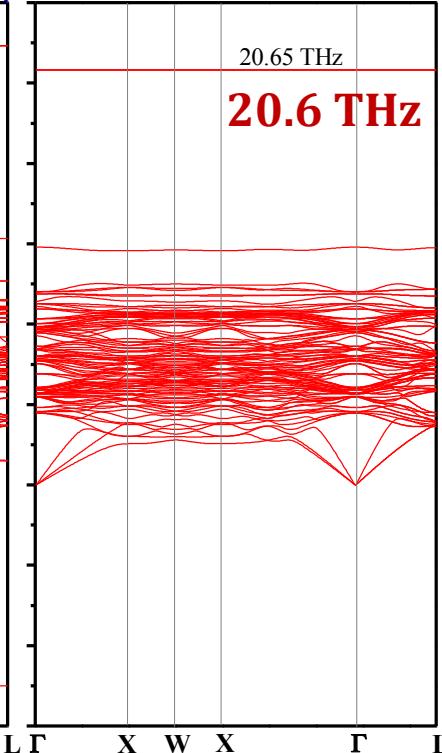
Ni-O_{oct}



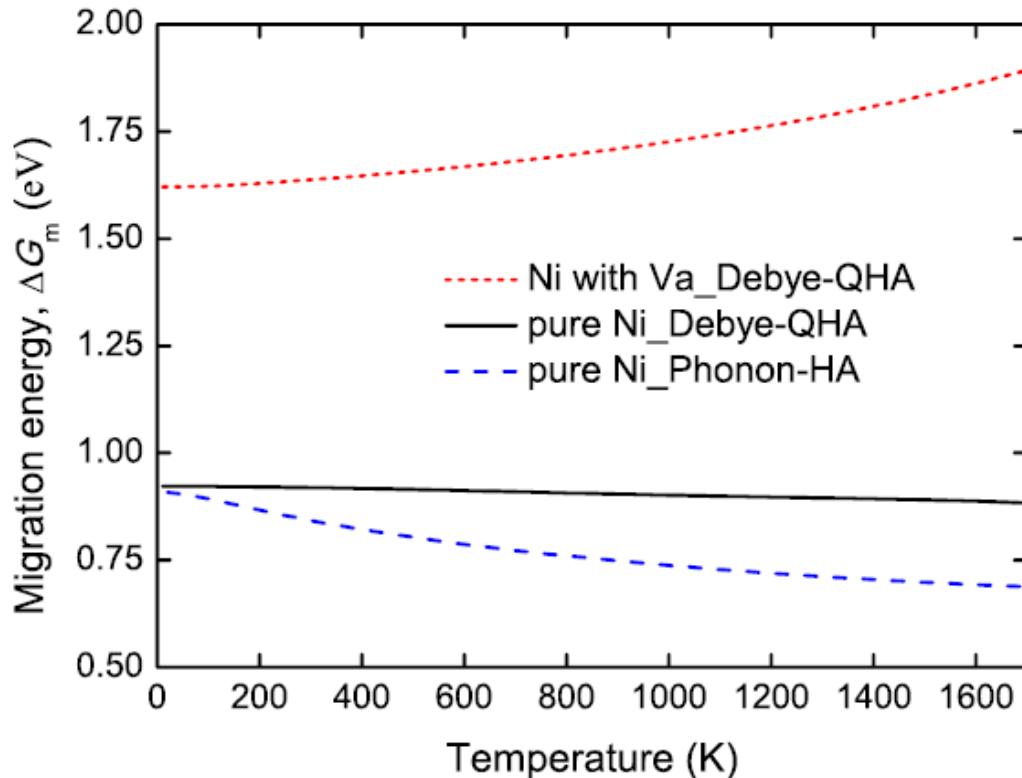
Ni-O_{saddle (oto)}



Ni-O_{tet}



Oxygen diffusion in fcc Ni – migration barrier at finite T



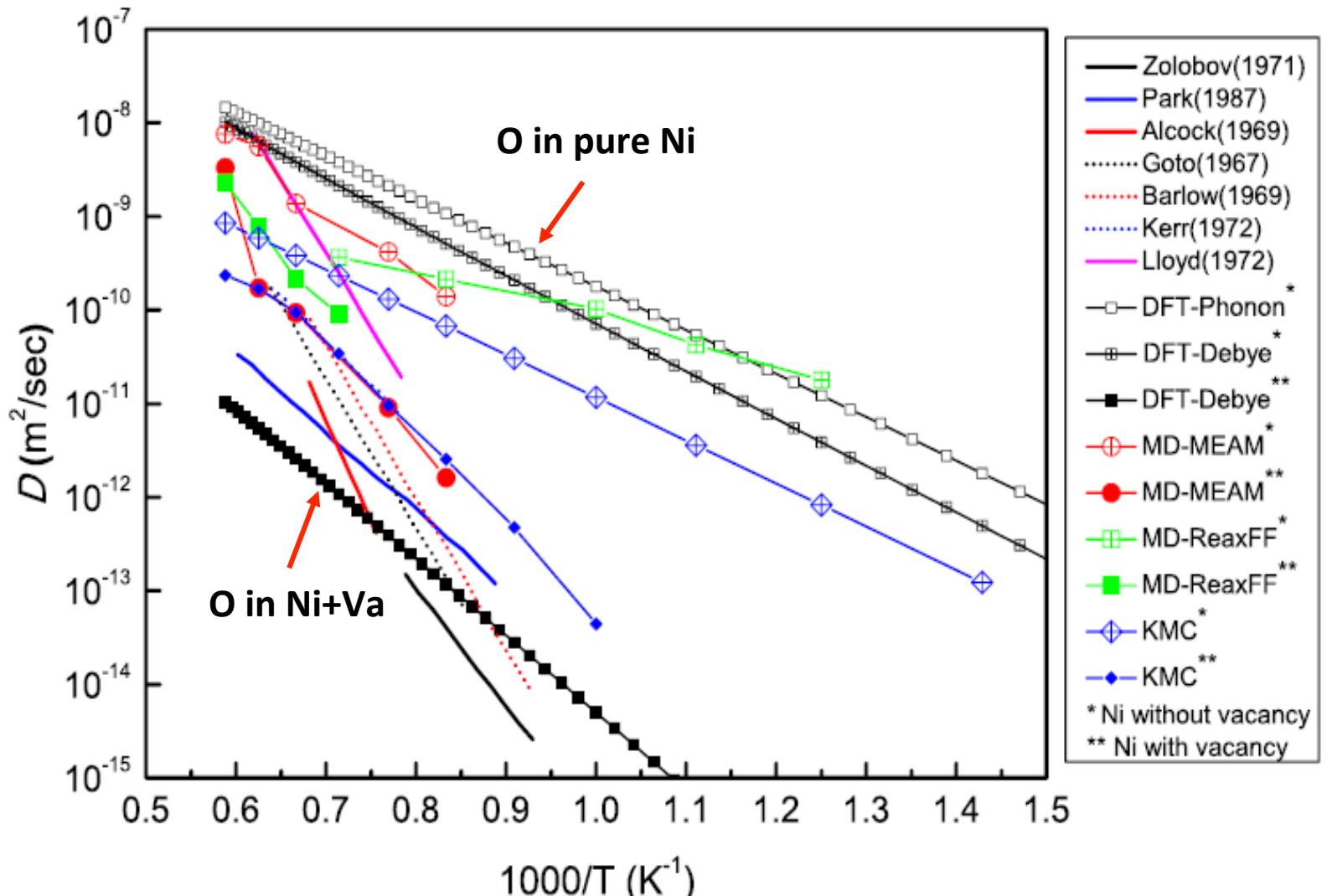
Jump frequency:

$$\Gamma = \frac{kT}{h} e^{-\Delta G_m / kT}$$

- **Phonon:** accurate but expensive
- **Debye:** efficient but less expensive

Migration free energies for $O_v-T_v-O_2$ pathway

Oxygen diffusion in fcc Ni – finite T diffusivity





Oxygen diffusion in fcc Ni -

Arrhenius parameters

$$D = D_0 \exp(-Q/kT)$$

TABLE III. Diffusion constants D_0 and activation energies Q predicted by different computational methods in comparison with available experimental data.

D_0 (m ² /s)	Q (eV)	T (K)	Method	Ref.
8.0×10^{-6}	0.92	700–1700	Without Va	DFT-Phonon ^a This work
7.4×10^{-6}	0.98	700–1700		DFT-Debye ^a This work
4.3×10^{-7}	0.91	700–1700		kMC ^a This work
6.4×10^{-6}	0.95	1200–1700		MD-MEAM ^a This work
2.3×10^{-8}	0.49	800–1400		ReaxFF ^a This work
8.8×10^{-7}	1.64	700–1700	With Va	DFT-Debye ^b This work
3.9×10^{-4}	1.82	1000–1700		kMC ^b This work
2.8×10^{-4}	1.95	1200–1700		MD-MEAM ^b This work
9.5×10^{-3}	2.25	1000–1400		ReaxFF ^b This work
4.9×10^{-6}	1.70	1123–1673	Expt	Potentiometric Park and Altstetter, 1987
7.9	3.20	1073–1473		Internal oxidation Barlow and Grundy, 1969
1.82	3.12	1173–1573		Internal oxidation Goto <i>et al.</i> , 1967
26.8	3.08	1273–1623		Internal oxidation Lloyd and Martin, 1972
8.93×10^3	4.28	1323–1473		Gravimetric Alcock and Brown, 1969
1.21×10^{-3}	2.49	623–1273		Desorption Zholobov and Malev, 1971
2.06×10^{-4}	1.89	1273–1573		Electrochemistry Kerr, 1972

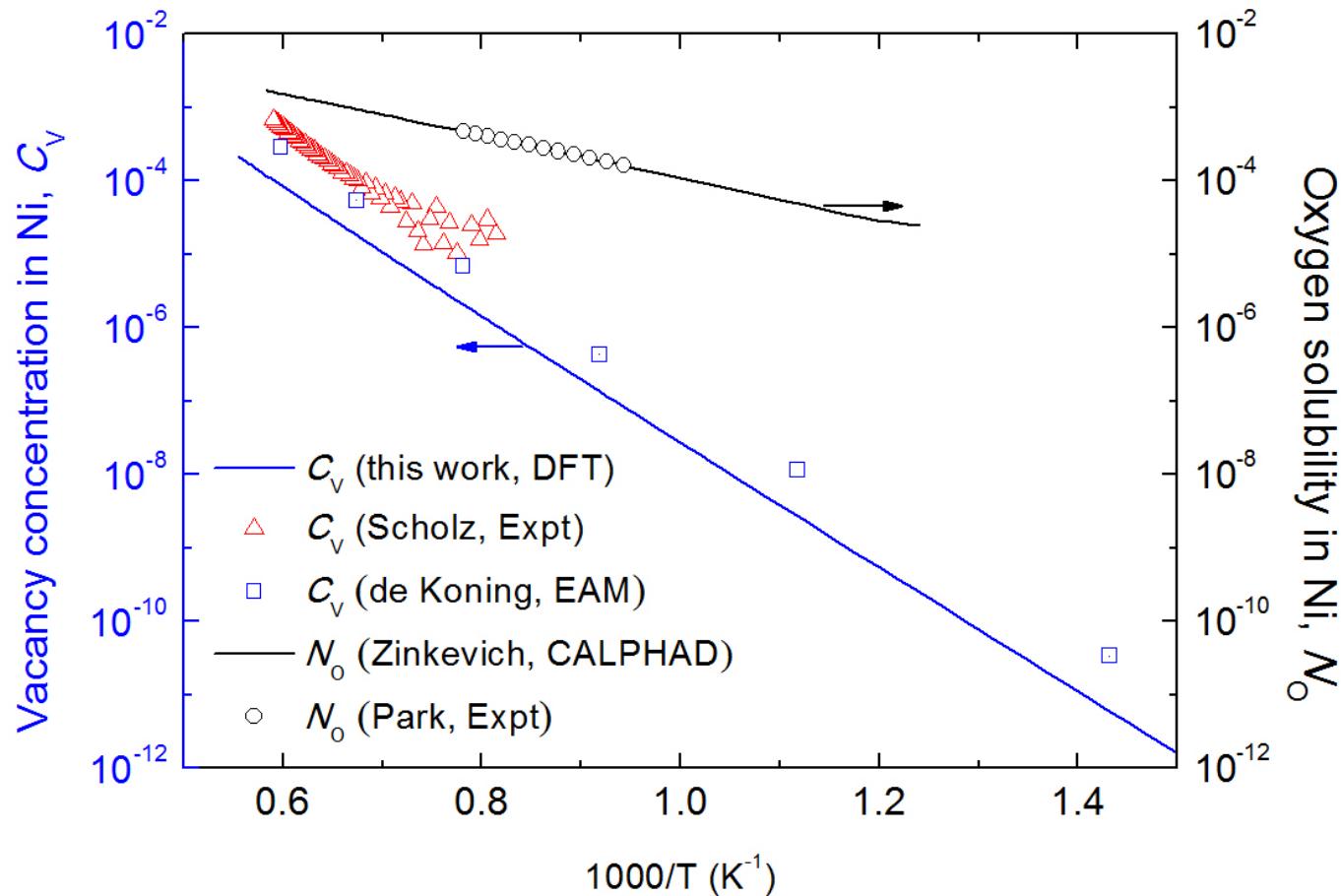
^aNi without vacancy.

^bNi with vacancy.

$$\begin{aligned} D_0 &: 8.8 \times 10^{-7} \sim 9.5 \times 10^{-3} \text{ (Expt: } 4.9 \times 10^{-6} \sim 8.93 \times 10^3\text{)} \\ Q &: 1.64 \sim 2.25 \text{ (Expt: } 1.7 \sim 4.28\text{)} \end{aligned}$$



Discussion – equilibrium O & Va contents in Ni



Vacancy concentration vs. oxygen solubility in Ni

PENNSTATE



Case 2 of 2

- Oxygen diffusion in bcc Fe with and without vacancy

Oxygen diffusion in bcc Fe

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Vacancy mechanism of oxygen diffusivity in bcc Fe: A first-principles study



S.L. Shang ^{a,b,*}, H.Z. Fang ^{a,b}, J. Wang ^{a,b,c}, C.P. Guo ^{a,b,d}, Y. Wang ^{a,b}, P.D. Jablonski ^{a,e}, Y. Du ^c, Z.K. Liu ^{a,b}

^a National Energy Technology Laboratory Regional University Alliance, U.S. Department of Energy, Pittsburgh, PA 15236, USA

^b Department of Materials Science and Engineering, The Pennsylvania State University, University Park, PA 16802, USA

^c State Key Laboratory of Powder Metallurgy, Central South University, Changsha, Hunan 410083, China

^d Department of Materials Science and Engineering, University of Science and Technology Beijing, Beijing 100083, China

^e National Energy Technology Laboratory, Department of Energy, Albany, OR 97321, USA

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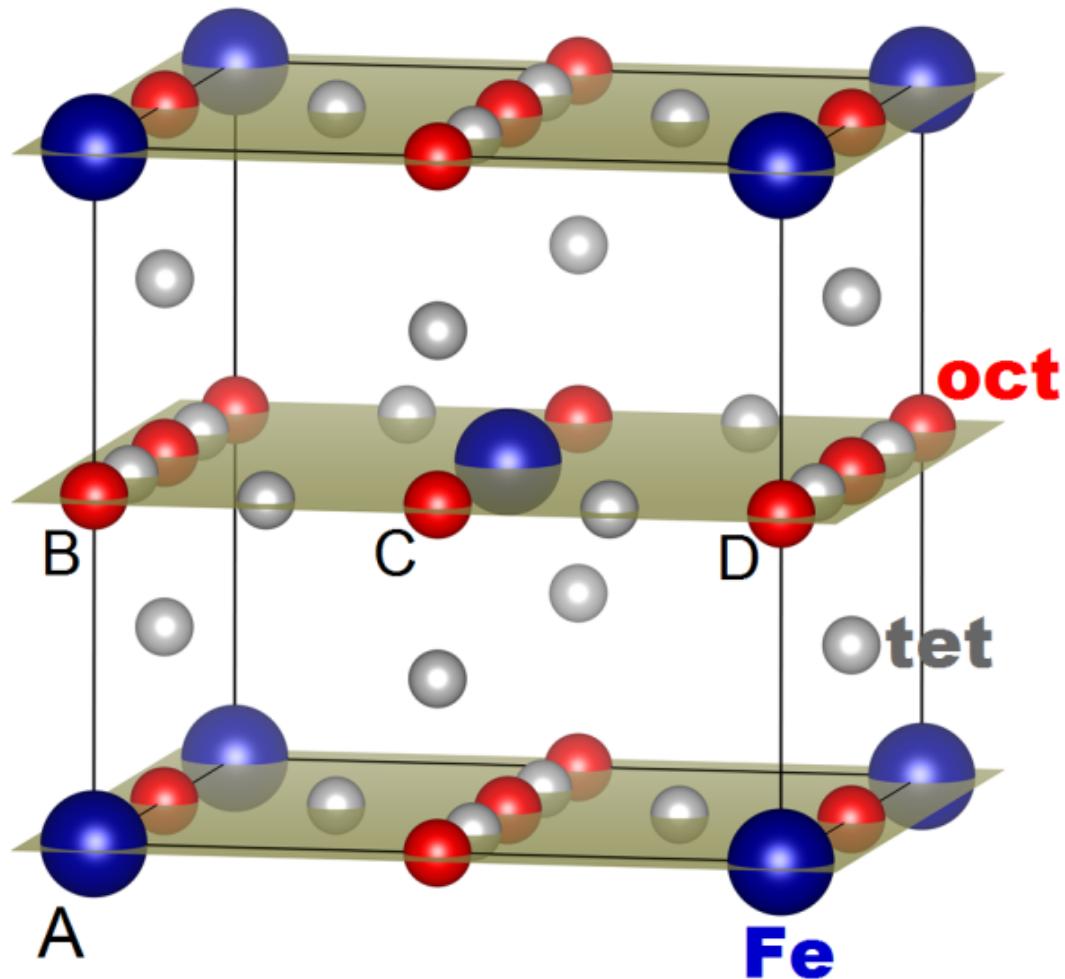
ABSTRACT

Diffusivity of interstitial oxygen (O) in bcc iron (Fe) with and without the effect of vacancy has been investigated in terms of first-principles calculations within the framework of transition state theory. Examination of migration pathway and phonon results indicates that O in octahedral interstice is always energetically favorable (minimum energy) with and without vacancy. It is found that vacancy possesses an extremely high affinity for O in bcc Fe, increasing dramatically the energy barrier (~80%) for O migration, and in turn, making the predicted diffusion coefficient of O in bcc Fe in favorable accord with experiments.

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bcc Fe & its *oct* & *tet* interstitial sites – structure



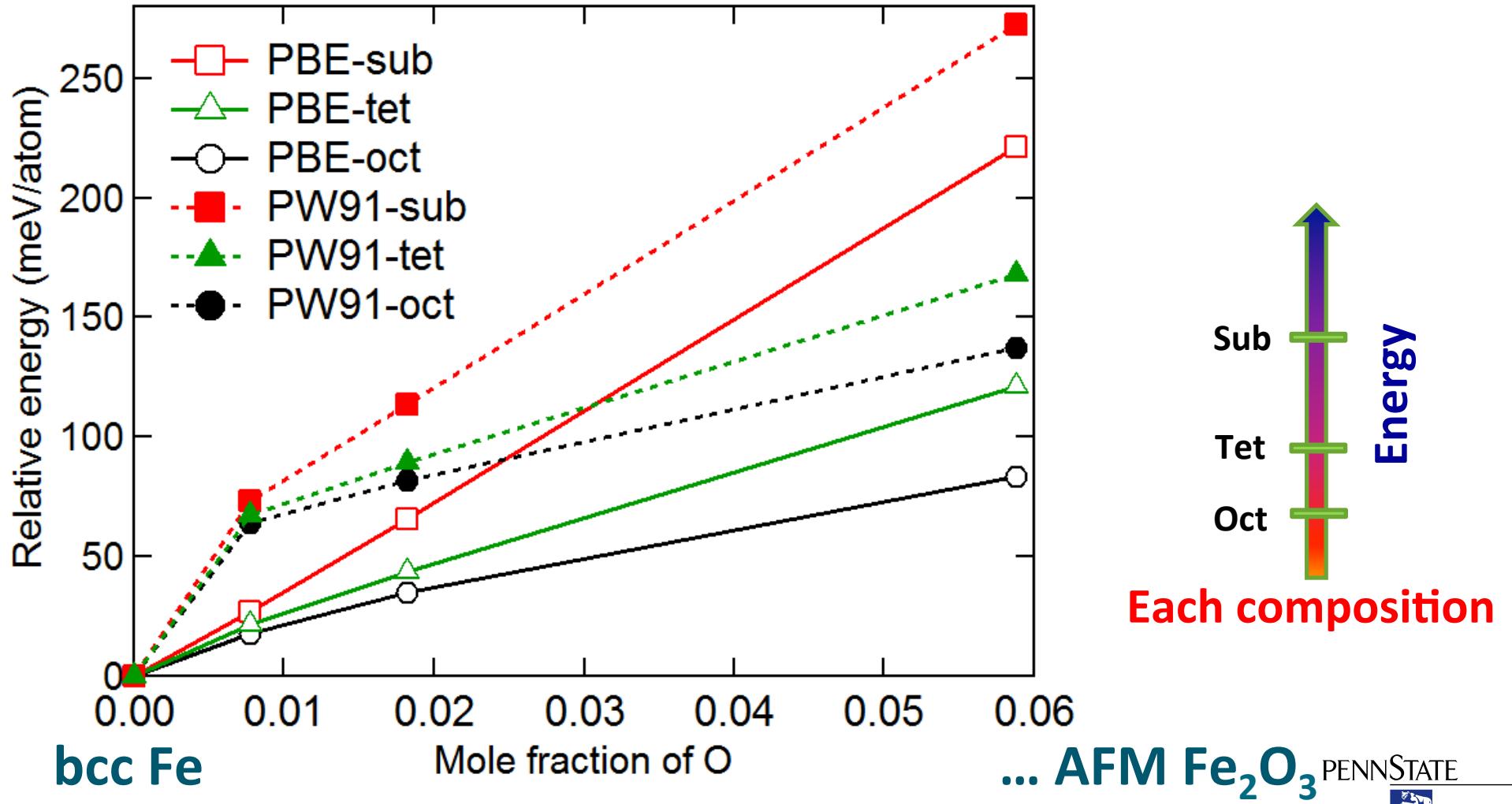
Fe: **2a** (0, 0, 0)

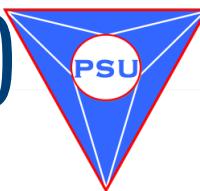
O-oct: **6b** (0 $\frac{1}{2}$ $\frac{1}{2}$)

O-tet: **12d** (0 $\frac{1}{2}$ $\frac{1}{4}$)

Site preference of O in bcc Fe

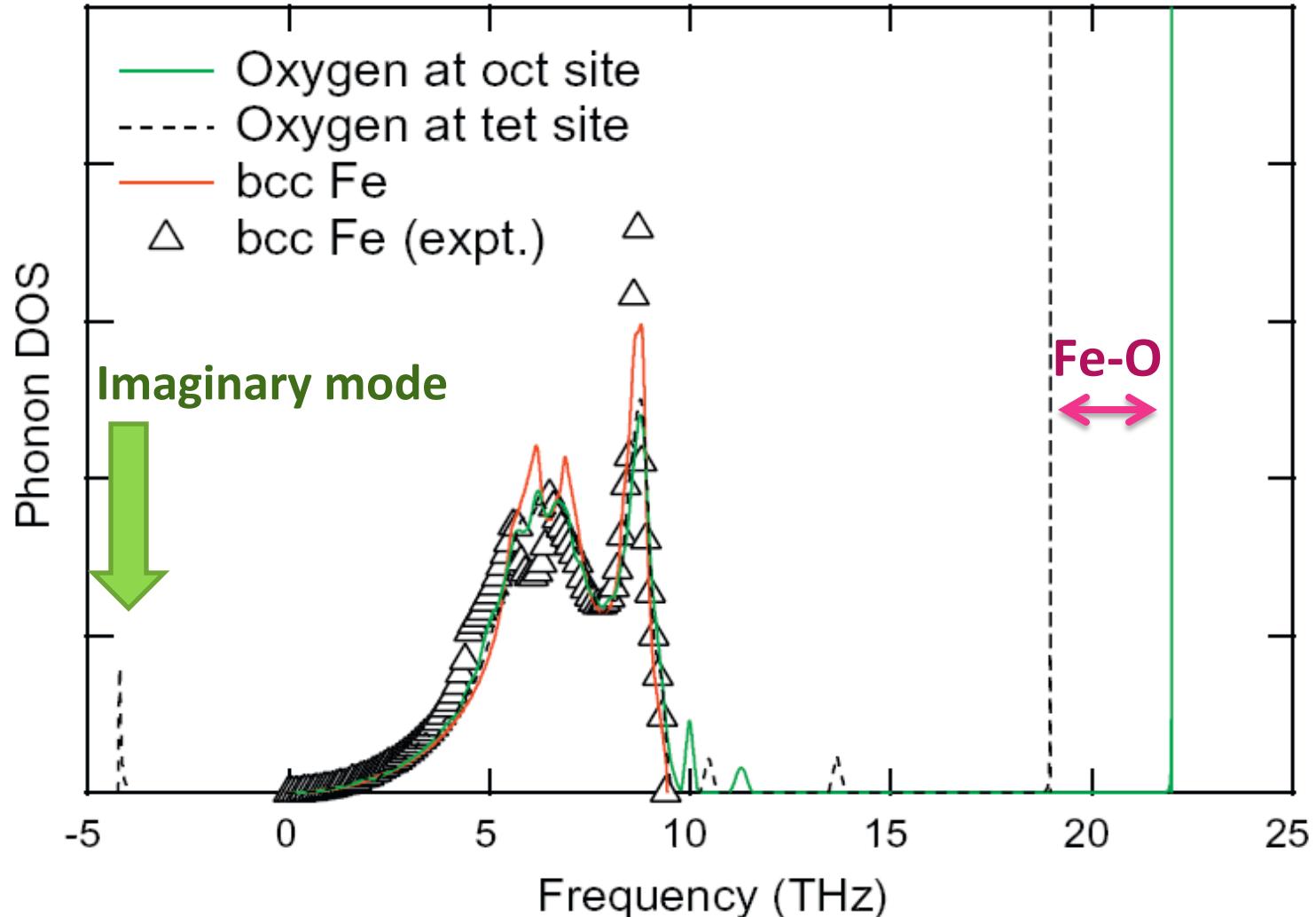
– w.r.t. bcc Fe and AFM $\alpha\text{-Fe}_2\text{O}_3$



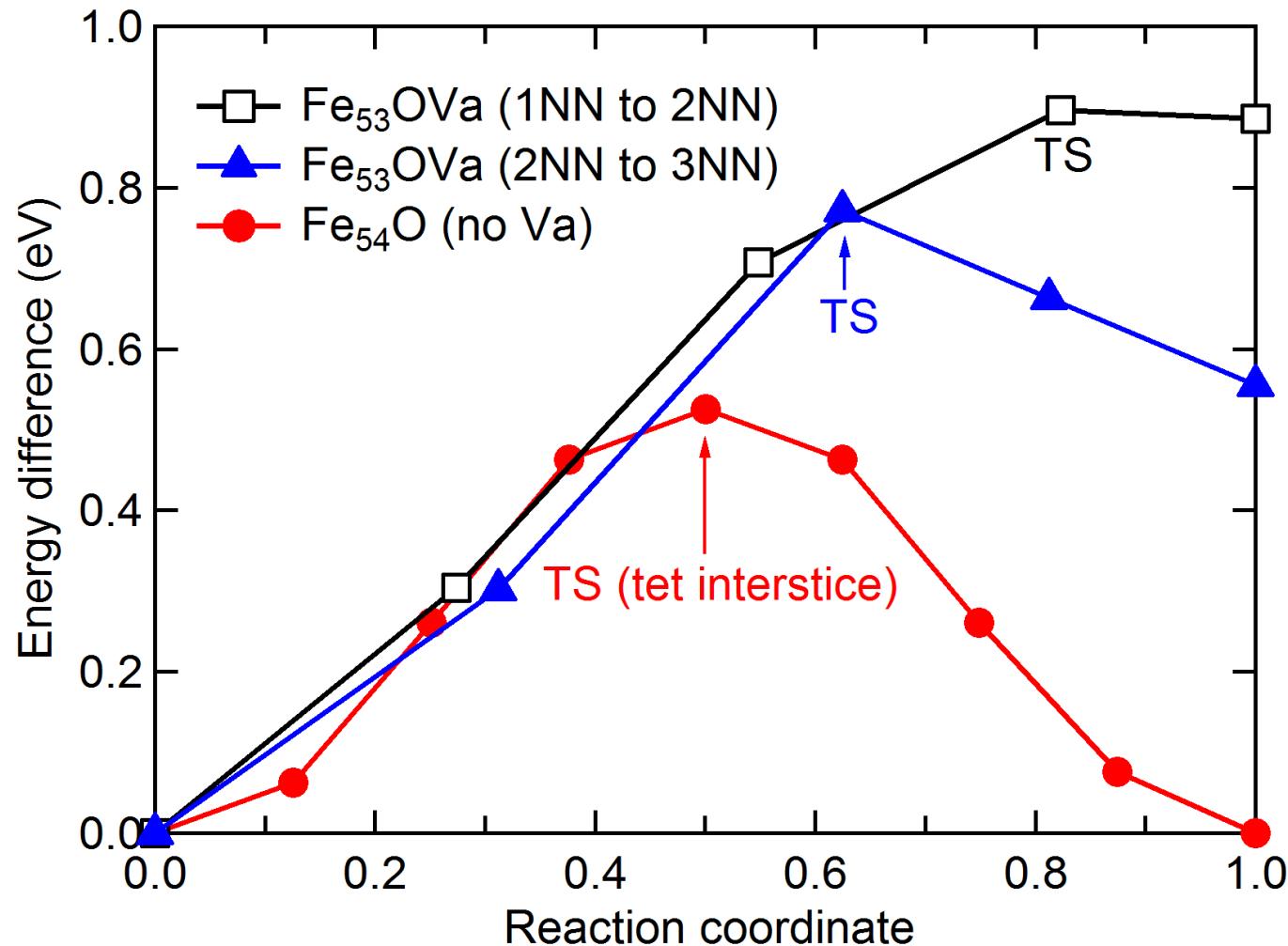


Phonon DOS of bcc Fe (without Va)

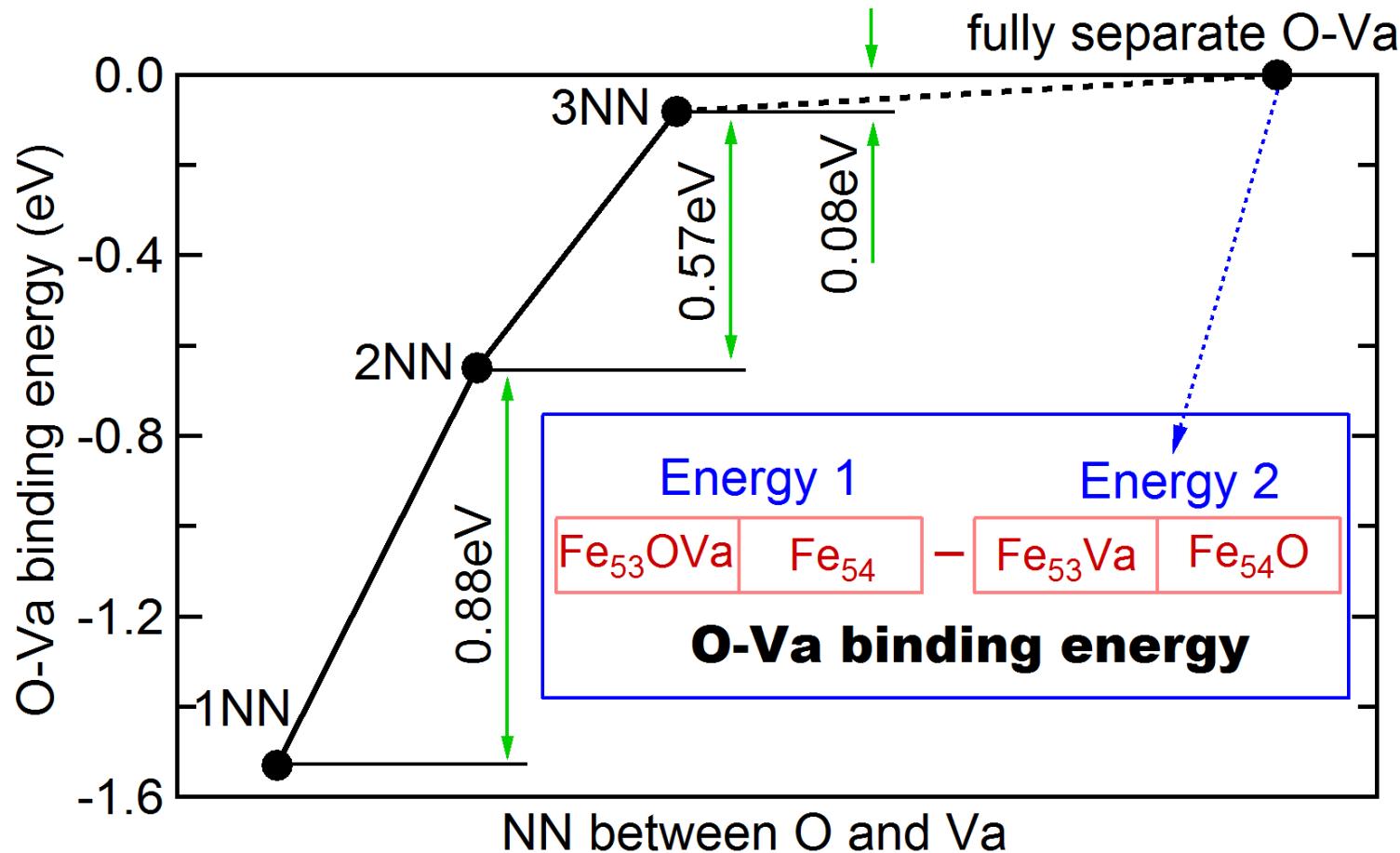
- O at oct & tet sites



Minimum energy pathways of O diffusion in bcc Fe **with/without** Va



0-Va binding energy in bcc Fe - using the $3 \times 3 \times 3$ supercell



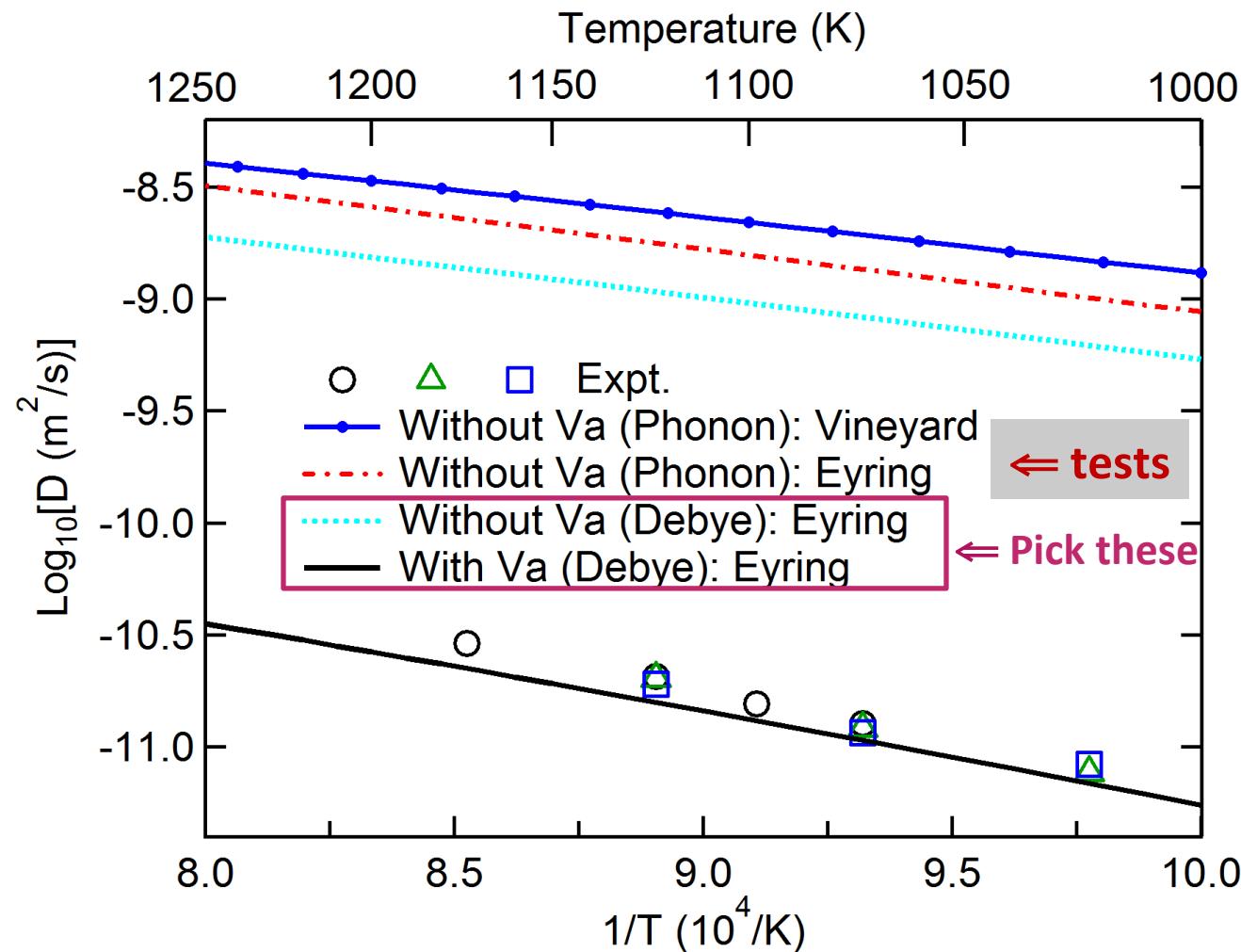
Activation energy of O diffusion in bcc Fe

- energy (kJ) between oct & tet sites

PW91: 2×2×2 cell	49.8	PBEsol: 3×3×3 cell	39.1
PW91: 3×3×3 cell	38.8	PBE-CINEB-bulk-relax (O-T)	51.4
PW91: 4×4×4 cell	38.9	PBE-CINEB-bulk-relax (O-T-O)	49.6
PBE: 2×2×2 cell	61.5	PBE-CINEB-bulk-fix (O-T-O)	52.1
PBE: 3×3×3 cell	47.0	PBE-CINEB-relax-Va (O-T-O)	98.8
PBE: 4×4×4 cell	47.8	Expt.: Swisher 1967	177
PBE: 3×3×3 (charged)	51.9	Expt.: Barlow 1969	167
LDA: 3×3×3 cell	16.1	Expt.: Takada 1986	86±6
			90±7

Arrhenius plots of O diffusion in bcc Fe

with/without Va $D = D_0 \exp(-Q/kT)$





Arrhenius plots of O diffusion in bcc Fe with/without Va (details)

$$D = D_0 \exp(-Q/kT)$$

System	ΔE_m	D_0	Q
Fe ₅₄ O (without Va) ^a	0.526 ^b	3.75×10^{-7} (e)	0.49 ^e
	0.503 ^c	5.72×10^{-7} (f)	0.56 ^f
		2.86×10^{-7} (g)	0.54 ^g
Fe ₅₃ OVa (1NN to 2NN) ^a	0.897 ^b	0.63×10^{-7} (g)	0.80 ^g
	0.892 ^d		
Fe ₅₃ OVa (2NN to 3NN) ^a	0.772 ^b		
Expt: O in dilute Al-Fe ^h		$(1.79_{-0.89}^{+1.76}) \times 10^{-7}$	0.89 ± 0.06
Expt: O in dilute Si-Fe ^h		$(2.91_{-1.57}^{+3.40}) \times 10^{-7}$	0.93 ± 0.07
Expt: O in dilute Ti-Fe ^h		$(3.78_{-1.83}^{+3.55}) \times 10^{-7}$	0.95 ± 0.06
Expt: O in dilute Si-Fe ⁱ		$(400 \pm 100) \times 10^{-7}$	1.73 ± 0.01
Expt: O in dilute Al-Fe ^j		37.2×10^{-7}	1.01
Expt: estimated values ^k			0.98 ± 0.1

Summary

- Exploration of site preference, diffusion pathways & barriers, and diffusivities of oxygen
- Existence of strong O-Va binding energy in metals
- Va-modified interstitial diffusion mechanism proposed to quantitatively predict diffusivity of oxygen in metals
- Case studied in fcc Ni and bcc Fe

