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(54) **OPTICAL FREQUENCY COMB-BASED
COHERENT LIDAR**

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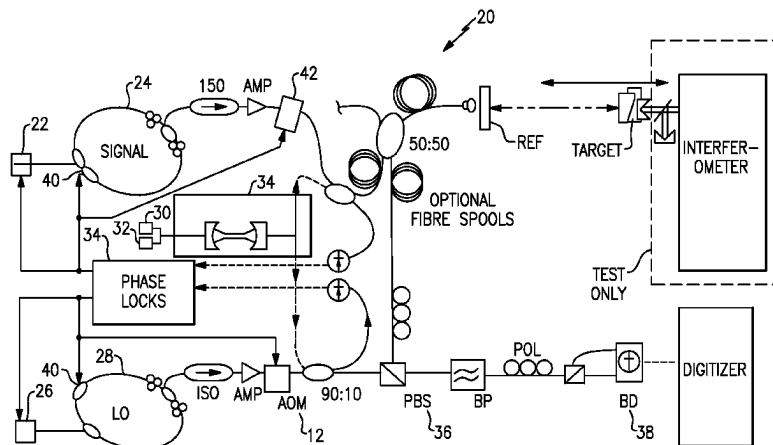
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(57) **ABSTRACT**

A coherent laser radar that uses two coherent femtosecond
fiber lasers to perform absolute ranging at long distance. One
coherent femtosecond fiber lasers acts as a source and the
other as a local oscillator for heterodyne detection of the
return signal from a cooperative target. The system simulta-
neously returns a time-of-flight range measurement for
coarse ranging and an interferometric range measurement for
fine ranging which is insensitive to spurious reflections that
can cause systematic errors. The range is measured with at
least 3 μm precision in 200 μs and 5 nm precision in 60 ms
over a 1.5 m ambiguity range. This ambiguity range can be
extended to 30 km through reversal of signal and LO source
roles.

13 Claims, 7 Drawing Sheets



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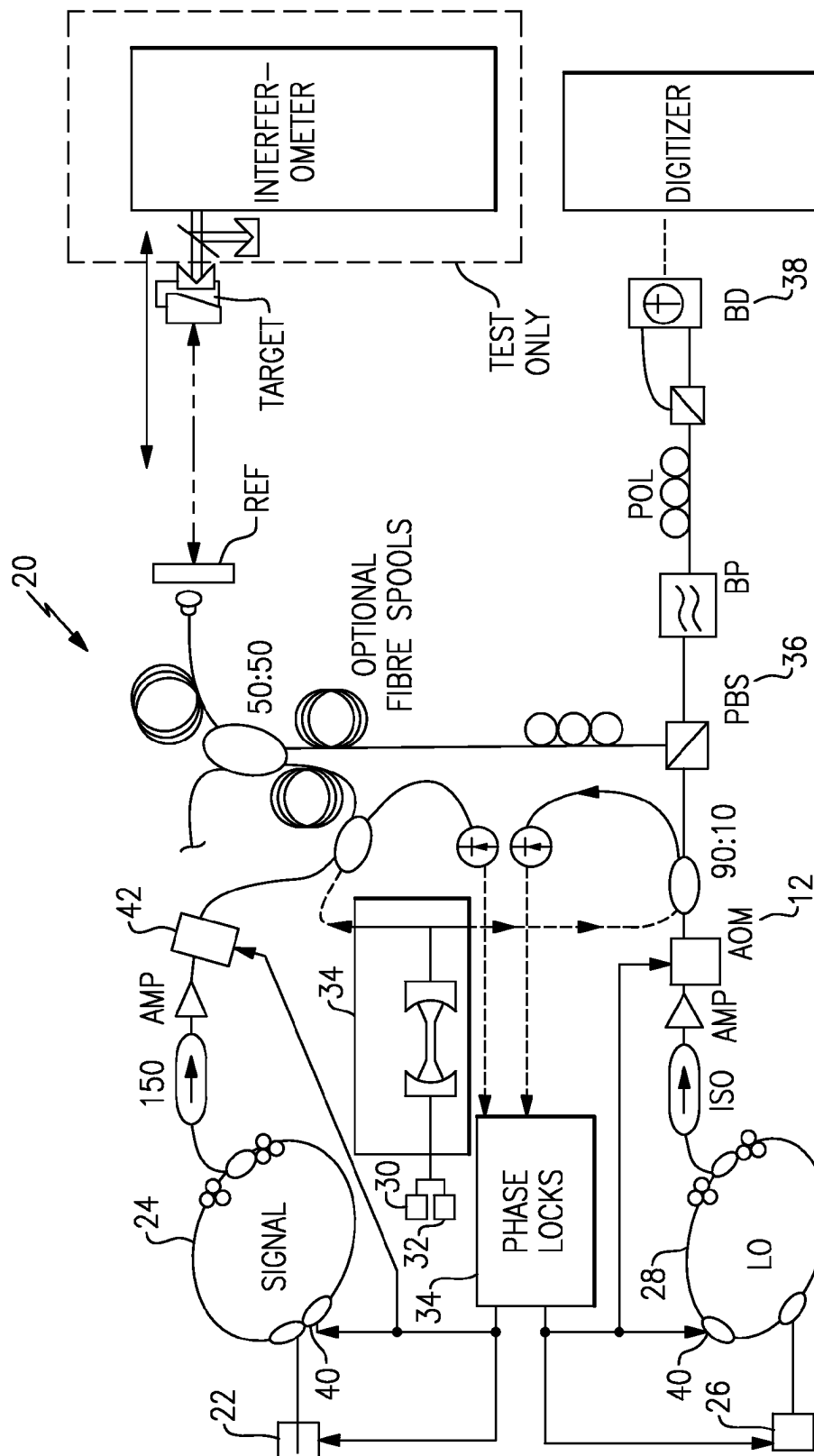


FIG. 1

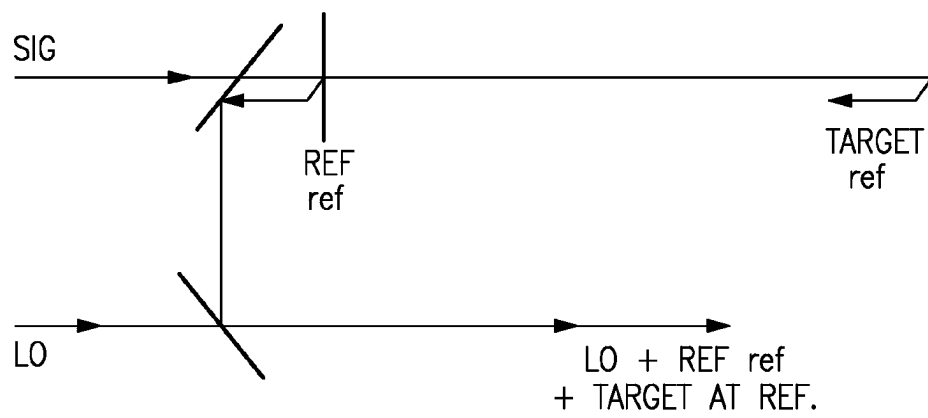


FIG. 2A

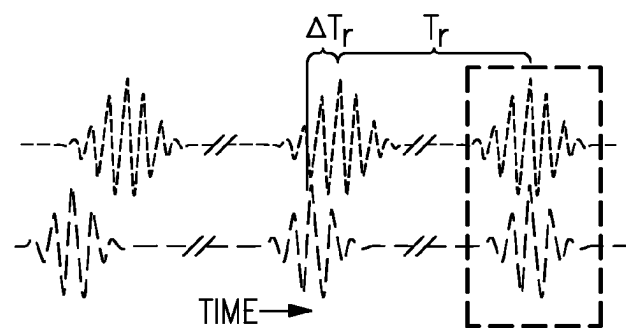


FIG. 2B

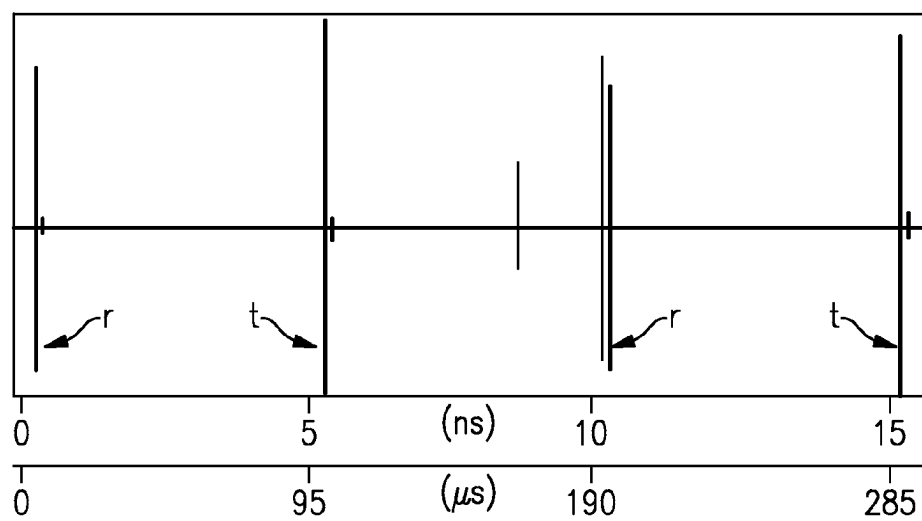
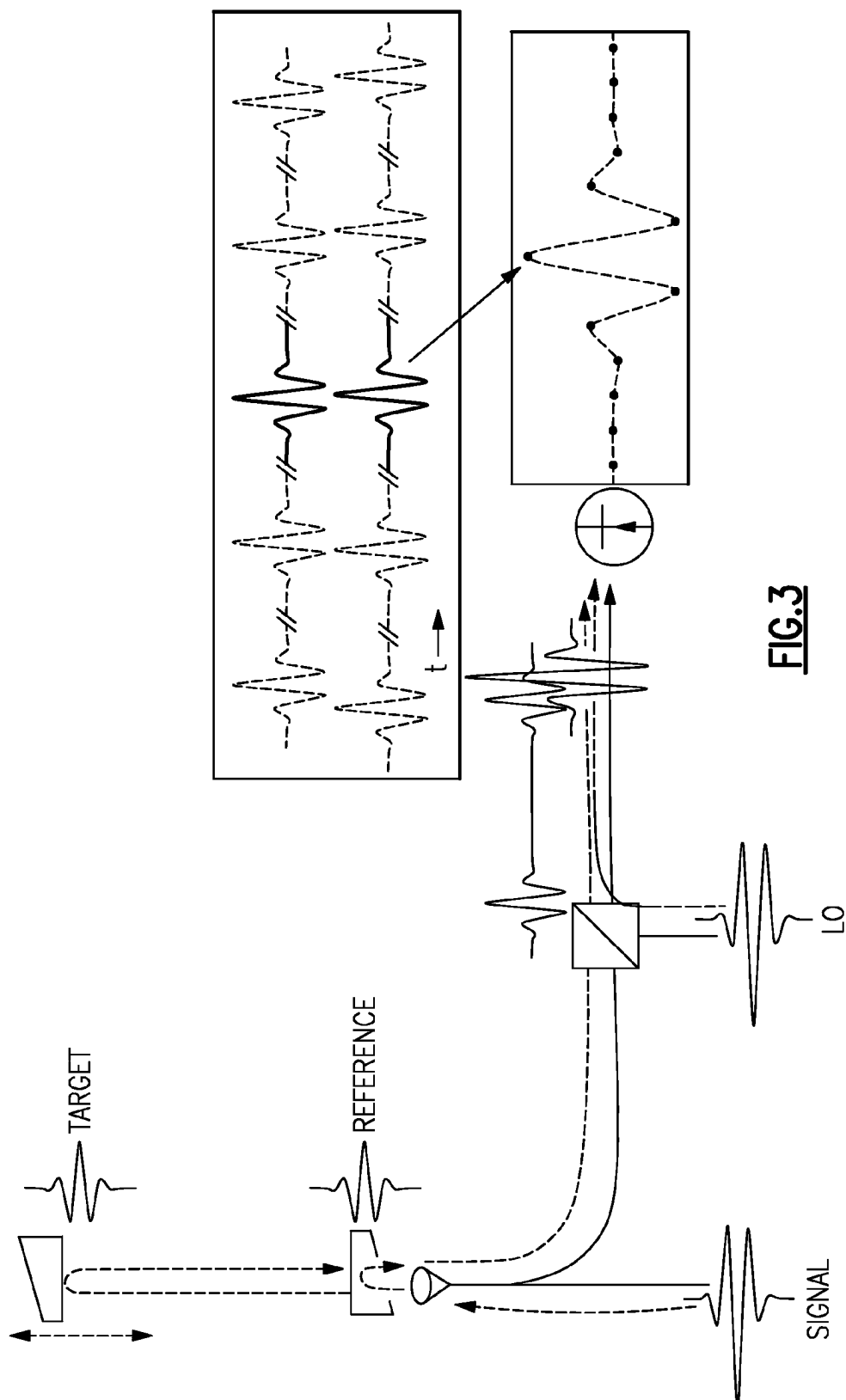
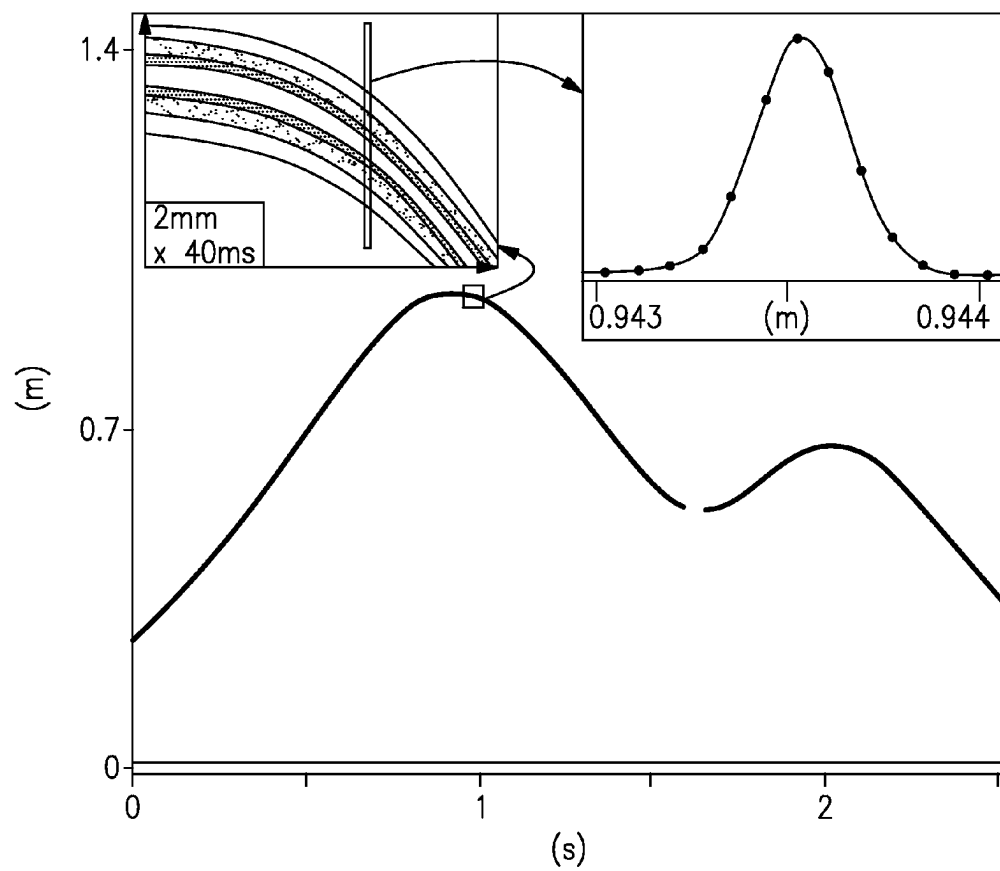
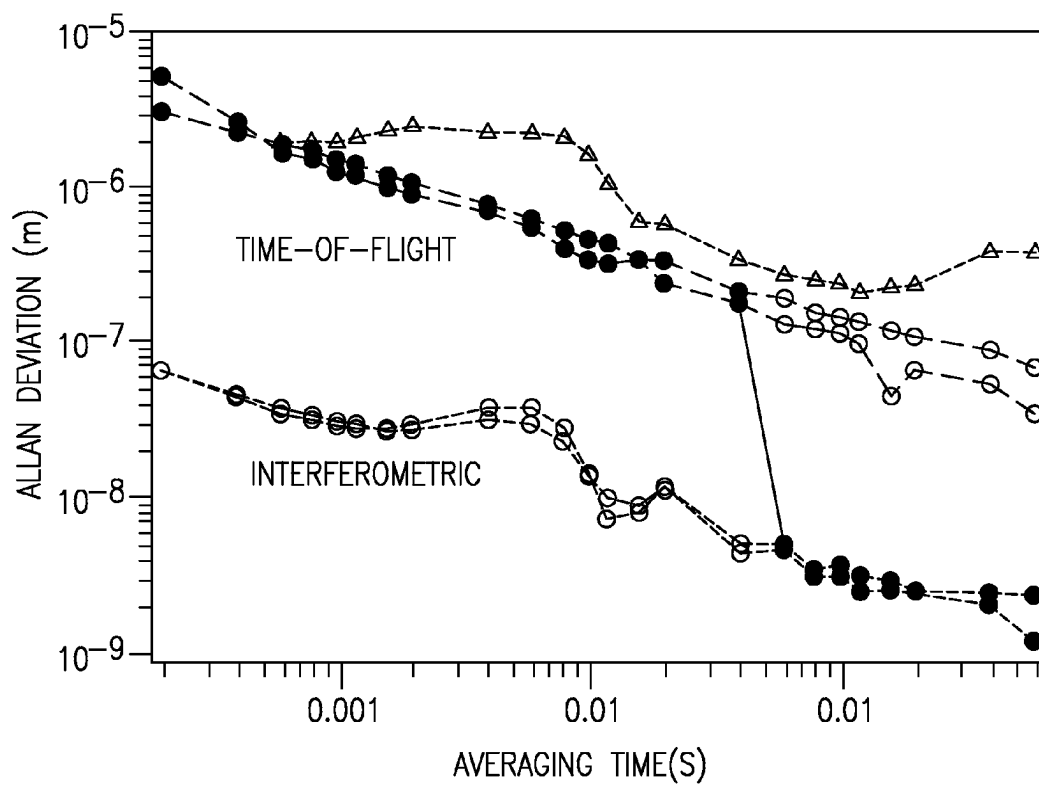
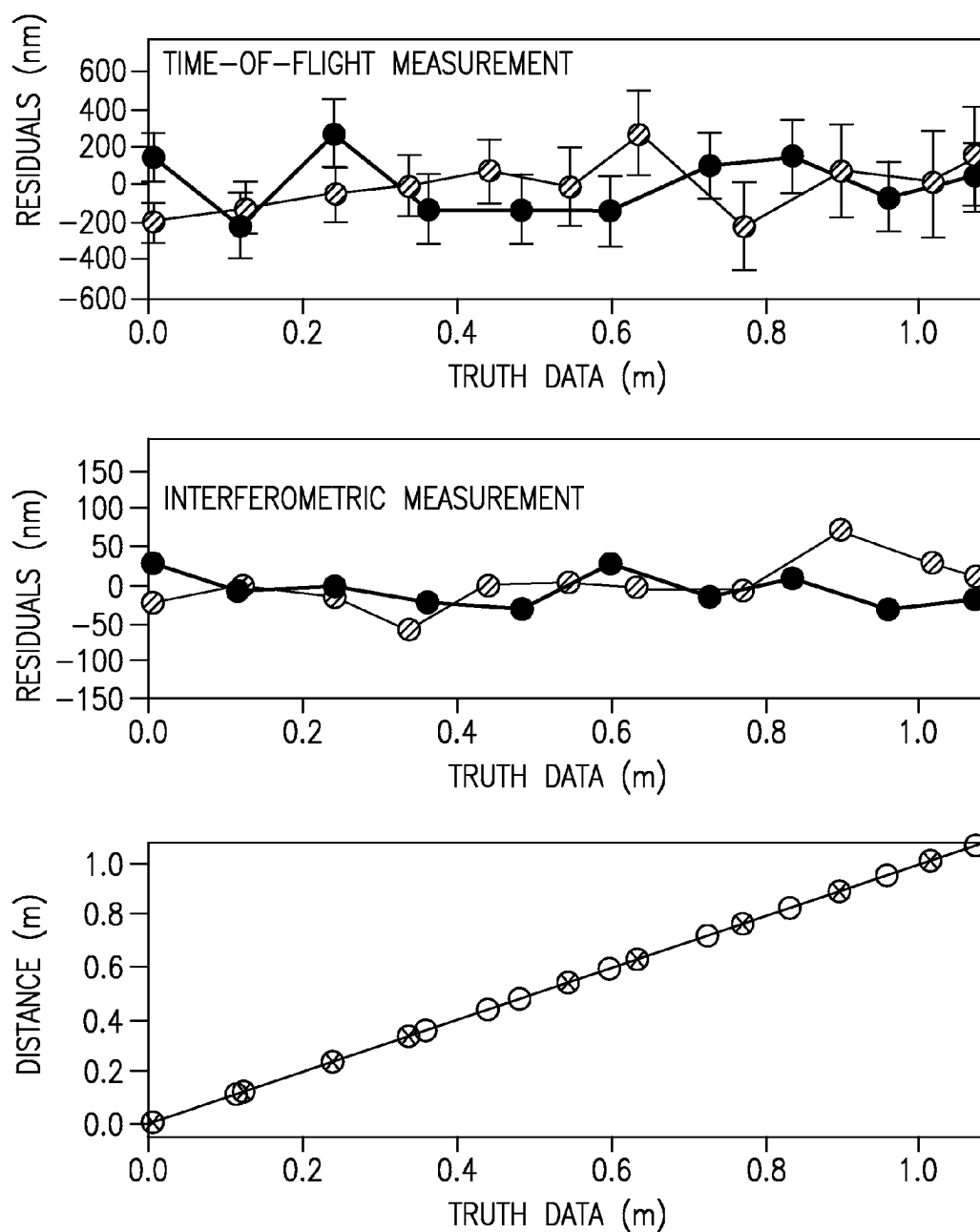


FIG. 2C



**FIG.4**

**FIG.5**

**FIG.6**

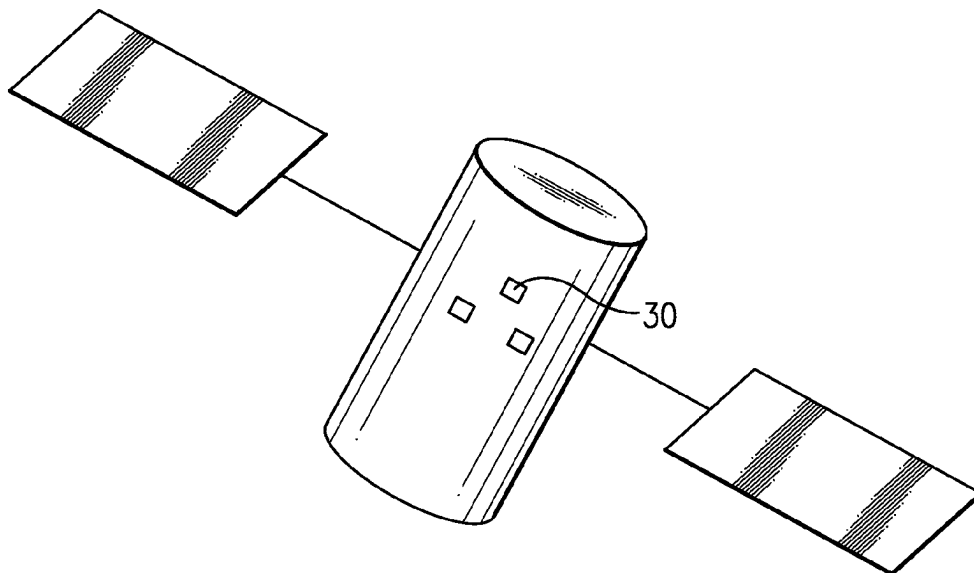


FIG. 7

OPTICAL FREQUENCY COMB-BASED COHERENT LIDAR

The present disclosure claims priority to U.S. Provisional Patent Application No. 61/347,095, filed May 21, 2010.

BACKGROUND

The present disclosure relates to a comb-based coherent LIDAR that provides a unique combination of precision, speed and large ambiguity range.

In intrasatellite ranging, and similarly in manufacturing applications, there are three critical parameters: precision/accuracy, ambiguity range and update rate. High precision is particularly important in maintaining the pointing during formation flying of satellites; for example, coherent combining of 1-m sub-apertures to form a synthetic aperture of 100-m diameter requires a relative pointing accuracy of less than ($\lambda/100$ m) rad for the sub-aperture on each satellite, which in turn requires distance measurements at the sub-aperture edges with less than $\lambda \times (1 \text{ m}/100 \text{ m})$ accuracy, or a few nanometers at optical wavelengths. The ambiguity range characterizes the measurement range window; longer distances are aliased back to within the ambiguity range. Larger ambiguity range requires less a priori distance knowledge. Finally, fast millisecond-scale update rates are needed for effective feedback.

Many of these requirements push or exceed the capabilities of current 'stand off' ranging technology. Generally speaking, laser ranging is the determination of the phase shift on a signal after traversing a given distance. Crudely, shorter-wavelength signals offer greater resolution, and longer wavelength signals offer greater ambiguity range. For instance, the widely used continuous-wave (c.w.) laser interferometer measures the phase of optical wavelengths to achieve sub-nanometer resolution. However, measurements are limited to relative range changes as the ambiguity range equals half the laser wavelength. Alternatively, laser radar (LIDAR) measures distance through pulsed or radio-frequency (rf)-modulated waveforms. For pulsed systems, the time-of-flight is measured

These systems offer large ambiguity ranges but with ~ 50 - $100 \mu\text{m}$ resolution. Multiwavelength interferometry (MWI) combines measurements at several optical wavelengths, which effectively generates a longer 'synthetic wavelength', and therefore a reasonable ambiguity range while maintaining sub-wavelength resolution. However, these systems are vulnerable to systematic errors from spurious reflections, and extending the ambiguity range beyond a millimeter can require slow scanning. Nevertheless, with extensive care in minimizing spurious reflections, the MSTAR system has successfully used MWI for sub-micrometer ranging. Femtosecond optical frequency combs offer an intriguing solution to intrasatellite ranging. From the early work by Minoshima, combs have been incorporated into precision ranging systems using various approaches.

Although frequency combs are normally considered in the frequency domain where they produce a comb of well-defined, narrow linewidth optical frequency lines, these same sources can also be viewed in the time-domain where they produce a train of well-defined, coherent optical pulses. These optical pulses can be very short in duration, i.e. have a large bandwidth, and can be arranged to have a high carrier phase coherence with an underlying optical cw "clock" laser. They very much resemble a coherent RADAR pulse train, except that the carrier frequency is shifted up into the optical region and their bandwidth can be significantly larger. As a

consequence, these sources are interesting for high-resolution coherent LIDAR systems.

One challenge of taking full advantage of these high-bandwidth, coherent sources in a coherent LIDAR system lies in effectively detecting the return signal. The optical pulses have very high bandwidth (i.e. THz or greater) and standard direct or heterodyne detection would require an equivalently large bandwidth receiver.

SUMMARY

A LIDAR system and method where a second coherent frequency comb is heterodyned against the return signal and effectively down-converts the full bandwidth of the return signal to baseband, where it can be detected by relatively low bandwidth detectors. In the frequency domain, this picture is equivalent to massively parallel heterodyne detection between the "teeth" of the signal and local oscillator (LO) frequency comb. In the time domain, it is equivalent to linear optical sampling. This technique allows measurement of absolute range with 1.5 meter ambiguity range to high precision by "handing over" an effective time-of-flight measurement to an interferometric measurement. The time-of-flight measurement provides $3 \mu\text{m}$ precision in $200 \mu\text{s}$; with averaging the measurement dropped below a quarter wavelength after 60 ms at which time the interferometric range measurement gave 3 nm precision. Such a LIDAR system may be suitable for precision positioning of satellites in a coherent formation

BRIEF DESCRIPTION OF THE DRAWINGS

Various features will become apparent to those skilled in the art from the following detailed description of the disclosed non-limiting embodiment. The drawings that accompany the detailed description can be briefly described as follows:

FIG. 1 is a general view that schematically illustrates a Light Detection and Ranging (LIDAR) system in an experimental set-up. Two erbium fibre frequency combs are each phase-locked to two c.w. reference lasers at 1,535 and 1,550 nm. The output of the frequency combs can be viewed in the time domain as a train of pulses or in the frequency domain as a regularly-spaced frequency comb. One fiber frequency comb acts as a signal frequency comb and one acts as a Local Oscillator (LO) frequency comb. The signal comb outputs a signal pulse train and the LO comb outputs an LO pulse train in time. For each laser an intercavity piezo-electric transducer and external acousto-optic modulator (AOM) provide modulation for one lock, and pump current modulation is sufficient for the second. A c.w. interferometer is used to monitor the relative target position only to provide truth data and is not part of the comb based measurement system. The signal pulse trains are combined with the LO on a polarizing beamsplitter (PBS), optically filtered with a 3-nm bandpass (BP) at 1,562 nm, directed to a balanced detector (BD), and finally digitized at 14 bits synchronously with the LO. The total detected signal power per reflection is approximately 0.4 mW, or 4 fJ per pulse. ISO, isolator; AMP, erbium fibre amplifier; POL, polarization control; 90:10 and 50:50 are splitter ratios. Optional (1.14-km) fibre spools are included in two configurations, matched pairs and bidirectional;

FIG. 2A is a schematic view of a high-repetition-rate 'signal' source which transmits pulses that are reflected from two partially reflecting planes (glass plates): the reference (r) and the target (t). The reference is a flat plate and yields two reflections, the first of which is ignored. Distance is measured

as the time delay between the reflections from the back surface of the reference flat and the front of the target.

FIG. 2B shows the signal pulses detected through linear optical sampling against a local oscillator (LO). The LO generates pulses at a slightly different repetition rate. Every repetition period (T_r), the LO pulse ‘slips’ by ΔT_r relative to the signal pulses, and samples a slightly different portion of the signal. Consecutive samples of the overlap between the signal and LO yield a high-resolution measurement of the returning target and reference pulses.

FIG. 2C is the measured voltage out of the detector in both real time (lower scale, 10 ns per data point) and effective time (upper scale, 500 fs per data point) for a target and reference plane separated by 0.76 m. A full ‘scan’ of the LO pulse across the signal pulse is accomplished every ~ 200 microsec in real time and every ~ 10 ns in effective time. Two such scans are shown to illustrate the fast, repetitive nature of the measurement. Also seen are several spurious reflections of the launch optics;

FIG. 3 schematically illustrates a Light Detection and Ranging (LIDAR) system;

FIG. 4 is a real-time image of the range versus time for a moving target (shown in false colour). Every 200 microsec, the system scans the entire 1.5-m ambiguity range of the system. Sequential scans are stacked horizontally to yield an image tracking the target motion. The signal at zero distance is the reference plane, and the moving signal represents the target. The upper left inset shows an expanded. The upper right inset shows a cross-section of the return signal (magnitude squared of the detected electric field). The signal width is set by the 0.42 THz signal $1/e^2$ bandwidth;

FIG. 5 shows the Allan deviation (precision) of the distance measurement versus averaging time. Three different measurements are shown: (i) an ~ 1 m reference-to-target distance without any fibre delay (lower black circles), (ii) an ~ 1 m reference-to-target distance with a 1.14-km fibre delay (upper black circles) and (iii) an ~ 1.14 km distance across a bidirectional fibre spool (triangles). The Allan deviation is evaluated over a 5-s data run that is composed of a series of individual data points at the 200 microsec update rate. For both reference-to-target distance measurements (with and without the fibre delay), both the time-of-flight and interferometric measurements are given. Solid circles show the experimental precision, which experiences a significant jump around 60 ms when the time-of-flight measurement is sufficiently stable compared to $\lambda_c/4$ and can be combined with the interferometric measurement. At shorter averaging times the interferometric data are still available, but provide only relative position information. For the 1.14-km distance measurement across the bidirectional fibre spool (red triangles), the precision of the time-of-flight distance is scaled to the group velocity of air. In this case, the uncertainty is dominated by actual length changes in the fibre and a 300-nm flicker floor from Rayleigh backscatter;

FIG. 6 are residuals of the measured time-of-flight and interferometric range measurements versus truth data from the commercial c.w. interferometer. The averaging period is 60 ms. Error bars are the standard deviation of the mean over the 60-ms period. Data taken with the fibre spool delay line are shown in black, and data without the spool are shown in crosshatch; as seen, the addition of the fibre delay has a negligible effect on the measurement. Although the interferometric data are much quieter, they become meaningful only when combined with the time-of-flight data to resolve the 780-nm range ambiguity; and

FIG. 7 is a schematic view of a satellite with multiple reference points such that a pitch and yaw of the satellite can

be determined through trilateration, where the angle is calculated from the distance to several reference points spanning the target.

DETAILED DESCRIPTION

FIG. 1 schematically illustrates a Light Detection and Ranging (LIDAR) system 20. The system 20 generally includes a first mode locked laser 24, also referred to as the signal comb, which transmits a signal pulse train reflected from a reference and a target (also illustrated in FIG. 3) and a second mode locked laser 28, also referred to as the LO comb, that generates a local oscillator (LO) pulse train. The combs can be phase locked to a first continuous wave laser 30 and a second continuous wave laser 32 that are themselves stabilized to an optical reference cavity 34. It should be understood that various laser types and systems may alternatively or additionally be provided. Furthermore, if only micrometer ranging is required (from the time-of-flight data), then the cw lasers need not be used to stabilize the optical carrier of the comb and the repetition rates can be monitored rather than phase-locked.

The pulsed nature of the comb is combined with the coherence of the carrier, allowing for a time-of-flight measurement simultaneously with an interferometric measurement based on carrier phase. This is implemented with the dual coherent frequency combs to achieve a nanometer level of precision with an ambiguity range of 1.5 m in 60 ms at low light levels and with high immunity to spurious reflections. The ambiguity range is readily extended to 30 km by reversing the role of the signal and LO combs 24, 28. The time-resolved signal also permits measurements between multiple reference planes in a single beam path. This host of features is unavailable in any other single system.

The combs 24, 28 may be of fiber ring laser design or other design. Signal pulse trains are combined with the LO pulse trains on a polarizing beamsplitter (PBS) 36, optically filtered with a 3-nm bandpass (BP) at 1,562 nm, directed to a balanced detector (BD) 38, then digitized. The total detected signal power per reflection is $\sim 0.4 \mu\text{W}$, or 4 fJ per pulse. ISO, isolator; AMP, erbium fibre amplifier; POL, polarization control; 90:10 and 50:50 are splitter ratios. Optional (1.14-km) fibre spools may be included in two configurations, in matched pairs and bidirectional.

The signal comb 24 and the LO comb 28 are each phase-locked to the two continuous wave reference lasers 30, 32 at, for example, 1,535 nm and 1,550 nm. Both cw lasers 30, 32 are stabilized to the same high-finesse optical reference cavity 34 via, for example, a Pound Dreyer Hall lock (PDH). The PDH may be implemented with fiber-optic components. With this approach, the optical reference cavity 34 effectively provides the overall timebase. If very high accuracy is needed, the frequency of the 1560 nm cavity-stabilized laser is measured by a fully self-referenced frequency comb (not shown) linked to a hydrogen maser. The two combs are phase-locked to the two cw reference lasers as described next. The lock setups for the two combs 24, 28 are essentially identical. The intercavity piezo-electric transducer (PZT) 40 and external acousto-optic modulator (AOM) 42 provide feedback for one lock, and feedback to the pump current 22, 26 is sufficient for the second. The comb light from each comb is combined with light from the stabilized cw lasers 30, 32, on the 90:10 splitter, with the 90% output directed to the target. The 10% port with the remaining comb light and the two cw lasers is spectrally separated on a 1 nm, 2-channel, grating filter and detected on two 100 MHz detectors. Each detector detects a beat signal between the comb and one cw laser.

The 1560 nm light stabilizes the comb carrier frequencies and the 1535 nm light stabilizes the comb repetition rates. Because optical bandwidth is relatively small compared to carrier frequency, the carrier stabilization lock may be the more critical of the two locks. For the carrier stabilization lock, the rf beat signal is divided by 8 to increase the capture range and then a digital phase lock loop (PLL) is used to lock the divided signal to a 10 MHz reference. The error signal is passed to a voltage controlled oscillator that drives the AOM and provides the high bandwidth (~100 kHz) feedback modulation. A second error signal is fed to an additional loop filter followed by a high voltage amplifier driving the intracavity PZT. This loop filter is tuned for high gain at DC and provides a slow, high dynamic range feedback. When locked, the integrated phase noise is <0.5 radians from 0.1 Hz-3 MHz.

The error signal for the lock at 1535 nm is generated using a mixer and tunable frequency synthesizer. The error signal is low pass filtered at 1.5 MHz and fed into a loop servo. The output of the loop servo is used to modulate the pump laser 22, 26 power to the combs. Our current feedback bandwidth is ~10 kHz limited by the laser dynamics combined with the erbium response. Although 10 kHz is sufficient, phase compensation could achieve still higher bandwidths.

With reference to FIGS. 2A-2C, the pair of stabilized femtosecond laser frequency combs 24, 28 have pulse trains of slightly different repetition periods (T_r and $T_r - \Delta T_r$). The pulse train from the signal comb 24 serves as the 'signal' source and samples a distance path defined by reflections off the target and the reference (also schematically illustrated in FIG. 3). The pulse train from the LO comb 28 serves as a broadband local oscillator (LO), and recovers range information in an approach equivalent to linear optical sampling. That is, a heterodyne cross-correlation is provided between the signal and LO combs. The heterodyne detection provides shot-noise limited performance so that even weak return signals can be detected and the information in the carrier phase is retained. Measurements can be made with slow detectors and electronics (50-100 MHz) and are readily mapped back into the original femtosecond timescale.

An entire scan of the signal return, shown in FIG. 2C, is completed every $T_{update} = T_r^2 / \Delta T_r$.

In one non-limiting embodiment, the combs 24, 28 operate at repetition rates of 100.021 and 100.016 MHz, giving $T_r \approx 10$ ns, $T_{update} = 1/5,190$ Hz ≈ 200 μ s and $\Delta T_r \approx 0.5$ ps. A 3-nm band-pass filter (BP; FIG. 1) limits the transmitted optical bandwidth to less than $1/(4\Delta T_r)$ in order to meet the Nyquist condition for sampling.

Given the digitized signal in FIG. 2C, the distance between target and reference reflections is calculated by the use of Fourier processing. Mathematically, the LO and signal electric field pulse trains are: $\sum_n e^{in\theta_{LO}} E_{LO}(t - nT_r)$ and $\sum_n e^{in\theta_S} E_S(t - n(T_r - \Delta T_r))$, respectively, where $E_{LO(S)}$ is the electric field of a single pulse, n is the pulse index and $\theta_{LO(S)}$ is the carrier-envelope offset phase.

For the n th pulse the detected voltage signal is proportional to the temporal overlap between the LO and delayed signal pulses, given by $V(t_{eff}) = \int E_{LO}^*(t) [E_S(t + t_{eff} - \tau_r) + e^{i\psi} E_S(t + t_{eff} - \tau_t)] dt$, where the effective time is $t_{eff} = n\Delta T_r$, ψ accounts for the π differential phase shift upon reflection as well as the relative Gouy phase, τ_r and τ_t are the delays on the reference and target pulses, respectively, and we assume for simplicity that $\theta_{LO} = \theta_S$.

The relative delay $\tau = \tau_t - \tau_r$ between the target and reference reflection peaks in FIG. 2C, are time gated separate contributions to $V(t_{eff})$ from the reference and target(s) to find $V^*(t_{eff})$ and $V^t(t_{eff})$. The Fourier transforms of the two are related by $\tilde{V}^t(\nu) = e^{i\phi(\nu) + i\psi} \tilde{V}^r(\nu)$, with the relative spectral phase of $\phi(\nu)$

$= 2\pi\nu\tau$. Converting from τ to measured distance L , and including the dispersion of the air path, gives the relative spectral phase:

$$\phi(\nu) = 4\pi L / \lambda_c + (4\pi L / v_{group})(\nu - \nu_c) \quad (1)$$

where ν_c is the carrier frequency, v_{group} is the group velocity at the carrier frequency, and λ_c is the carrier wavelength, calculated at measured atmospheric conditions.

A linear fit $\phi = \phi_0 + b(\nu - \nu_c)$ gives the time-of-flight measurement through $L_{tof} = b(v_{group}/4\pi)$ and the high-precision interferometric distance measurement through $L_{int} = (\phi_0 + 2\pi m)(\lambda_c/4\pi)$. The $2\pi m$ ambiguity reflects the inherent $\lambda_c/2$ range ambiguity in any interferometric measurement. From equation (1), the system 20 is essentially an MWI with many simultaneously transmitted wavelengths (equal to the number of transmitted comb lines)

This approach is robust to systematic shifts for several reasons. First, the time gating eliminates shifts due to spurious reflections outside of the +30 ps (+4.5 mm) range window, which can be seen, for example, in FIG. 2C at 8.5 ns and shortly after the reference and target returns. In standard MWI, these spurious reflections are a significant systematic error, because the measured range is effectively a weighted average of all returns. (Either polarization multiplexing or physically separate beam paths are required for the target and reference in MWI.) Second, there are no assumptions about the particular pulse shape. Third, effects due to dispersion are apparent as deviations in the spectral phase, $\phi(\nu)$, from a linear slope. Fourth, strong self-consistency checks are possible through comparisons of the time-of-flight measurements centred at different carrier frequencies (by tuning the optical bandpass filter) and of the averaged time-of-flight and interferometric range measurements.

Rapid update rates are important in applications such as formation flying or large-scale manufacturing, where the range information is used within a feedback system to orient the components. In one example, it takes 200 λ s to scan a 1.5-m ambiguity range, and therefore it is possible to capture the motion of a moving target, as is shown in FIG. 4, where sequential 200 μ s traces are placed side by side to map the position of the moving target.

A Hilbert transform is used to show only the magnitude squared (intensity) of the detected signal. From the time-of-flight, the absolute distance between the reference and target can be determined to within 3 μ m for each trace (FIG. 5). Note the dropout of information at ~1.6 s due to misalignment of the moving cart; for a standard c.w. interferometer with a range ambiguity of one wavelength, such a dropout would ruin the measurement, but here the system 20 readily reacquires the absolute range. A target that moves out of the ambiguity range is readily tracked with an unwrapping algorithm that should allow one to track an object moving as fast as 3,700 m s⁻¹.

High-Precision Absolute Distance Measurements Using Both Time-of-Flight and Interferometric Range Measurements.

For relatively slow moving targets below approximately 10 μ m/s, the time-of-flight measurement can be averaged down sufficiently to hand over to the more precise interferometric range measurement, as described after equation (1). A series of experiments to demonstrate this capability by comparing the results to 'truth' data supplied by a standard fringe counting interferometric distance meter are described below.

In the first experiment, the reference-to-target distance at discrete steps over a ~1 m track was recorded. In satellite-to-satellite positioning the measurement could be at a range of a kilometer or longer. Two additional experiments with a fibre

spool to extend the measurement range were also performed. Unlike space, optical fibre is nonlinear, dispersive, birefringent and backscatters the incident light. To counter nonlinearities, the outgoing signal pulse was chirped (necessary in any case for chirped-pulse amplification). To counter the dispersive effects, ~700 m of higher dispersion fibre was added at the output of the LO laser but within the phase-locked loop, so that the effective delay on the two comb outputs remained, but their relative dispersion was approximately equal at the detector. The most important remaining effect was Rayleigh backscatter from the outgoing pulse, because it cannot be time-gated from the signal. In one experiment, this was avoided by using two identical ~1.14-km spools—one for the outgoing and one for the return signal—and measured the reference-to-target distance at the end of the 1.14-km fibre delay. These data mimic those needed to make remote measurements of the pointing of a satellite, or the angle of a machined surface, through trilateration.

With reference to FIG. 6, no difference in system performance was found for the reference-to-target distance measurements with and without the 1.14-km fibre delay. For test data and the truth data, the range was calculated for identical atmospheric conditions (air temperature, pressure and humidity) so that they shared a common 1×10^{-7} uncertainty from variations in the atmospheric conditions. An independently measured 1 nm drift in the air path (due to temperature change) was subtracted out of the data. At 60 ms averaging, the statistical error on the time-of-flight measurement was below 200 nm, with a systematic error evaluated at less than 100 nm. Because this uncertainty is below $\pm \lambda/2$, the distance measurement could be handed over to the interferometric range measurement, which had only a 20-30 nm scatter versus the truth data, consistent with the dominant uncertainty calculated from the estimated $\sim 0.1^\circ \text{C}$. temperature variations between the two air paths.

FIG. 5 shows the precision of both the time-of-flight and interferometric measurements versus averaging time. Both with and without the fibre delay, the precision of the time-of-flight distance is $\sigma_{tof} = 3 \mu\text{m} (T_{update}/T)^{1/2}$, where T is the averaging period. This scatter is about twice that expected from the measured white noise on the signal and results from ~20 fs residual timing jitter between the combs 24, 28. The precision of the interferometric distance is roughly $\sigma_{int} = 100 \text{ nm} (T_{update}/T)^{1/2}$, reaching 5 nm at 60 ms and continuing to drop below 3 nm at 0.5 s. It is limited by the residual carrier phase jitter between the pulse trains from the combs 24, 28. For both range measurements, tighter phase-locking or post-correction of the data from monitoring of the error signals will improve the precision.

In another experiment, the dual fibre spools were replaced with a single, bidirectional spool and measured the delay between a reference reflection before the 1.14-km fibre spool and the target. The relative uncertainty, also shown in FIG. 5, is limited by Rayleigh backscattering, which gave rise to a 'flicker' noise floor of ~300 nm, which was too large to permit a confident 'handover' to the interferometric range measurement. At longer times, the uncertainty increases due to actual $\sim 1 \mu\text{m s}^{-1}$ fibre length changes from temperature effects. The comb repetition rate of $T_r^{-1} \approx 100 \text{ MHz}$ sets the ambiguity range of our system to $R_A = T_r v_{group}/2 = 1.5 \text{ m}$, adequate for most practical situations, but lower than the fibre delay, equivalent to 1.8 km of air. To remove this ambiguity, the roles of the signal and LO lasers were switched for use with the Vernier effect. Because of the difference in repetition rates, a distance measured normally and a second distance measured with the lasers switched differs by $m\Delta R_A$, where $\Delta R_A = \Delta T_r v_{group}/2$ is the difference in ambiguity ranges, and m

is an integer giving the number of ambiguity ranges by which the true distance exceeds R_A . Therefore, with the two measurements, m can be found and any ambiguity resolved up to the now larger ambiguity range of $v_{group}/(2\Delta f_r) = 30 \text{ km}$, where $\Delta f_r = (T_{update})^{-1} \approx 5 \text{ kHz}$ is the difference in comb repetition rates. Applying this technique to the fibre spool, a fibre length of 1,139.2 m was measured in good agreement with a standard optical time-domain reflectometry measurement (using the same group index) of $1,138.4 \pm 1 \text{ m}$.

The limit to the fractional accuracy in the time-of-flight and interferometric range measurements is ultimately the fractional accuracy in the rf timebase and optical frequency, respectively. Here, an rf time base is provided by the hydrogen maser that can support better than 1×10^{-13} fractional ranging resolution, that is, 3 nm in 30 km or below the systematic uncertainty. The fractional accuracy of the carrier frequency will depend on the underlying c.w. reference lasers, which can be stabilized to the calibrated reference cavity, a molecular reference, or a self-referenced frequency comb. (Here, the carrier frequency was stabilized to the previously described reference cavity 34, which had ~30 kHz wander as measured with a self-referenced frequency comb.) The fractional accuracy provided by the optical reference cavity 34 or molecular reference is sufficient for nanometer-scale measurements at short ranges, or for differential range measurements at long ranges. Nanometer scale absolute ranging at long distances requires a fully self-referenced comb, with the associated technical complexity. For extreme precision at very long ranges, the entire system may be based on an optical clock, to provide an increase of greater than 1×10^4 in accuracy, with effectively unlimited operation distances.

The data above illustrates the ability of this system 20 to measure the range between multiple reflections over a large range window (ambiguity range) at a short update rate and at long ranges. Equally important are the built-in checks on hidden systematics by comparing time-of-flight and interferometric range measurements. Although the detection and processing is straightforward, the technical challenge lies in the dual, coherent frequency comb sources, particularly for satellite applications. The comb source requirements here are simpler in two regards than those for the fully self-referenced octave-spanning combs needed to support current optical clocks. First, although the combs are phase-locked to hertz-level linewidths here, the absolute linewidth actually need only be below Δf_r to cover the full 1.5-m ambiguity range. (Narrower linewidths are required only for extremely long-range operation beyond $v_{group}/\Delta f_r = 30 \text{ km}$.) Second, the comb output need only span 10-20 nm optical bandwidth rather than a full octave. The basic technique is not limited to fibre-based mode-locked lasers, and other passively mode-locked lasers may alternatively provide a more robust and compact system. Third, if only the time-of-flight data is needed, the optical carrier of the combs can be allowed to drift.

As stated earlier, for either large-scale manufacturing or formation flying, it is often the angle or pointing of the target that is most critical. The pitch and yaw of the object can be determined through trilateration, where the angle is calculated from the distance to several reference points 30 spanning the target (FIG. 7).

Multiple satellites flying in a precision formation can effectively act as a single distributed instrument and provide entirely new capabilities for space-based sciences. Formations would enable higher-resolution searches for extraterrestrial planets by providing a large synthetic aperture, enable direct imaging of a black hole by supporting an X-ray telescope distributed across satellites, or enable tests of general

relativity through accurate measurements of satellite spacing in a gravitational field. The formation acts as a single instrument only if the relative spacing and pointing of the satellites is tightly maintained, which is made possible by comparing distance measurements between multiple reference points on the satellites and feeding back to the satellite position and pointing. This approach is well suited to carrying out this measurement, because the time gating of multiple reflections allows one to collect and process data for all trilateration reference points on the same photodiode and digitizer, saving both space and power. The wide ambiguity range and high update rate can support rapid reconfiguration of the target objects.

For active stabilization of absolute distances, or pointing angles, the time-of-flight measurement may beneficially be reduced to below $\lambda_c/4$ in a single scan so that interferometric precision could be reached in a single update time, T_{update} . The time-of-flight measurement uncertainty improves through increase in the signal pulse bandwidth up to the available ~ 4 THz source bandwidth (assuming a corresponding reduction in the residual comb jitter). However, the pulse bandwidth BW is constrained by the relationship $BW < T_{update}/2T_r^2$ (to achieve adequate sampling), so a larger bandwidth implies a lower update rate. This constraint can be removed by coherently combining multiple detection channels at different wavelength bands across the source. With this approach, nanometer-level precision is possible at less than a millisecond update rate.

The linear optical sampling arrangement is relatively straightforward, but contains a few technical subtleties. In the linear optical sampling picture, signal and LO pulses arrive at the detector at a rate of ~ 100 MHz and with varying delays between them. The signal and LO pulse overlap is integrated over the detector response time ($\sim 1/100$ MHz) to yield a single voltage sample, which is then synchronously digitized at the LO repetition rate. Sampling is performed, for example, with a 110 MHz balanced detector that allows amplitude noise suppression in the relatively high power LO. To eliminate detector ringing effects, the signal is low-pass filtered at 50 MHz (which also eliminates any nonsensical signal above the Nyquist frequency set by the LO sampling rate). Once the signal and LO are combined on the same detector, the measurement is insensitive to electronic phase shifts as long as the system remains linear. Saturation of detectors and amplifiers is avoided.

For time-of-flight distance measurements with resolution smaller than 200 nm, a few pitfalls may arise. In virtual time, the LO samples the signal laser $T_r/\Delta T_r = 19,260$ times per signal repetition period, T_r . This gives an effective sample step size of 78 μm . For ease in processing one might be tempted to set up the system so that after every 19,260 samples the signal and the LO have the same phase, however, systematics related to the pulse shape limit ability to measure the centre of peaks below a precision of 300 nm with a 78- μm step size. Instead, the phase locks are arranged so that the pulses overlap every 19,260 \times 51 samples. Averaging sequential frames then effectively subdivision of the 78- μm step size by 51, which is sufficient to achieve sub-100-nm resolution. It was also found that the use of interleaved ADCs (analog-to-digital converters) in the digitizer may lead to scatter in the data.

For the interferometric measurement, the phase factor ψ , is defined herein as $\psi = \pi + \phi_{Gouy}$. The factor π arises from the reflection from the air-glass interface at the target. The reference pulse occurs from a glass-air interface and suffers no such phase shift. The Gouy phase shift is dependent on the

actual distance and we use a separate measurement of the beam Rayleigh range to calculate this phase shift from the time-of-flight distance data.

In an experiment, the processing is not performed in real time. Rather, the raw data is collected and post-processed. The processing is done on a 'scan-by-scan' basis. The 19,260 data points from each ~ 200 - μs -long scan are first high-pass-filtered at 5 MHz and then searched for the three largest peaks, the first two of which are from the front and back surfaces of the reference flat and the last one from the target. Two copies of the data are then generated: one with a 60-ps time window (the exact width is not important) around the appropriate reference reflection and one with a 60-ps time window about the target reflection. The 19,260 data points are truncated to a length that has only low numbers in its prime factorization (for example, 19,200) and fast Fourier transformed (FFT). The spectral phase is extracted from the FFT and fit across ~ 0.4 THz of bandwidth to equation (1) and further processed as described in the text. The most computationally intensive operation is the FFT, and real-time processing of the data should be possible with modern field programmable gate arrays.

As with any ranging LIDAR system, the achievable distance resolution, ΔR , is

$$\Delta R \sim \frac{c}{2 \times BW \times SNR} \quad (0.1)$$

where c is the speed of light, BW is the transmitted optical bandwidth, and SNR is the signal-to-noise ratio. For absolute ranging, there are also contributing systematic errors. This equation applies to time-of-flight pulse measurement, where the BW is inversely proportional to pulse width, and reflects how well one can estimate the pulse center. It also applies to the interferometric measurement provided the BW is replaced by the carrier optical frequency, ν , and the SNR by the fractional pulse-to-pulse phase noise $\delta\phi/(2\pi)$. In our system, a "single" measurement takes an update time $t_{update} = 1/\Delta f_r$, which is the time for the LO comb to fully sweep across the signal comb.

In this case, there is a tradeoff between bandwidth and update rate,

$$BW < t_{update} \frac{f_r^2}{4} = t_{update} \frac{c^2}{16R_A^2} \quad (0.2)$$

where f_r is the repetition rate, and $R_A = c/(2f_r)$ is the ambiguity range. This equation can be derived in the time domain by requiring the LO pulse never "misses" the signal pulse or in the frequency domain by requiring there is never any aliasing of multiple optical heterodyne beats onto a single rf beat (See FIG. 1). Assuming the SNR is close to shot-noise limited at $SNR \sim \eta \sqrt{P_{return}}/f_r$, where P_{return} is the total detected return power and η is the detector responsivity, then Eq. (0.1) and (0.2) yield a "figure-of-merit",

$$\Delta R \times t_{update} \sim \frac{2c}{\eta \sqrt{P_{return}}} f_r^{-3/2}, \quad (0.3)$$

i.e. the product of resolution and update rate, which we would like to be as small as possible. The overall performance

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improves dramatically with increasing repetition rate, at the cost of a lower ambiguity range.

The source bandwidth typically far exceeds the detection bandwidth, BW. Therefore, multiple detectors can be utilized to take advantage of the full available bandwidth as a linear array of photodetectors may be used after the grating-based (or other) spectral filter to detect multiple spectral bands at once. In that case, Eq. (0.1) has an extra factor of $1/M$, where M is the number of detectors (or spectral channels, each with bandwidth BW that satisfies Eq. (0.2)) and the SNR per channel is modified to $SNR \sim \eta \sqrt{P_{return}} / M f_r$ to yield,

$$\Delta R \times t_{update} \sim \frac{2c}{\eta \sqrt{P_{return}} \sqrt{M}} f_r^{-3/2} \quad (0.4)$$

Once ΔR has dropped below a quarter wavelength, then the interferometric data provides much more precise ranging. In practice, the resolution of (0.4) is reached only if the systematics are sufficiently low, which requires either active locking of the two combs, or, alternatively, accurate monitoring of their relative coherence and a software correction.

The nm-scale ranging that is possible with the combs is excessive for many terrestrial applications. However, the system is flexible in terms of trading off precision for speed, as indicated in Eq. (0.1) to (0.4). While these scaling laws relate to the measurement of a single range, the system can also be configured so that the M channels each measure the range to a different point on the surface. In that case, the same basic scaling of (0.3) applies to each channel independently, and P_{return} is the power per measurement point. Each channel retains the ability to “handover” the time-of-flight range measurement to an interferometric measurement with nm-level precision. This ability to measure the absolute range at nm-level precision and with reasonable (cm to meter) ambiguity range is a feature not easily found in current 3D laser radar systems. Such a system would provide much higher potential range resolution than current Geiger-mode APD based 3D imaging systems although it would require significantly more return power and would therefore only be suitable at shorter distances and when such increased range resolution is warranted.

The disclosed non-limiting embodiment provides ranging application of a dual-comb coherent LIDAR system. However, the coherent LIDAR can alternatively operate as a high-resolution vibrometer as well. In fact, the broad comb bandwidth lends itself to spectral averaging to reduce the limiting effects of speckle on high-resolution vibrometry signals. The use of dual combs relaxes the significant range restriction of the previous comb-based vibrometer, and one could conceivably configure the system to loosely lock onto a target for high-resolution vibrometry data, although such a system would require more complete control of the combs than demonstrated here. Another potential application of coherent frequency combs is in the area of synthetic aperture LIDAR particularly given the analogy of this source to a coherent pulsed RADAR source.

The highly coherent frequency combs have the potential to provide very high resolution coherent LIDAR systems for ranging and vibrometry, as well as a host of other applications including synthetic aperture LIDAR and multispectral LIDAR. To fully exploit their broad bandwidth does require an increase in transmit power over a single wavelength system but the dual-comb system here is completely compatible with chirped pulse amplification. Unlike “metrology-grade” frequency combs, the sources here are considerably simpler

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since they do not require a full octave of output bandwidth. Nevertheless, more work is required to either further improve the robustness of fiber-based combs or alternatively, to base these systems on other types of compact and robust modulated lasers.

The frequency comb based LIDAR disclosed herein offers a host of powerful features—precision, stability, speed, large-range ambiguity, low light level operation, multiplexing capabilities, flexibility and spurious reflection immunity—that, as a whole, are unavailable in existing approaches.

It should be understood that relative positional terms such as “forward,” “aft,” “upper,” “lower,” “above,” “below,” and the like are with reference to the normal operational attitude of the vehicle and should not be considered otherwise limiting.

It should be understood that like reference numerals identify corresponding or similar elements throughout the several drawings. It should also be understood that although a particular component arrangement is disclosed in the illustrated embodiment, other arrangements will benefit herefrom.

Although particular step sequences are shown, described, and claimed, it should be understood that steps may be performed in any order, separated or combined unless otherwise indicated and will still benefit from the present disclosure.

The foregoing description is exemplary rather than defined by the limitations within. Various non-limiting embodiments are disclosed herein, however, one of ordinary skill in the art would recognize that various modifications and variations in light of the above teachings will fall within the scope of the appended claims. It is therefore to be understood that within the scope of the appended claims, the disclosure may be practiced other than as specifically described. For that reason the appended claims should be studied to determine true scope and content.

What is claimed is:

1. A method of comb-based coherent Light Detection and Ranging (LIDAR) comprising:

generating a signal comb that transmits a signal pulse train reflected from a reference and a target to generate a target signal from the target and a reference signal from the reference;

generating a local oscillator (LO) comb that transmits a LO pulse train at a slightly different repetition rate than the signal pulse train with a coherent optical carrier;

detecting the reflected signal pulse train through linear optical sampling against the LO pulse train in which consecutive samples of an overlap between a signal pulse from the target signal and the reference signal, and a LO pulse from the LO pulse train yields a high-resolution measurement of the target signal and reference signal;

Fourier transforming the target signal to extract the optical phase vs. frequency of the target signal;

Fourier transforming the reference signal to extract the optical phase vs. frequency of the reference signal;

determining a frequency dependant optical phase difference between the optical phase of the target signal and the optical phase of the reference signal;

fitting the optical phase difference to a straight line;

obtaining a time-of-flight distance measurement from a slope of the straight line and the known group velocity of light at the frequency of the coherent optical carrier; obtaining an interferometric distance measurement between the target and the reference from an intercept of the straight line and the wavelength of the coherent optical carrier.

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2. A method as recited in claim 1, wherein the time-of-flight measurement is repeated and averaged until a desired uncertainty of the time-of-flight distance measurement is achieved.

3. A method as recited in claim 1, wherein the time-of-flight measurement is repeated until an uncertainty of the time-of-flight distance measurement is below half the wavelength of optical carrier.

4. A method as recited in claim 3, wherein the time-of-flight measurement is repeated to remove an ambiguity of the interferometric distance measurement.

5. A method as recited in claim 1, wherein generating the signal comb that transmits the signal pulse train reflected from the reference and the target with a first mode locked laser.

6. A method as recited in claim 5, wherein generating the local oscillator (LO) comb that transmits the LO pulse train with a second mode locked laser.

7. A method as recited in claim 6, wherein the coherent optical carrier is generated by phase locking the signal comb and the LO comb with a common frequency reference that is a continuous wave laser.

8. A Light Detection and Ranging (LIDAR) comprising:
a first mode locked laser which generates a signal comb that transmits a signal pulse train reflected from a reference and a target subject to a distance measurement relative the reference;

a second mode locked laser that generates a local oscillator (LO) comb;

a first continuous wave laser and a second continuous wave laser stabilized to an optical reference cavity, said first continuous wave laser phase locked at a first tooth of the signal comb and the LO comb, said second continuous wave laser phase locked at a second tooth of the signal

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comb and the LO comb different than the first tooth to relate a distance between the reference and the target to a known distance of said optical reference cavity;

a beamsplitter including first and second output ports, said beamsplitter being arranged to receive and combine said signal comb, or said LO comb with said first continuous wave laser and said second continuous wave laser and direct a first output from said first output port to said target; and

wherein said beamsplitter is arranged to direct a second output from said second output port to a 2-channel, optical filter with each channel centered on one of said continuous wave lasers.

9. A LIDAR as recited in claim 8, wherein the signal comb and the LO comb are phase locked with said first continuous wave laser such that their optical carriers are mutually coherent and phase locked with said second continuous wave laser such that their repetition rates have a slight, known difference.

10. A LIDAR as recited in claim 8, further comprising two detectors that are each operable to detect a beat signal between a respective one of said combs and a respective one of said continuous wave lasers.

11. A LIDAR as recited in claim 10, further comprising phaselock electronics arranged to receive said second output from said detectors and generate an error signal for feedback to said frequency combs.

12. A LIDAR as recited in claim 8, wherein said beamsplitter is a polarizing beamsplitter.

13. A LIDAR as recited in claim 12, wherein said beamsplitter is a polarization insensitive beamsplitter of any ratio with regard, respectively, to said first output and said second output.

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