

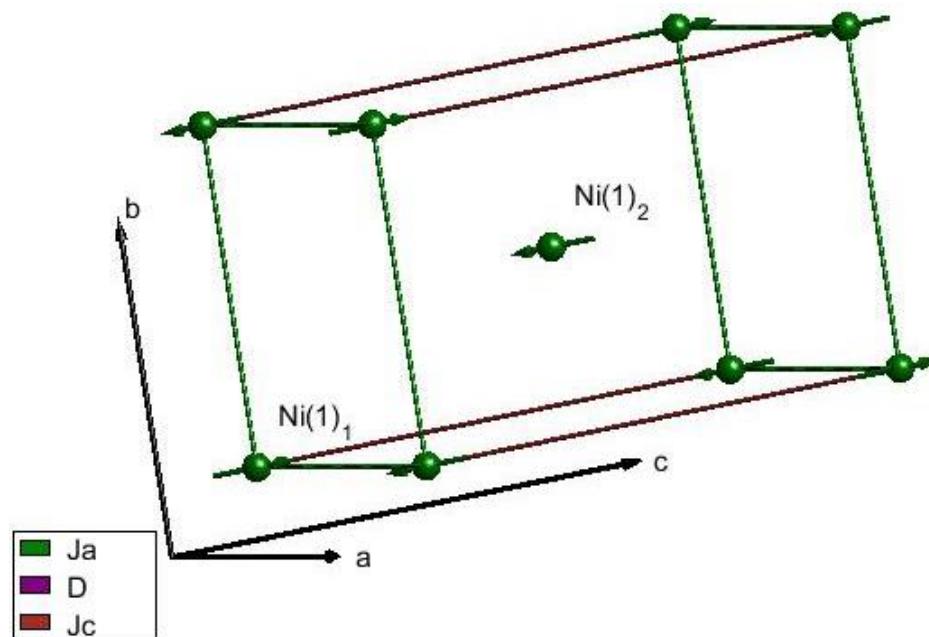
Neutron Scattering studies of the Crystal and Magnetic Structures of Molecular Magnets

Danielle Villa

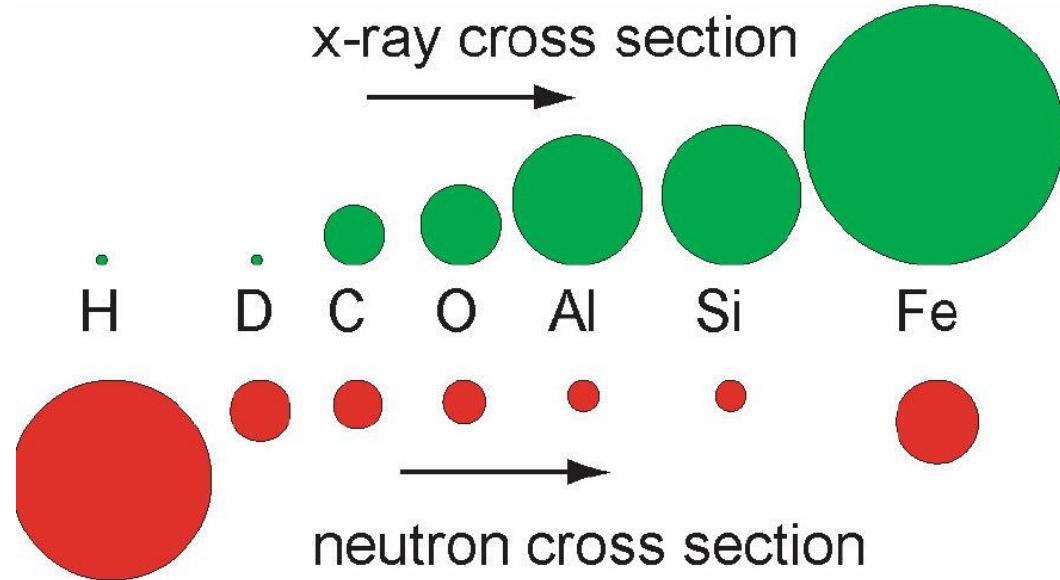
Mentors: Craig Brown and Qingzhen Huang

Molecular Magnets

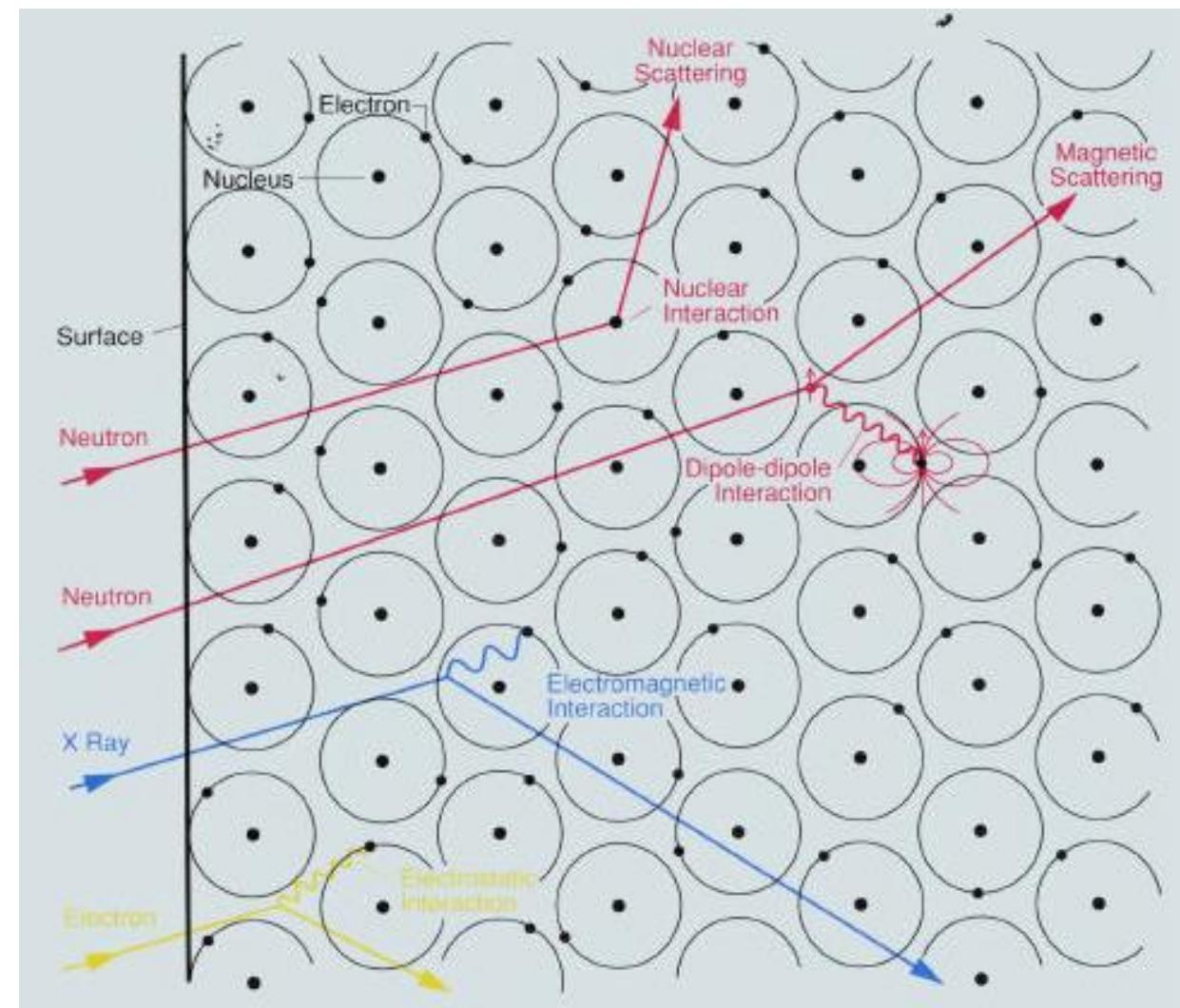
- ▶ Molecular Quantum Magnets
 - ▶ Molecules that exhibit magnetic behavior at low temperatures
 - ▶ Antiferromagnetism 
 - ▶ $\hat{H} = \sum D(\hat{S}_i^{z^2}) + \sum J_{ij}\hat{S}_i \cdot \hat{S}_j$
 - ▶ D-> Single Ion Anisotropy
 - ▶ J-> Magnetic Exchange Interaction
- ▶ Quantum Theory
- ▶ Information Storage
- ▶ Quantum Computing



Neutrons



"Dynamics and Neutron Scattering" by John R.D. Copley



Neutron Scattering "LECTURE 1: Introduction & Neutron Scattering "Theory"" by Roger Pynn

Neutron Scattering

Diffraction



BT-I High Resolution Powder Diffractometer

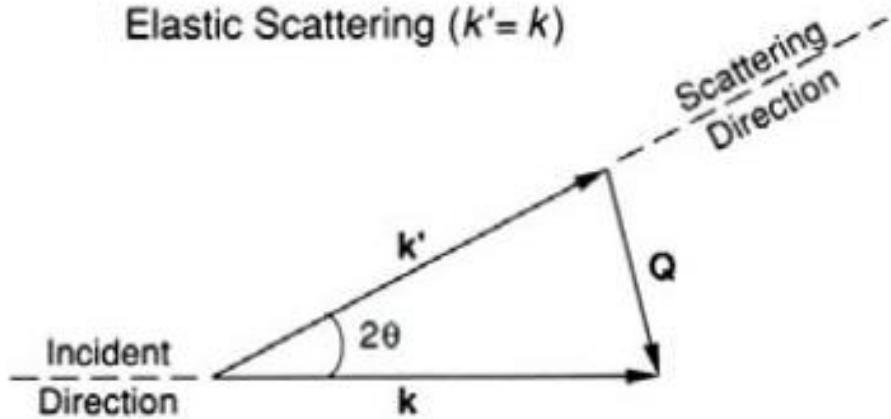
Inelastic



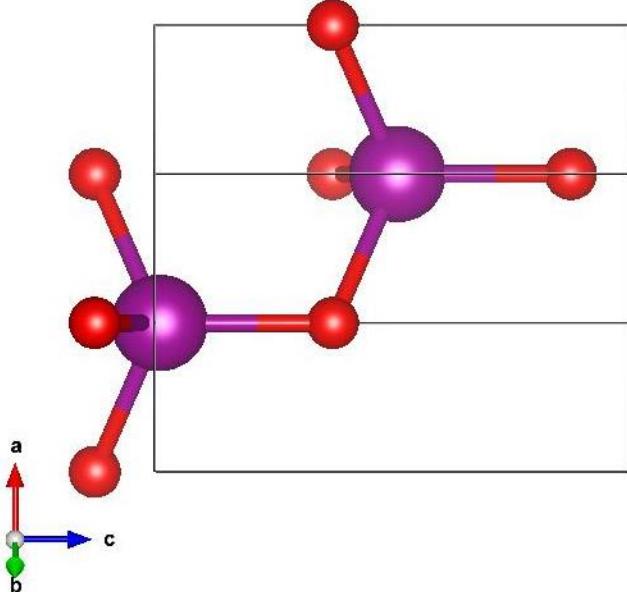
Disk Chopper Time of Flight Spectrometer

Powder Diffraction

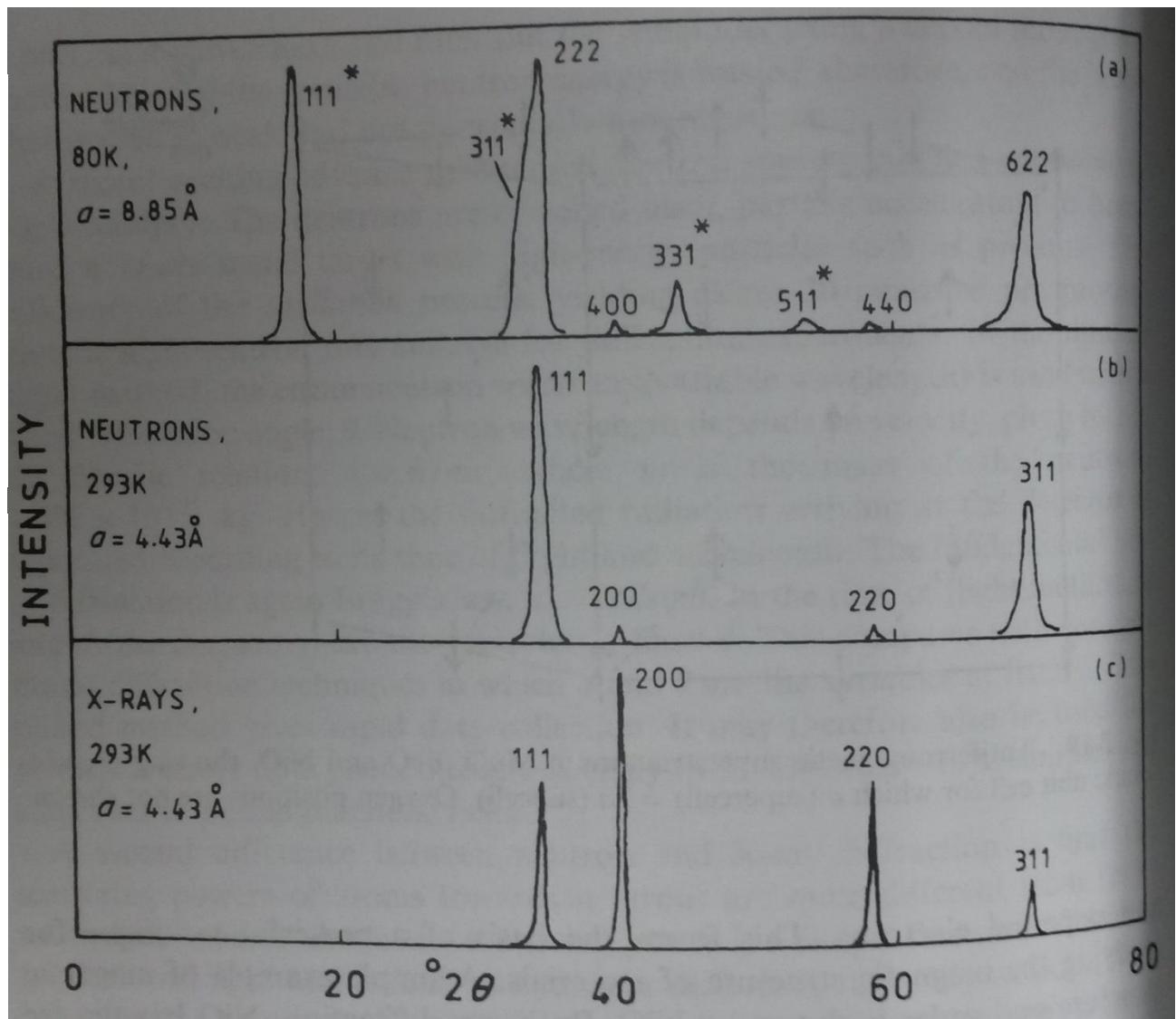
Elastic Scattering ($k' = k$)



https://en.wikipedia.org/wiki/X-ray_crystallography



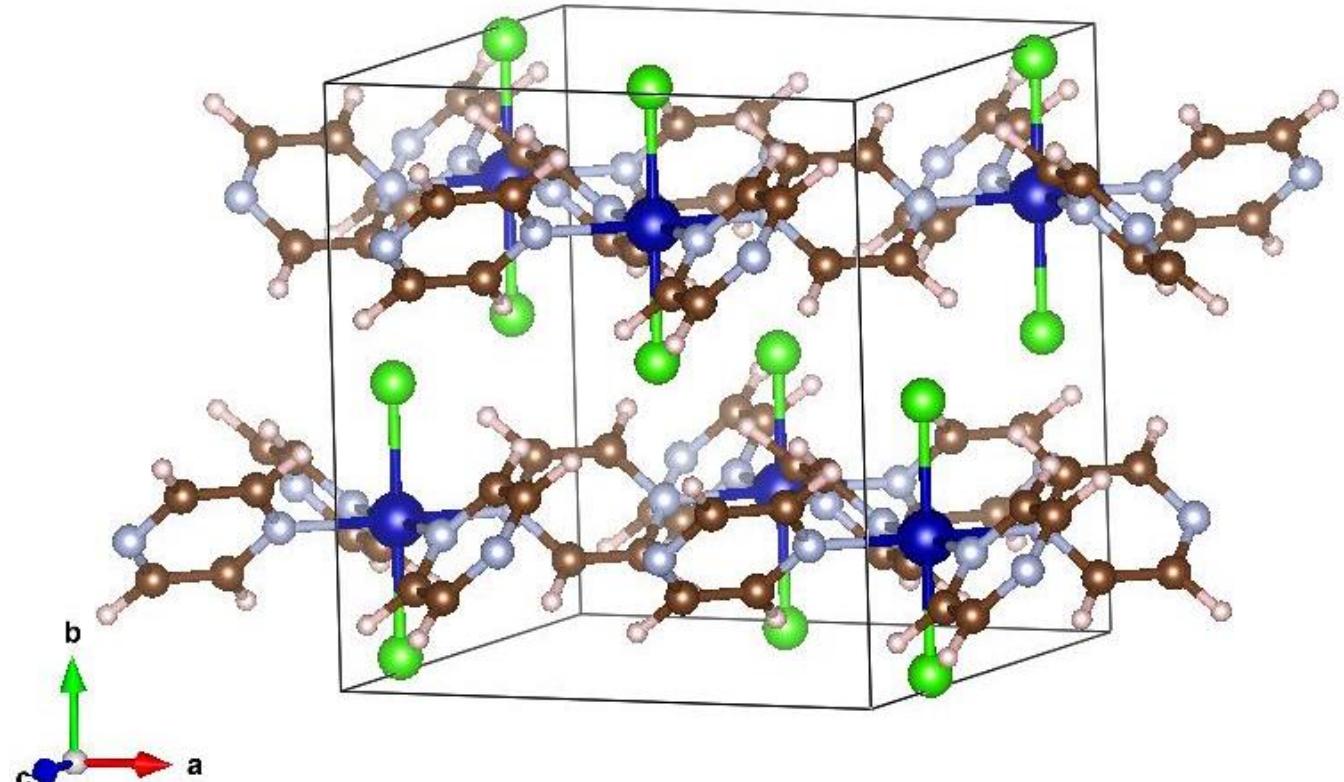
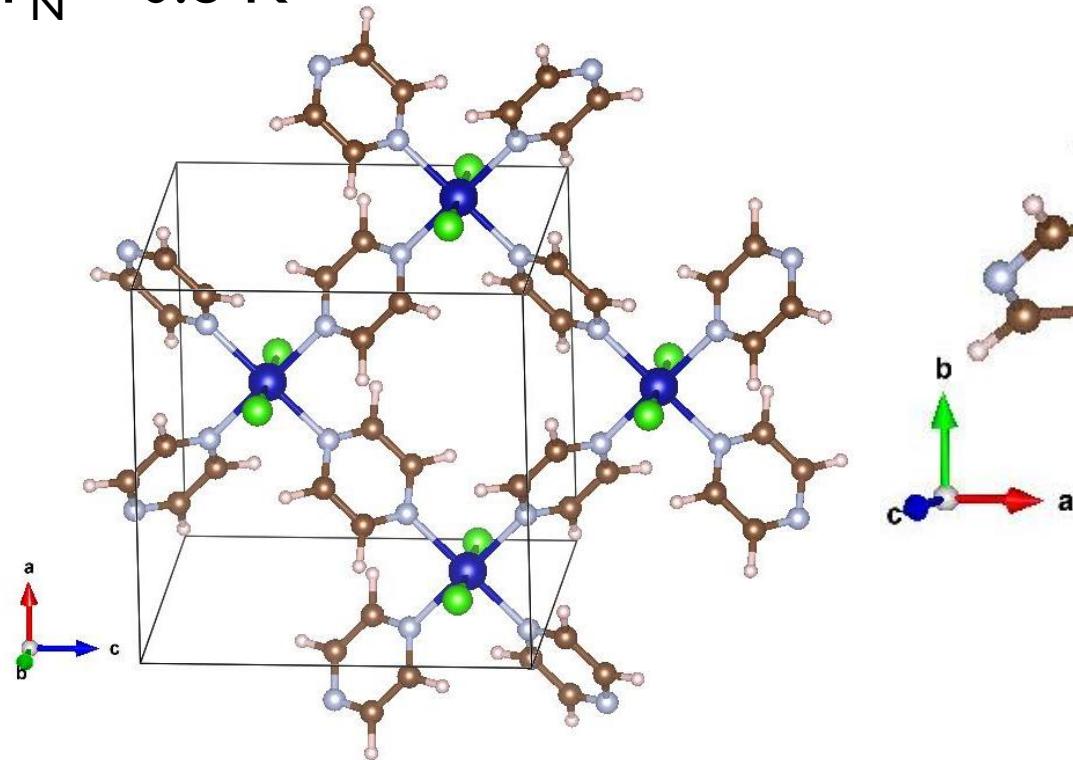
Structure from Ki Min Nam, et al. J. Am. Chem. Soc., 2012, 134 (20), pp 8392–8395

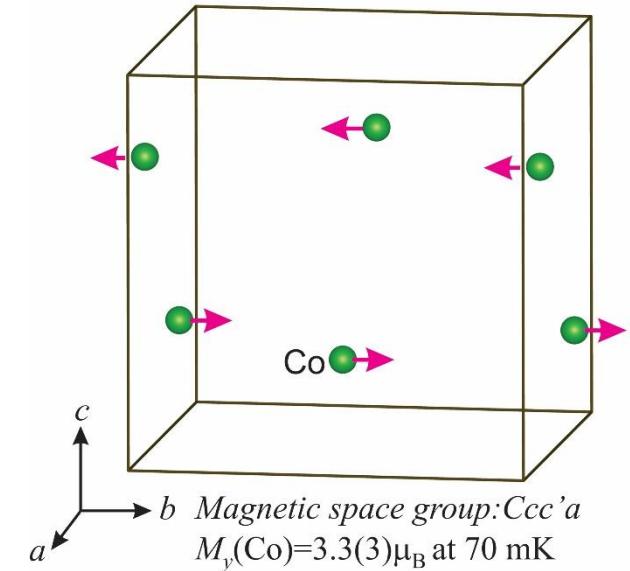
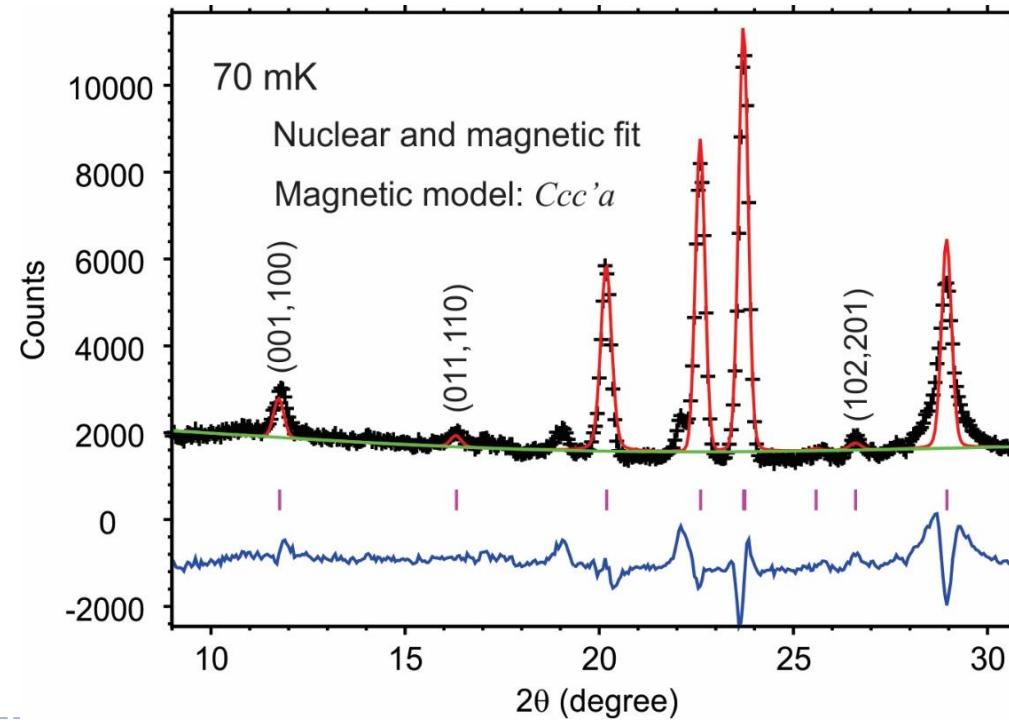
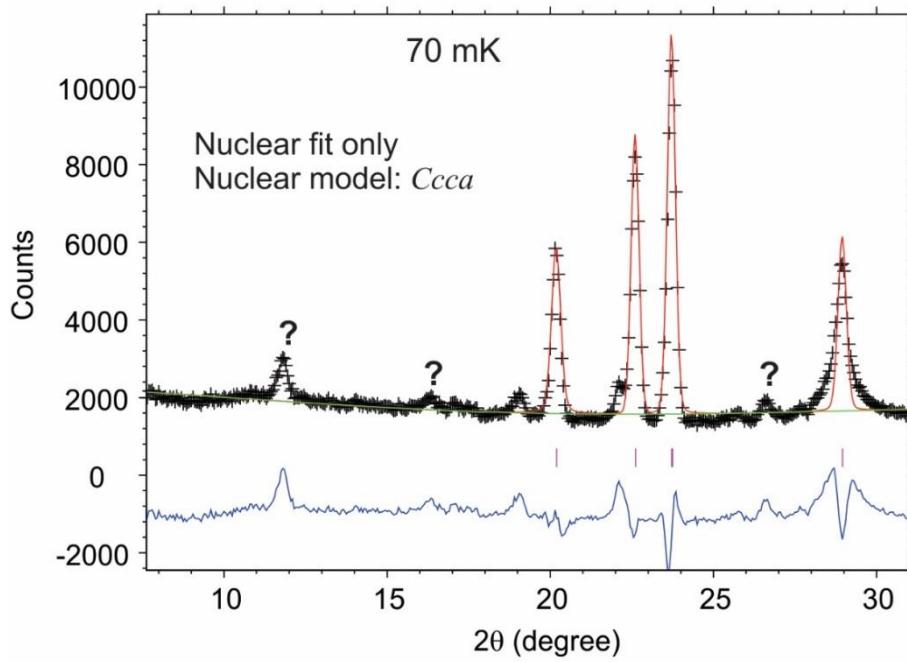
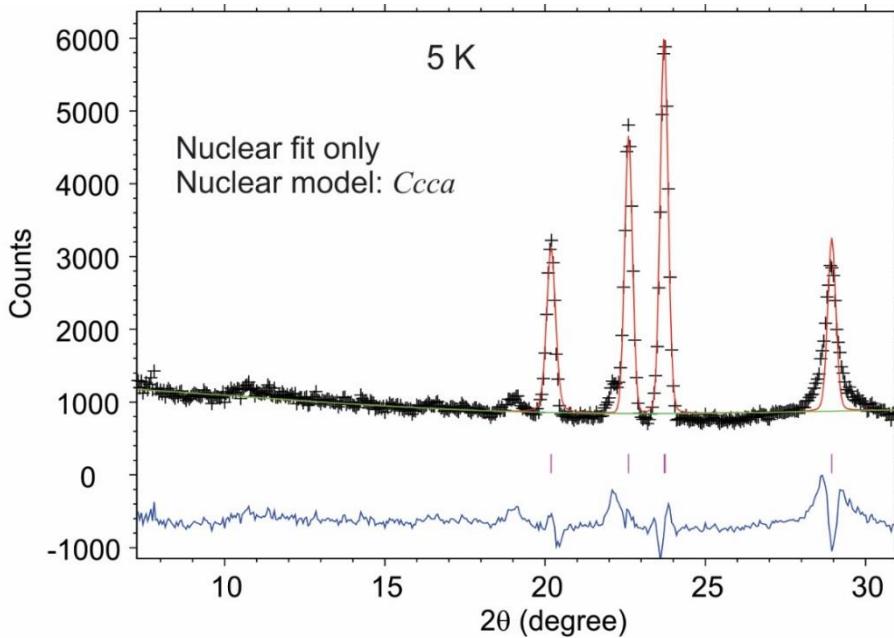


Basic Solid State Chemistry by Anthony R West (pg. 164)

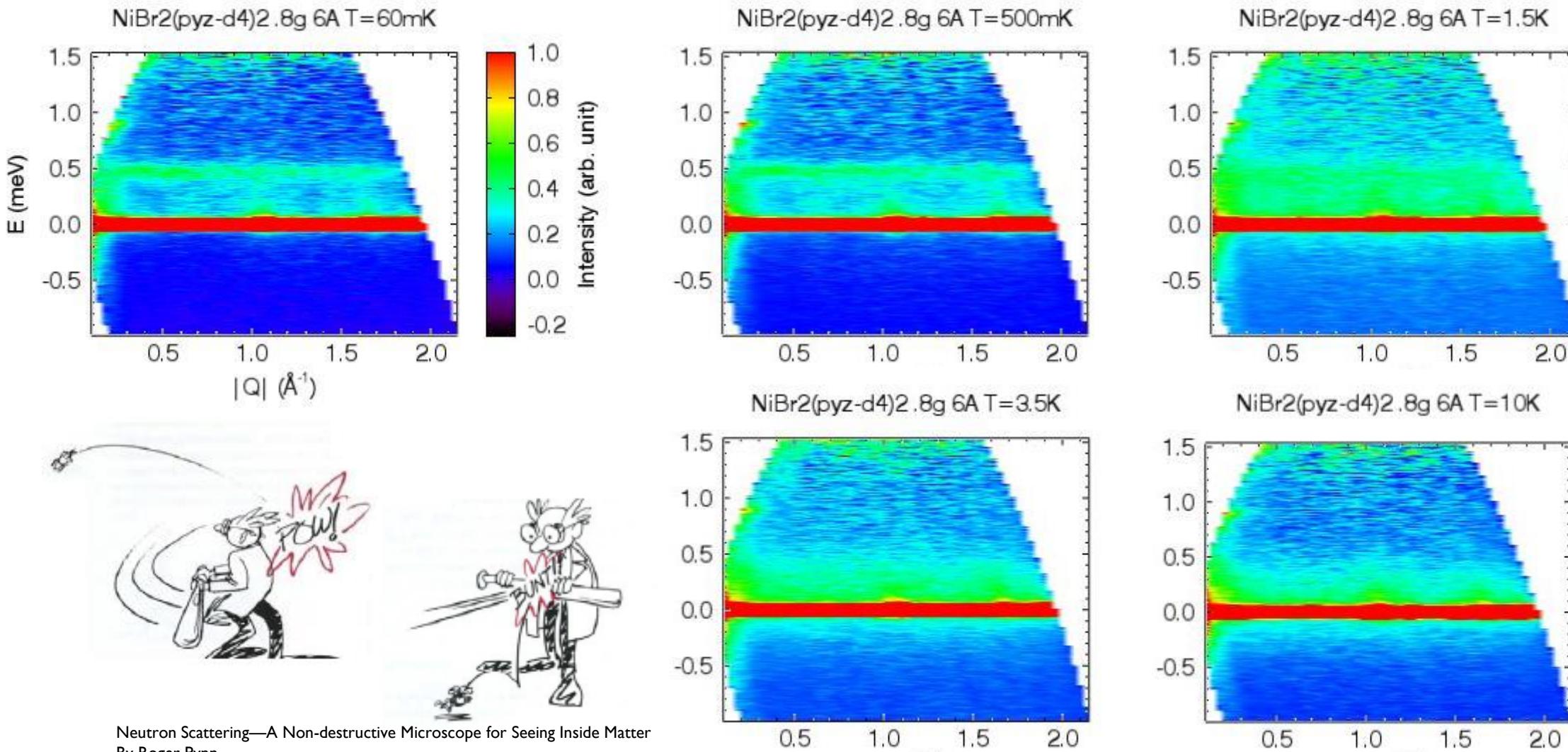
$\text{CoCl}_2\text{pyz}_2$

- ▶ $S = 3/2$
- ▶ Orthorhombic
- ▶ $T_N = 0.8 \text{ K}$





Inelastic Scattering



$\text{NiX}_2(\text{pyz})_2$

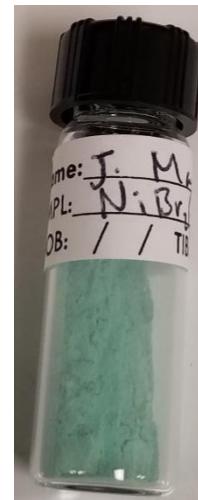
$\text{NiCl}_2(\text{pyz})_2$

- ▶ S=I
- ▶ Tetragonal

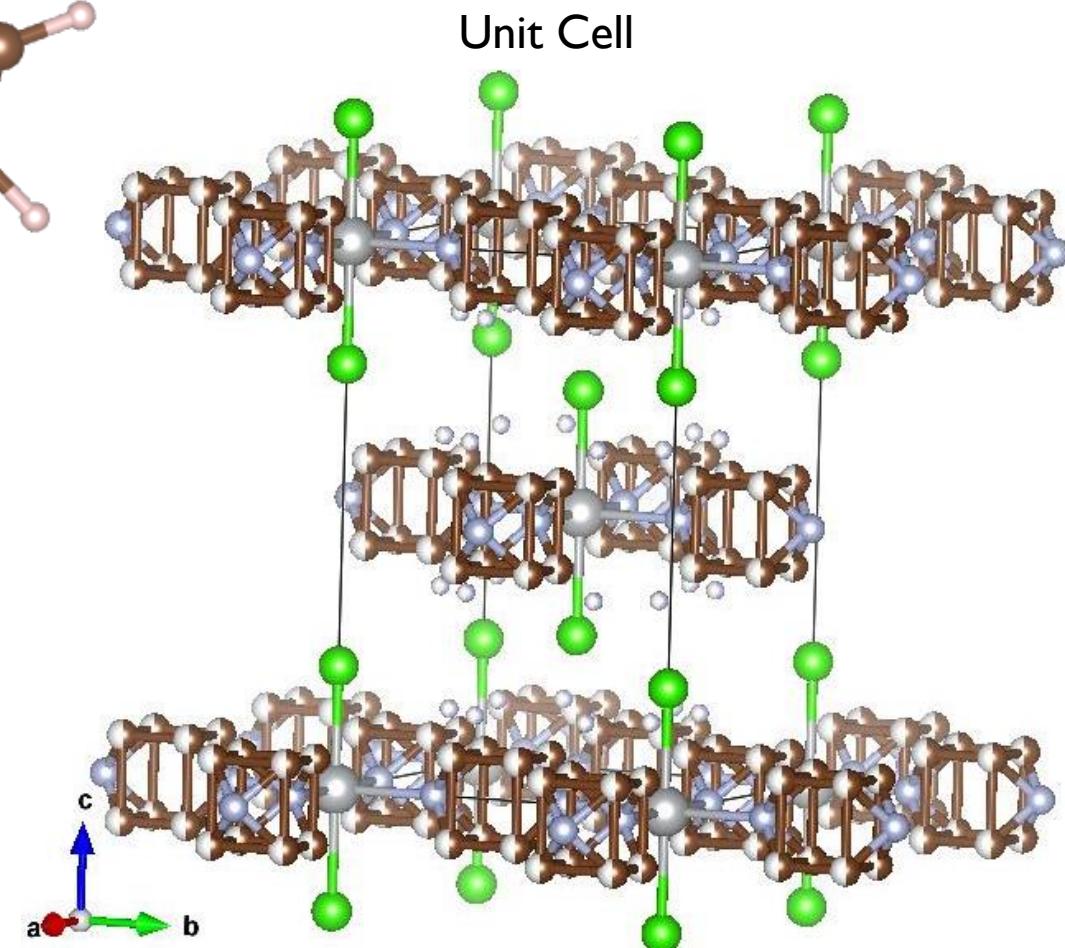
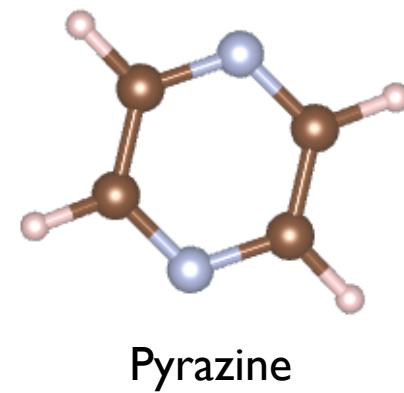
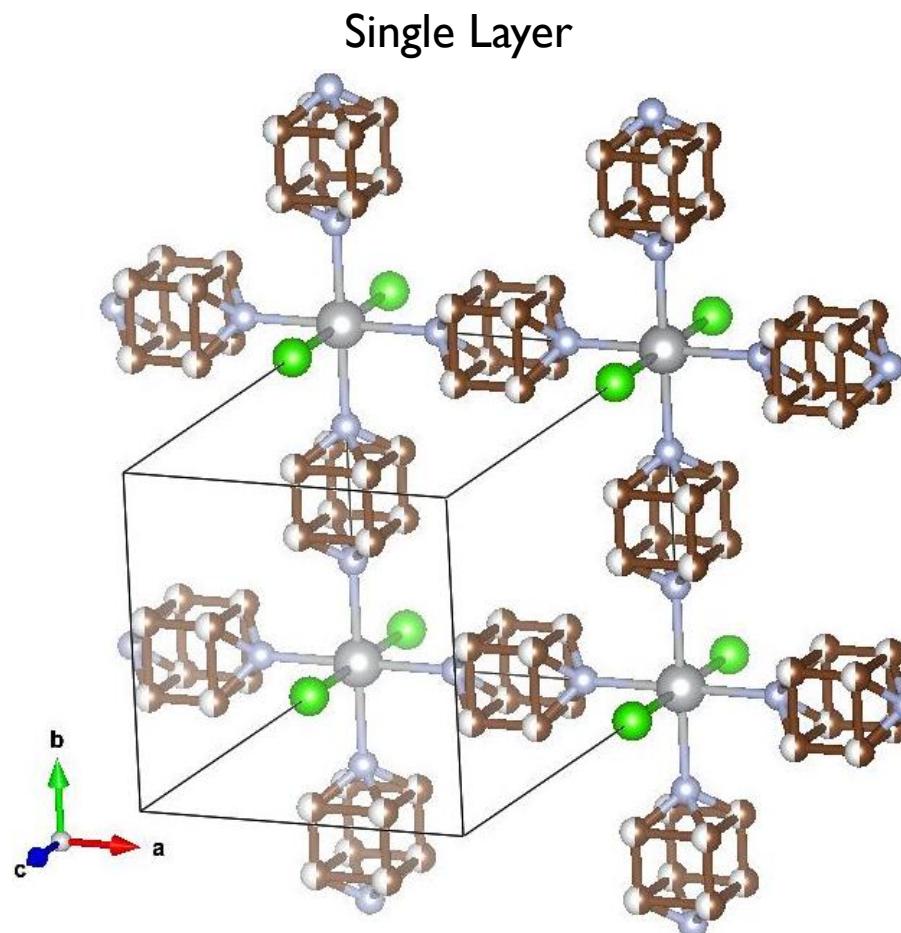


$\text{NiBr}_2(\text{pyz})_2$

- ▶ S=I
- ▶ Tetragonal



$\text{NiX}_2(\text{pyz})_2$



$\text{NiX}_2(\text{pyz})_2$

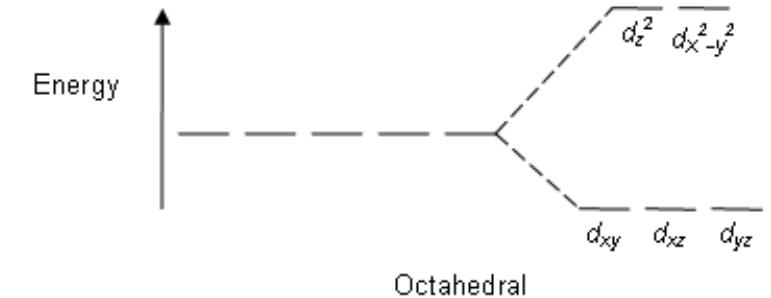
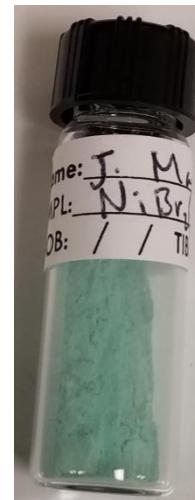
$\text{NiCl}_2(\text{pyz})_2$

- ▶ S=1
- ▶ Tetragonal



$\text{NiBr}_2(\text{pyz})_2$

- ▶ S=1
- ▶ Tetragonal

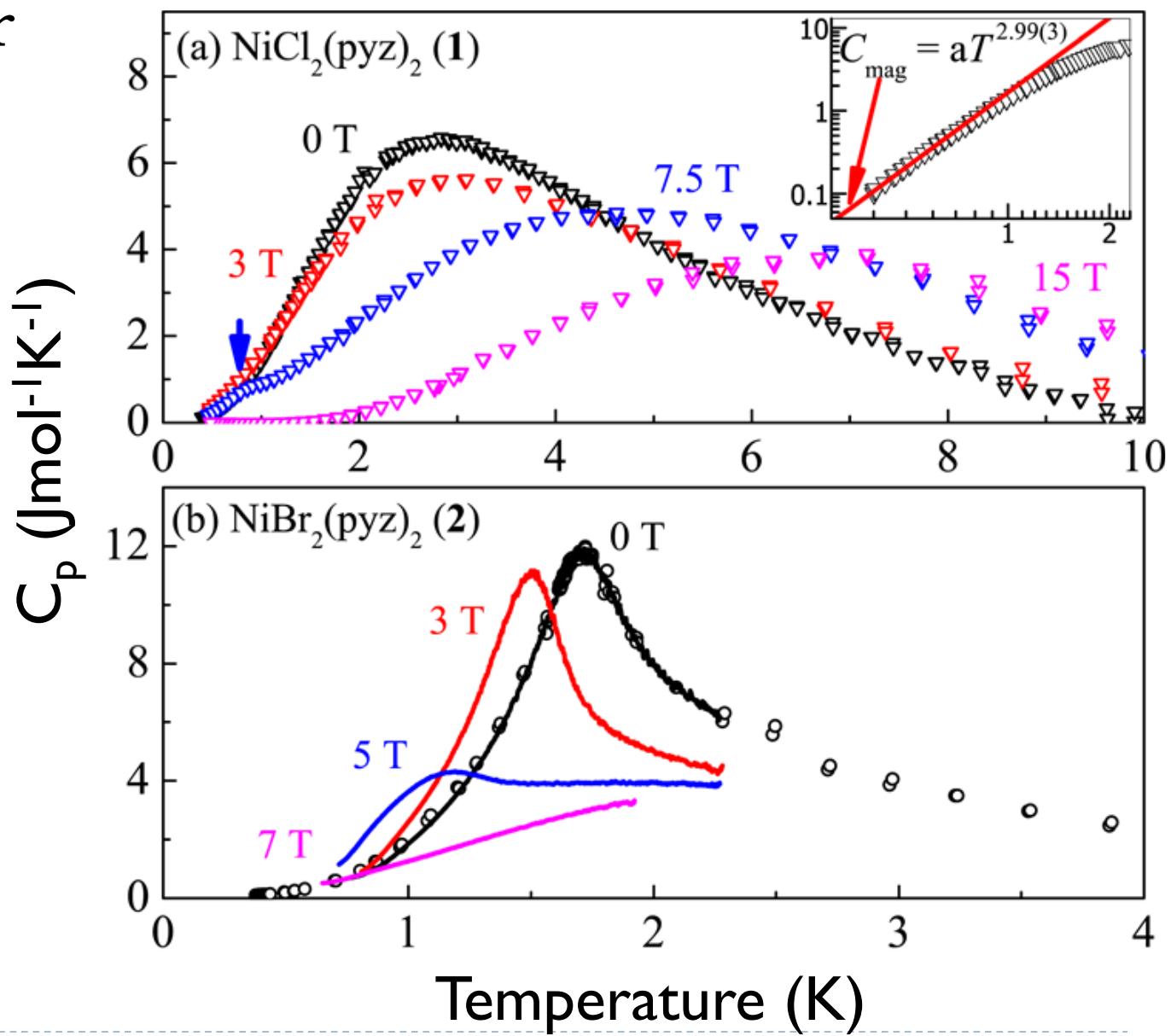
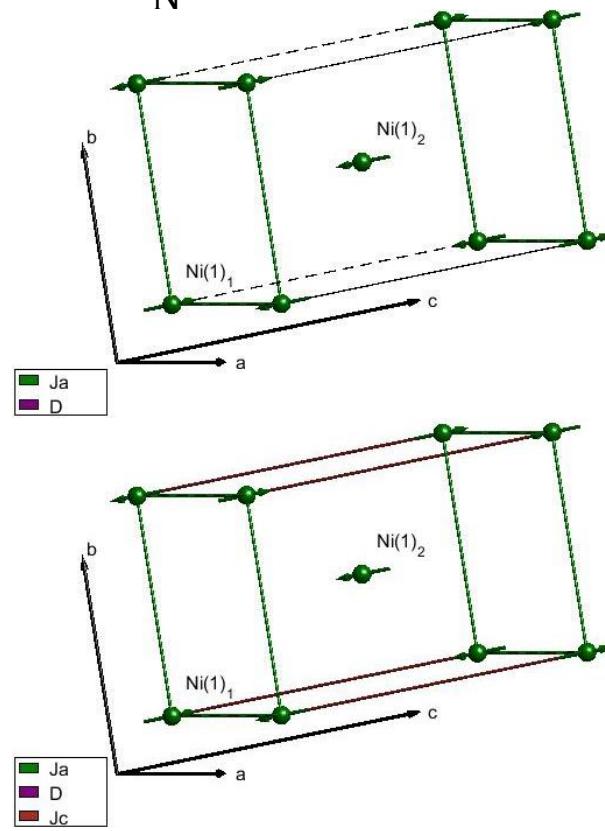


17
F
Cl
Br
I



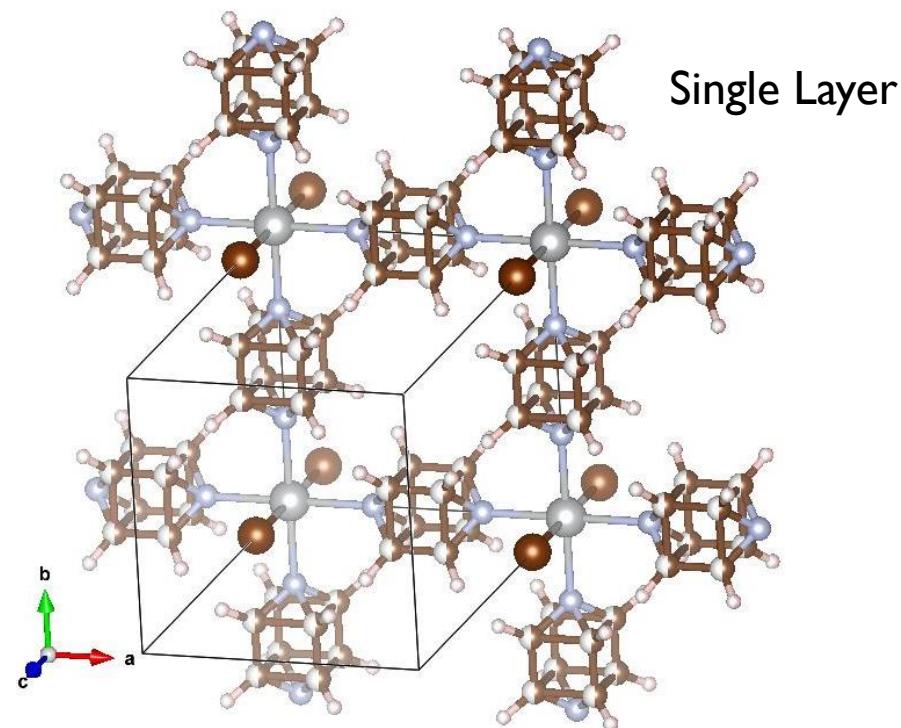
Long Range Magnetic Order

- NiBr_2
 - $T_N = 1.9 \text{ K}$
- NiCl_2
 - $T_N < 0.08 \text{ K}$

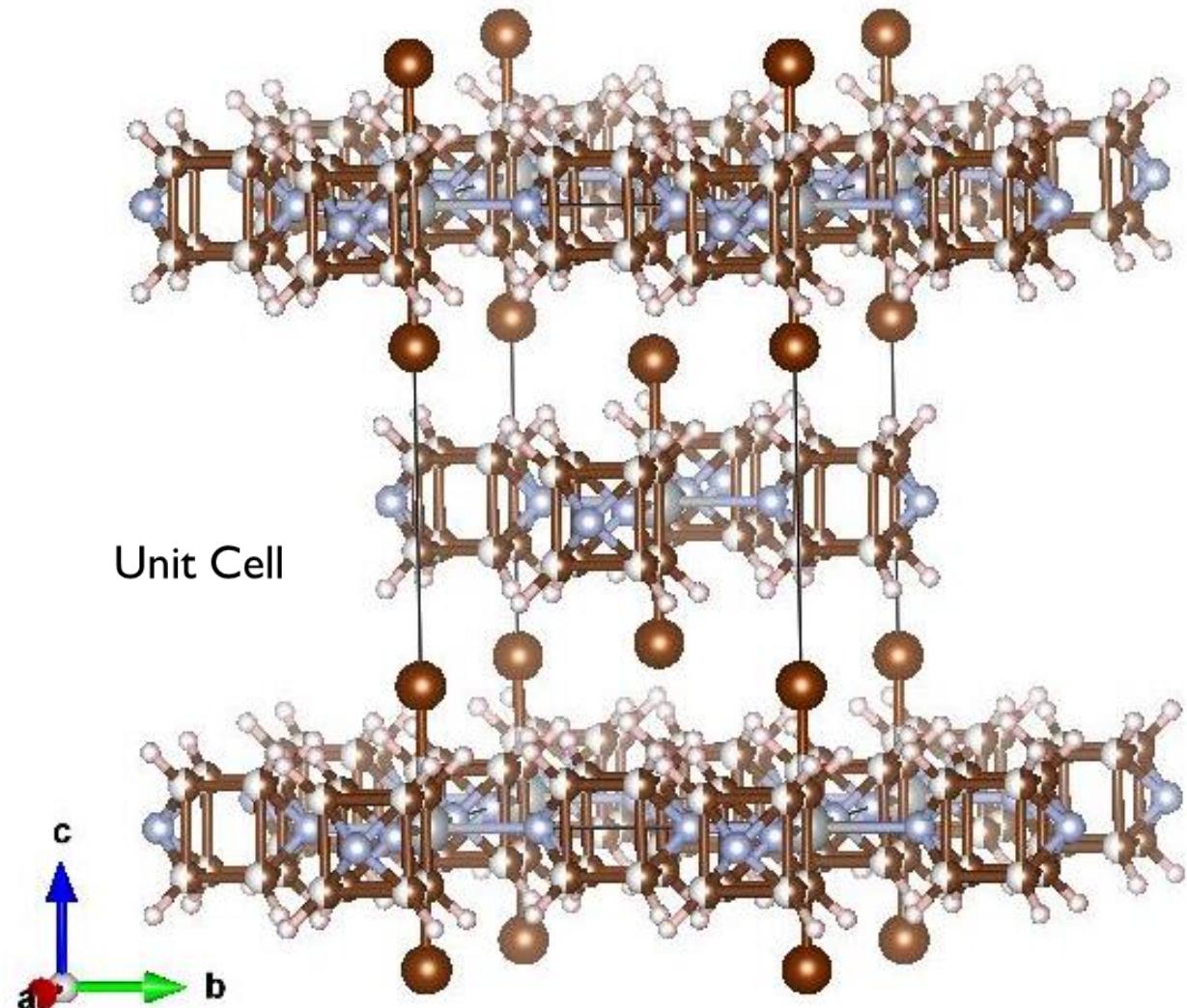


$\text{NiBr}_2\text{pyz}_2$

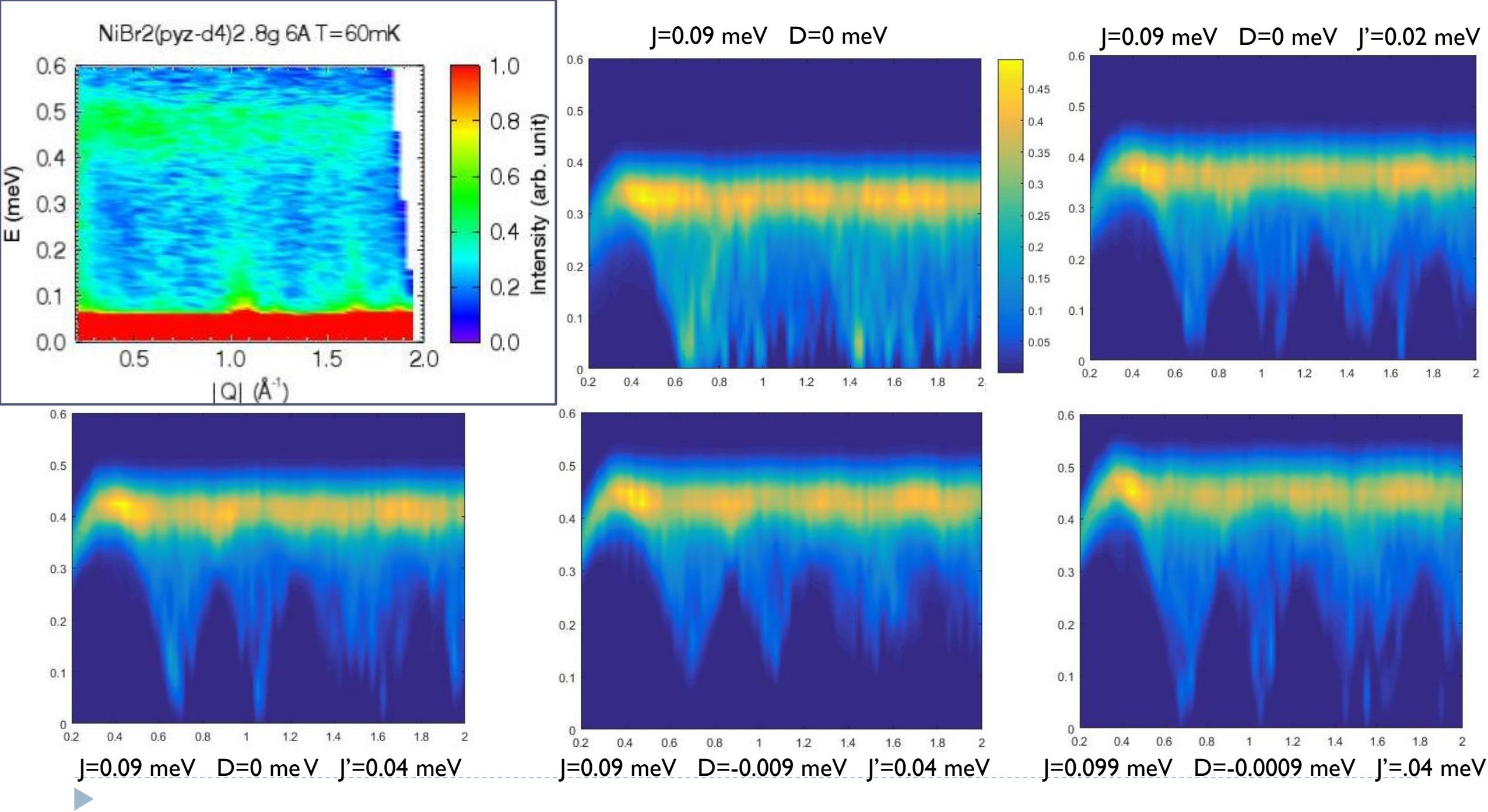
- $S=1$
- Tetragonal
- $T_N=1.9 \text{ K}$



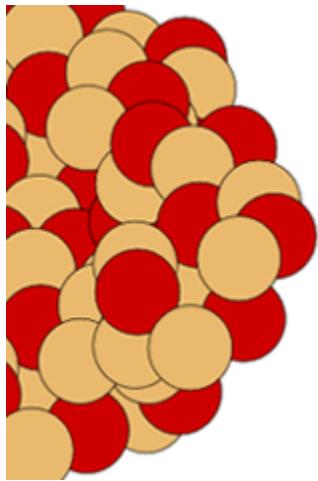
Single Layer



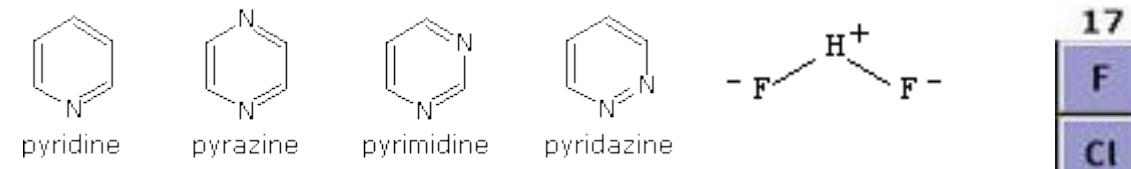
Unit Cell



A Need



Neutron



17
F
Cl
Br
I



Thank You

- ▶ Qingzhen Huang
- ▶ Craig Brown
- ▶ Jamie Manson
- ▶ Joe Dura
- ▶ Julie Borchers
- ▶ Dan Neumann
- ▶ Steve Disseler
- ▶ William Ratcliff
- ▶ Juscelino Leao
- ▶ Yegor Vekhov
- ▶ Alan Ye
- ▶ Terry Udovic
- ▶ Jeff Lynn
- ▶ Yamali Hernandez
- ▶ Jamie Brambleby
- ▶ Paul Goddard



Non-NIST Experimental Partners

- ▶ The University of Warwick
- ▶ Advanced Photon Source, Argonne National Laboratory
- ▶ National High Magnetic Field Laboratory, Los Alamos National Laboratory
- ▶ Swiss Muon Source, Paul Scherrer Institut
- ▶ ISIS Pulsed Neutron and Muon Source

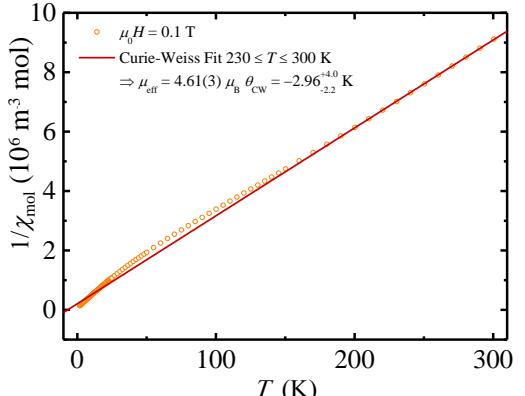
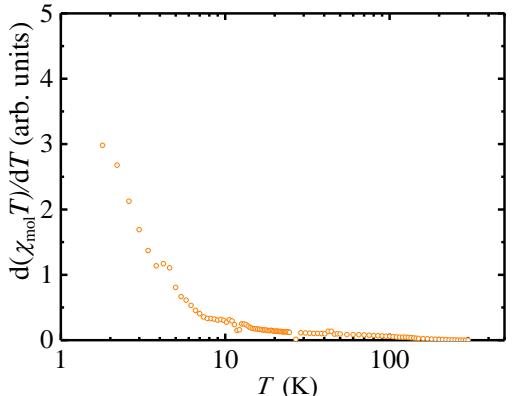
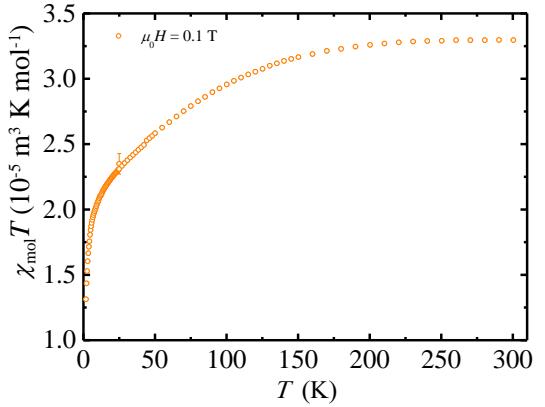
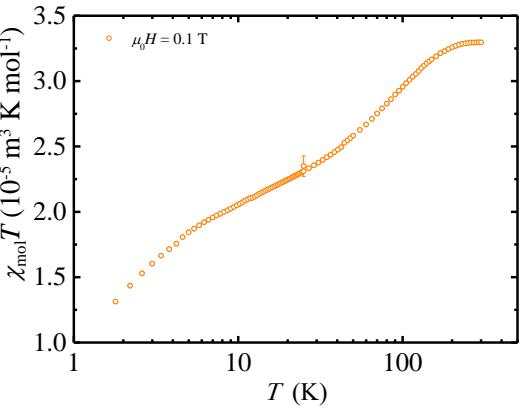
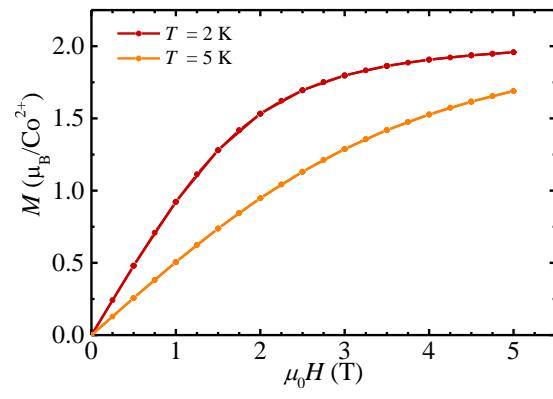
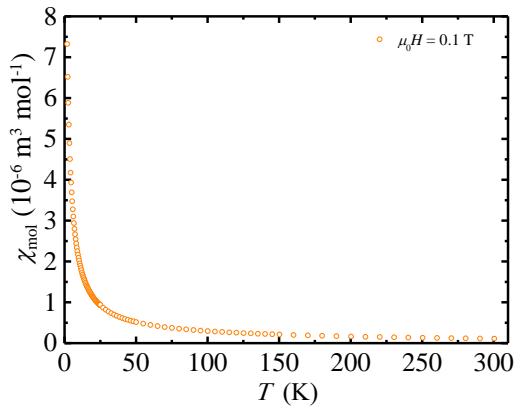
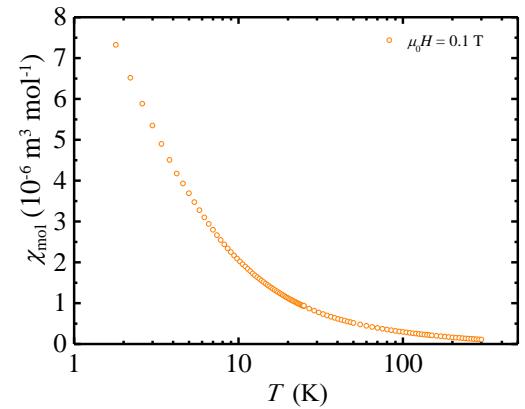


Software Used for Data Analysis

- ▶ **Neutron Diffraction**
 - ▶ CMPR
 - ▶ GSAS
- ▶ **Inelastic Neutron Scattering**
 - ▶ DAVE
 - ▶ SpinW
- ▶ **Visualization**
 - ▶ VESTA

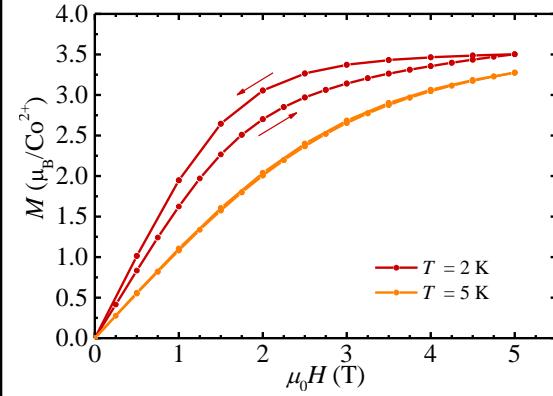


$\text{CoCl}_2\text{pyz}_2$



JLM4-088: $\text{CoCl}_2(\text{pyz-D}_4)_2$
Sample dispersed in
Vaseline

Comparison:
Sample without Vaseline



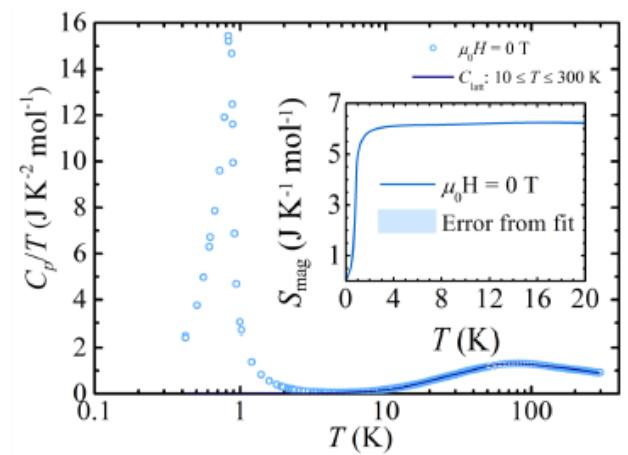


Fig. 1 Heat capacity (plotted as C_p/T) vs. temperature for $\text{Co}_2\text{Cl}_2(\text{pyz-d}_4)_2$ (JLM4-088). A Peak at $0.85(1)$ K indicates a transition to long range order. The solid line is a fit to a model of one Debye plus two Einstein modes in the range $10 \leq T \leq 300$ K. Inset: Subtracting the lattice contribution from the data, the entropy change to 20 K is determined to be consistent with $R \ln 2 = 5.8 \text{ J K}^{-1} \text{ mol}^{-1}$, indicating the sample behaves as an effective spin $1/2$ system at low temperatures.

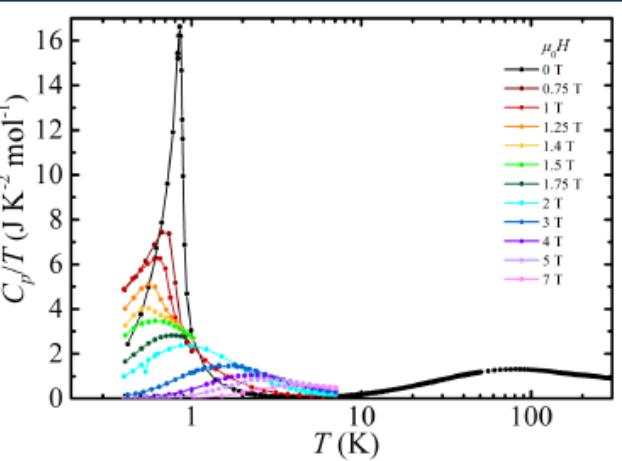


Fig. 3 Field dependence of C_p/T vs. T . The λ -peak moves to lower temperatures as the field is increased, indicative of a transition to an AFM ordered ground state. T_N is too low to be experimentally accessible for applied fields $\mu_0 H \geq 1.5$ T. At higher fields a broad feature emerges which moves monotonically to higher temperatures with applied field.

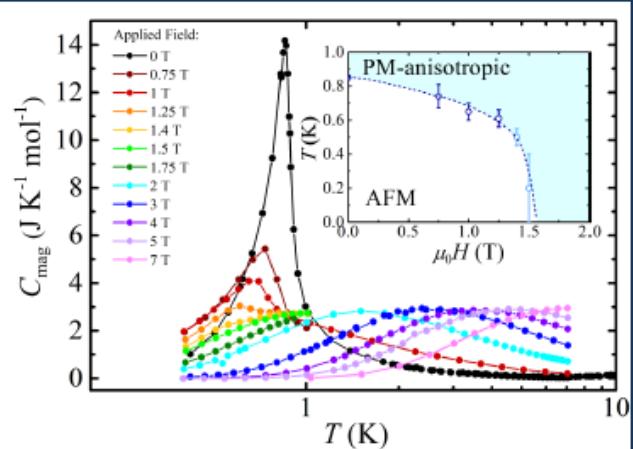


Fig. 4 Magnetic heat capacity (C_{mag}) vs. T , following subtraction of the lattice fit. The broad feature (for $\mu_0 H \geq 3$ T) has a constant amplitude, and a peak position which increases linearly with field. This is highly indicative of a Schottky anomaly due to the field-induced splitting of the doublet ground state in Co^{2+} . Inset: phase boundary separating the paramagnetic state (with easy-plane anisotropy) with the AFM ordered state. Faded points are most easily seen in C_p/T vs. T (Fig. 3).

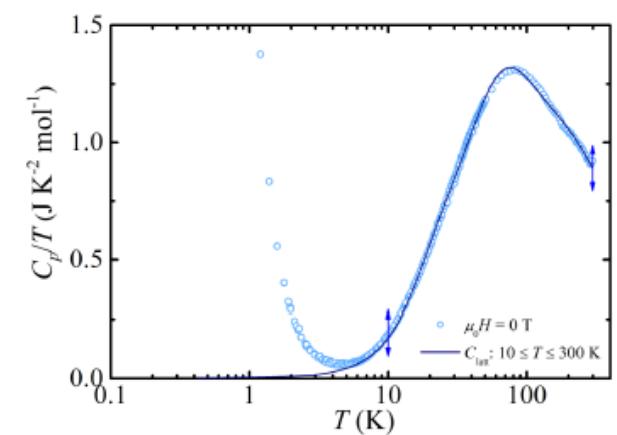
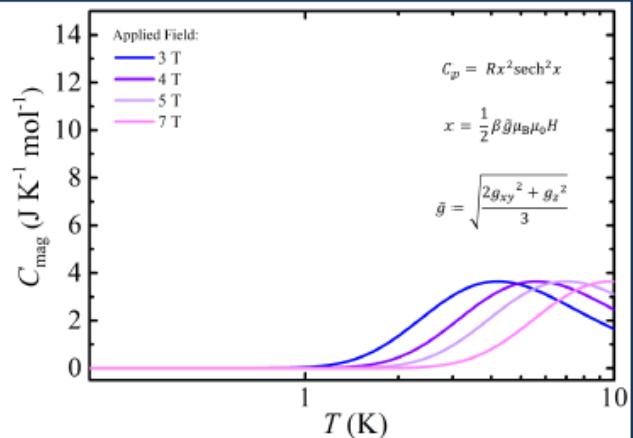
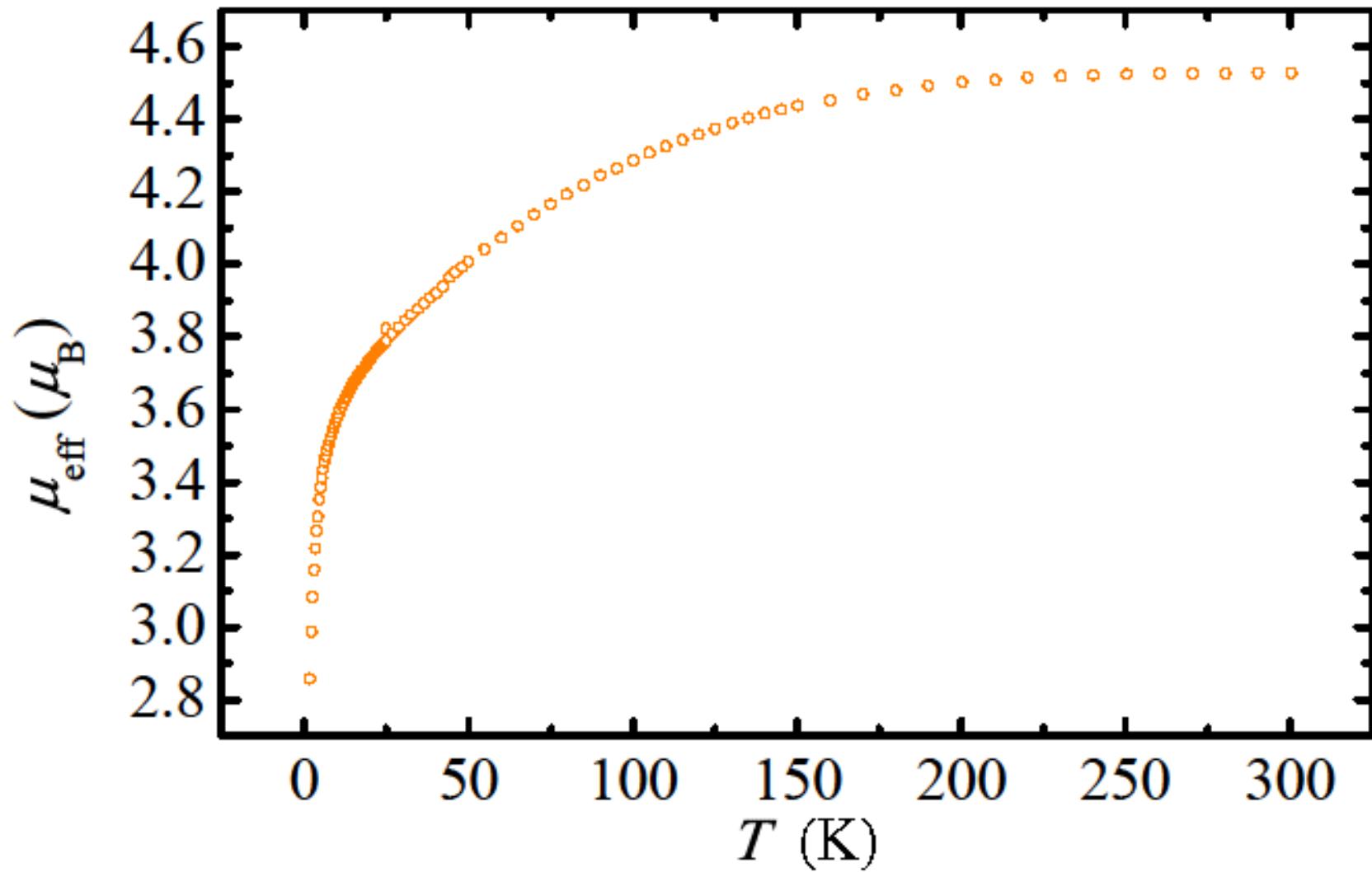


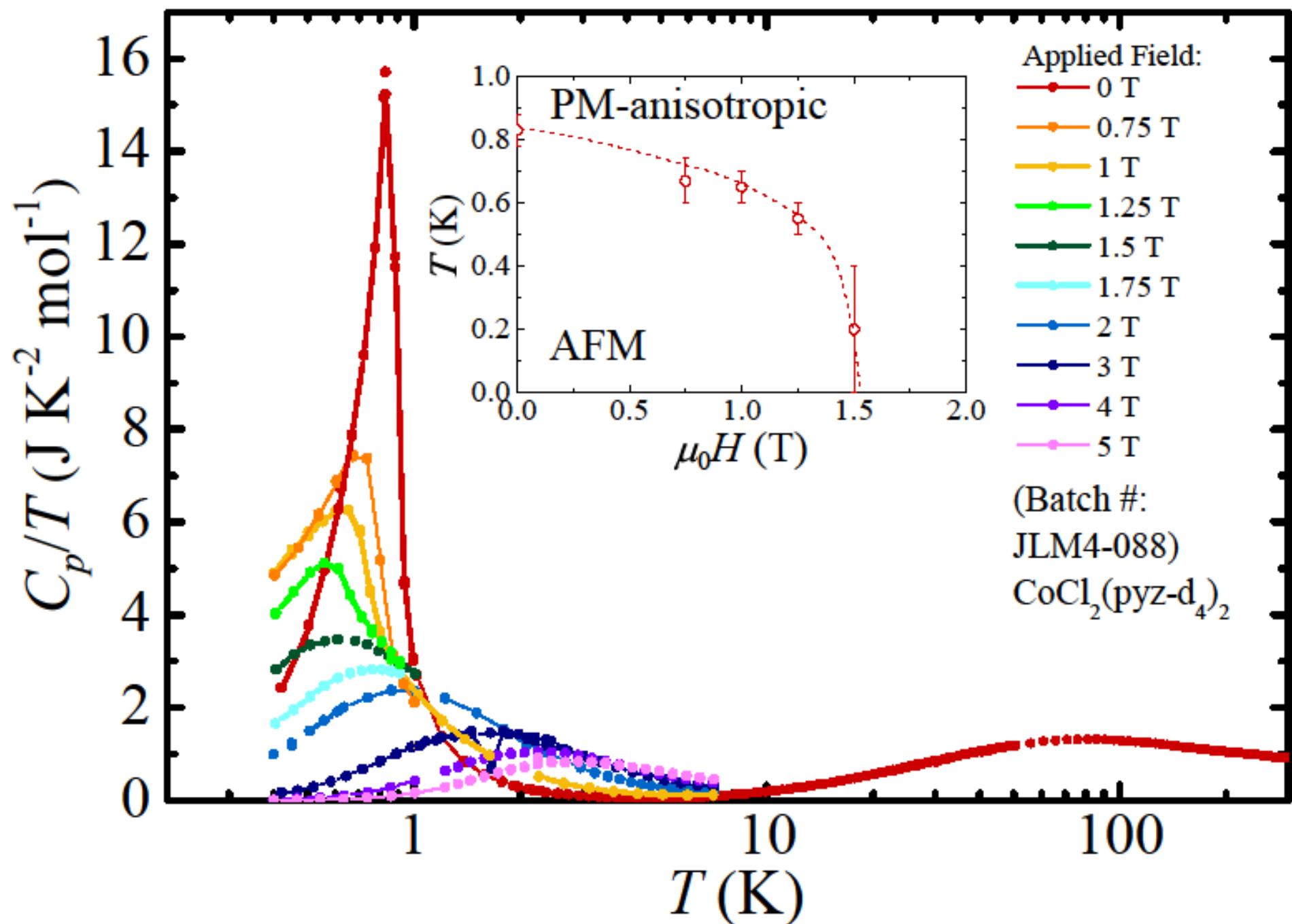
Fig. 2 Heat capacity vs. temperature, showing the lattice fit to the data. The fit only has a small dependence on the lower bound used for the fit, giving rise to the small error in the calculated entropy change (Fig. 1 Inset).

Fig. 5 Simulated molar heat capacity (C_{mag}) vs. T , for a two level system in an applied field. The result is a Schottky anomaly, with the equation given in the figure. The g -factor was taken to be the powder average g -factor, assuming the published^[1] values for the fully hydrogenated phase of $g_{xy} = 5.98$ and $g_z = 1.97$. This model is a good representation of the measured data (Fig. 4), having the approximately the correct amplitude and field dependence for the maxima.

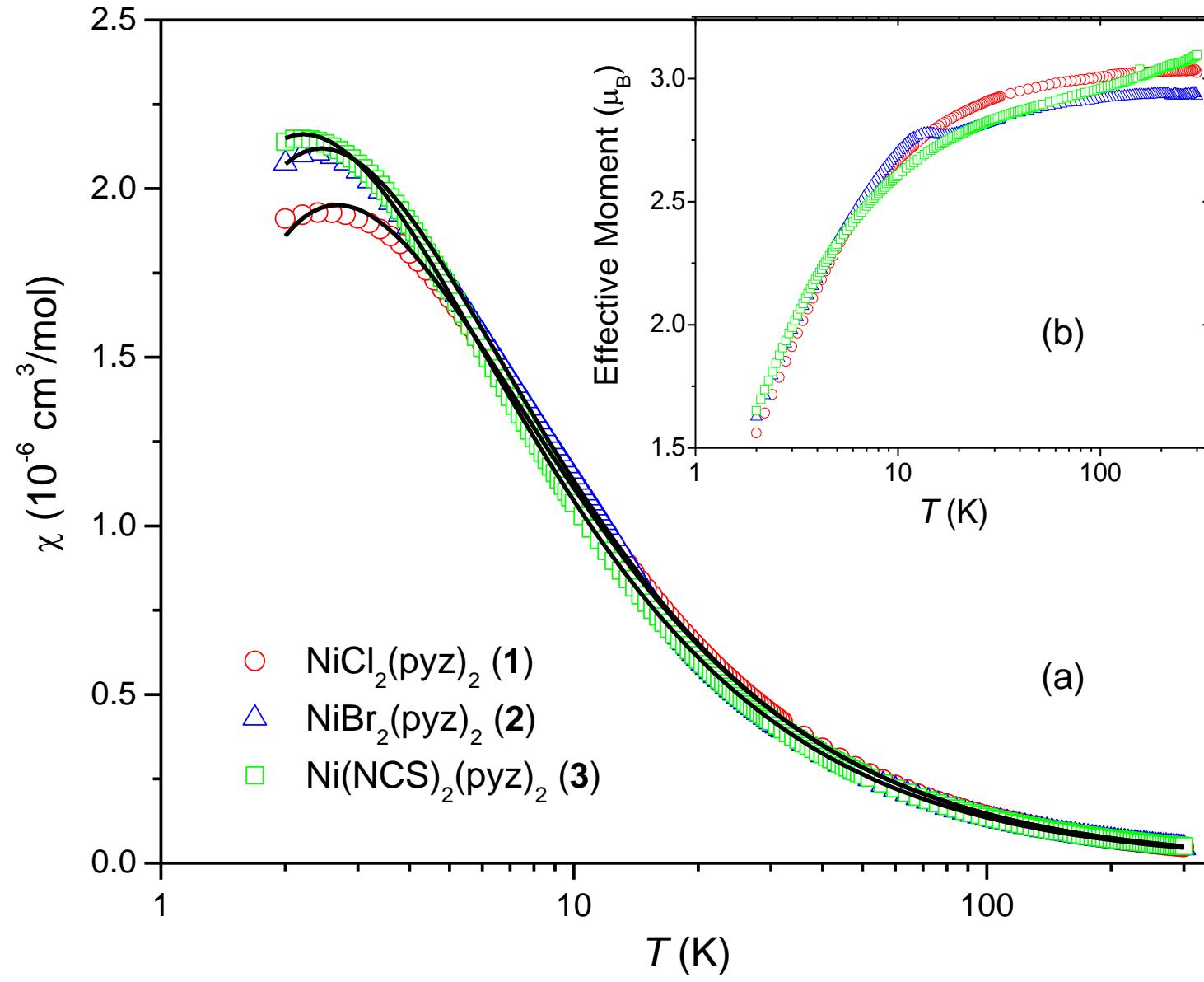
[1] R. L. Carlin *et al.*, PRB **32**, 7476 (1985).

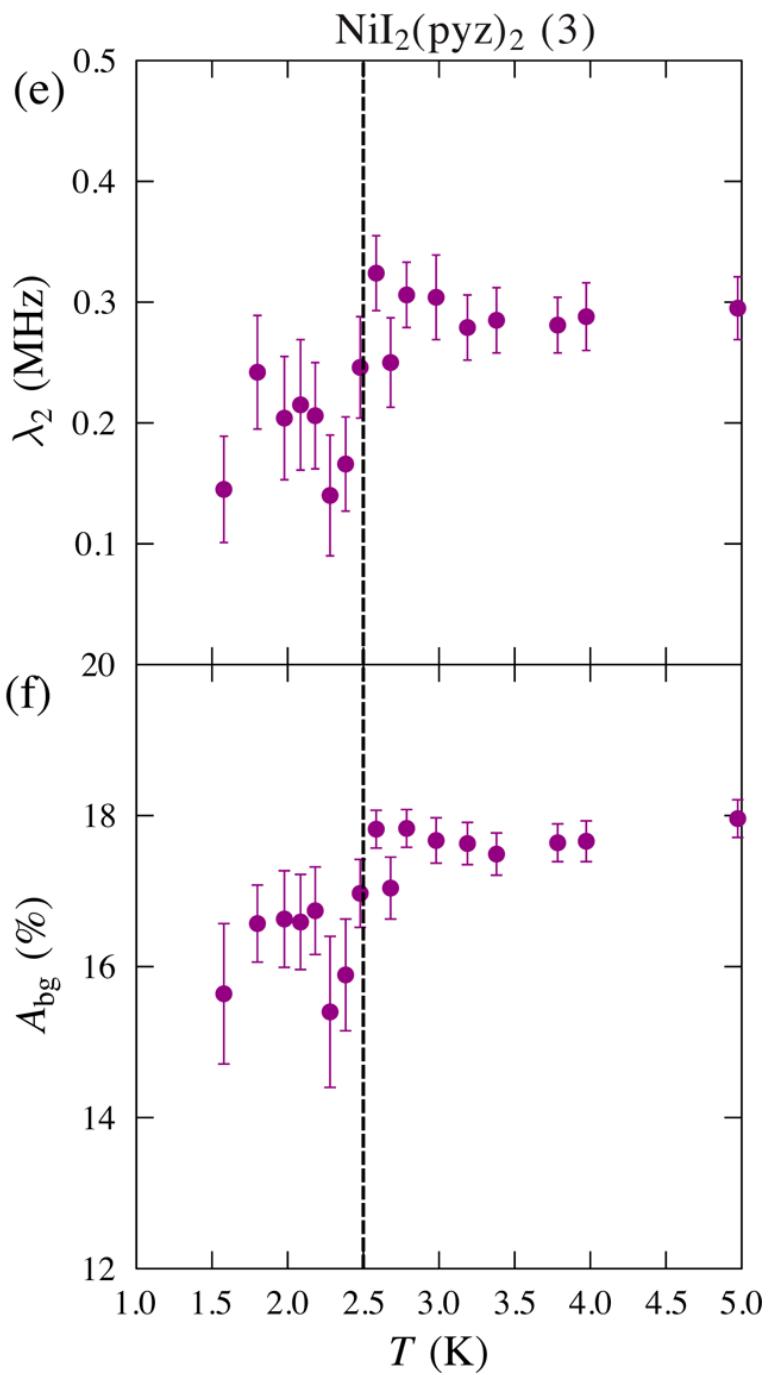
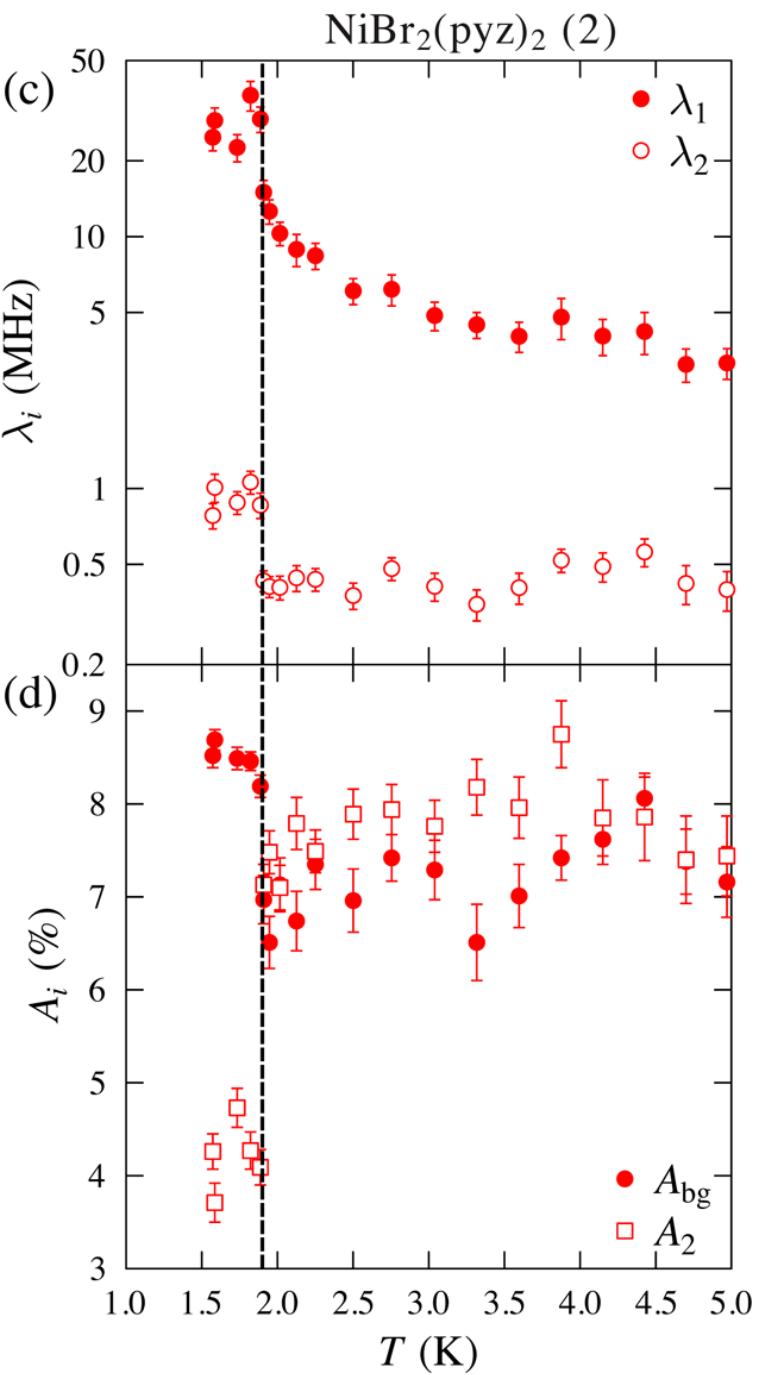
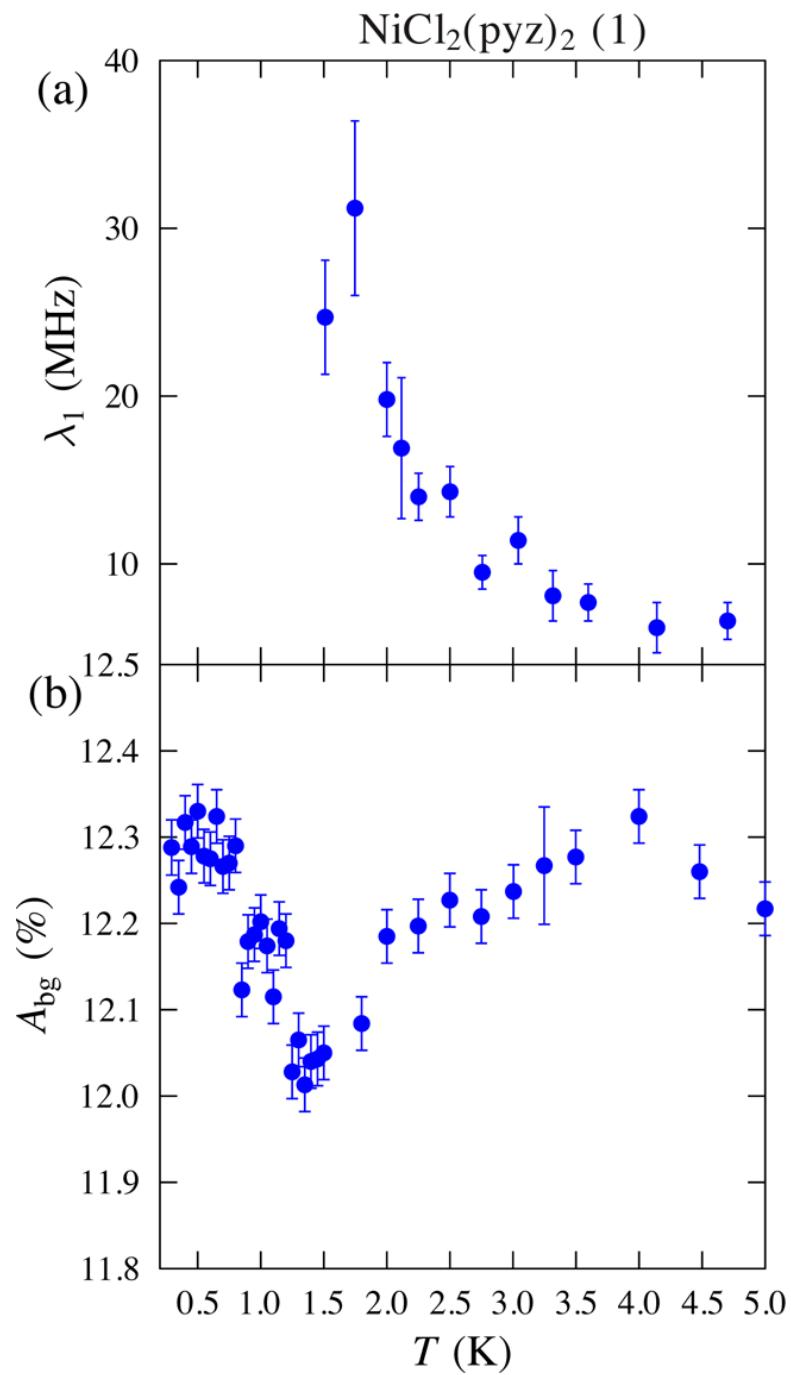






$\text{NiBr}_2\text{pyz}_2$





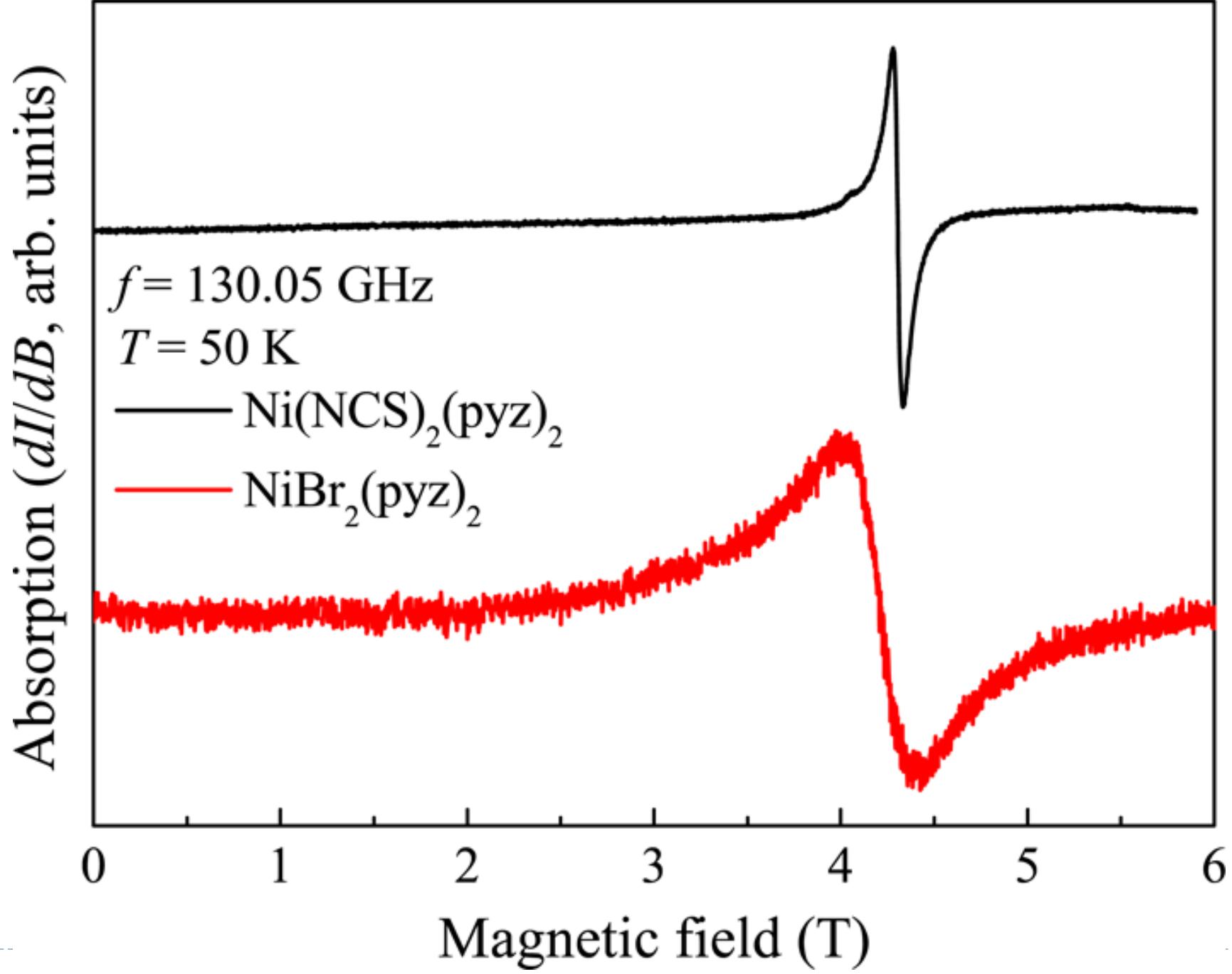


Fig. 1

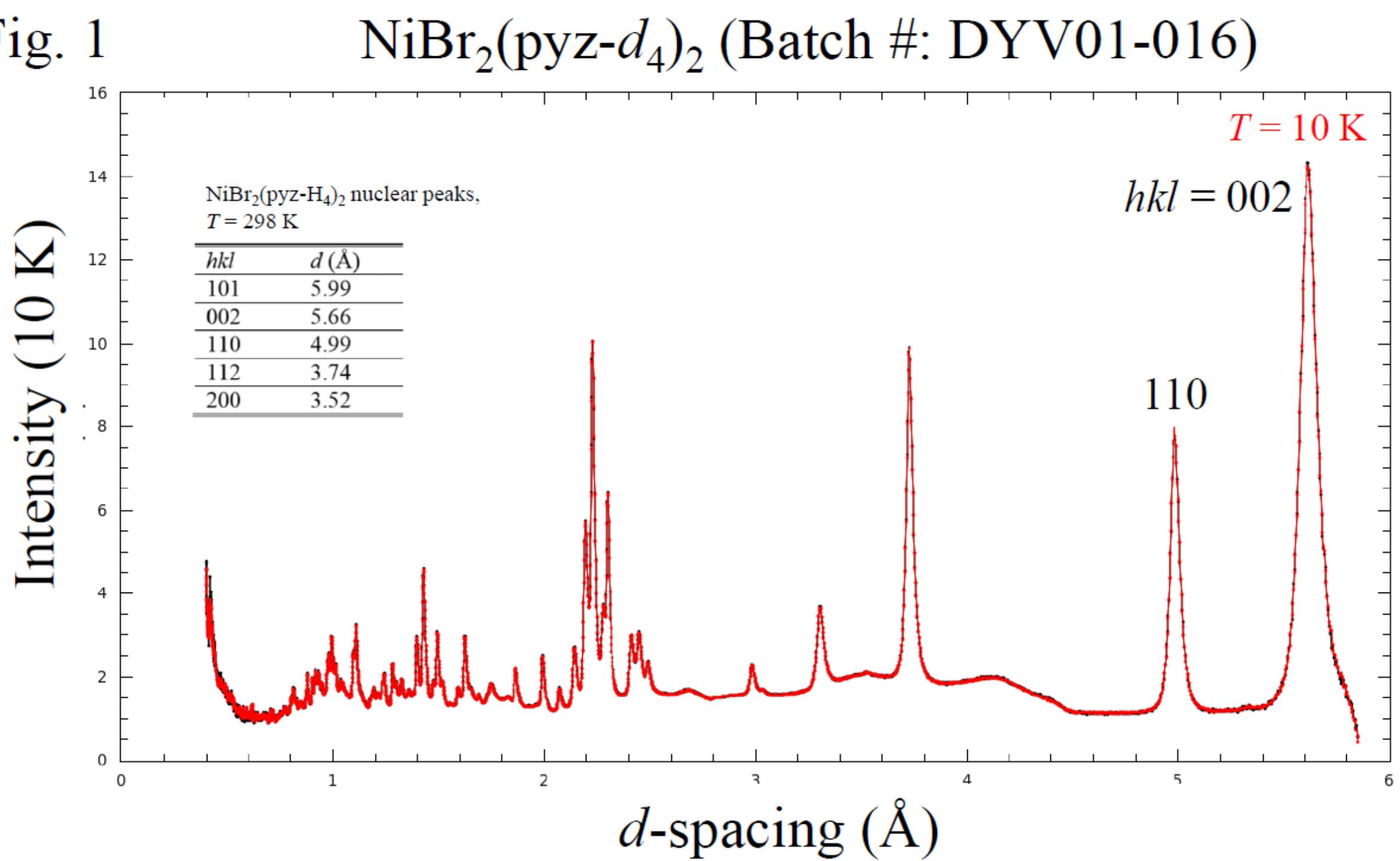


Fig. 2

Intensity (0.3 K - 10 K)

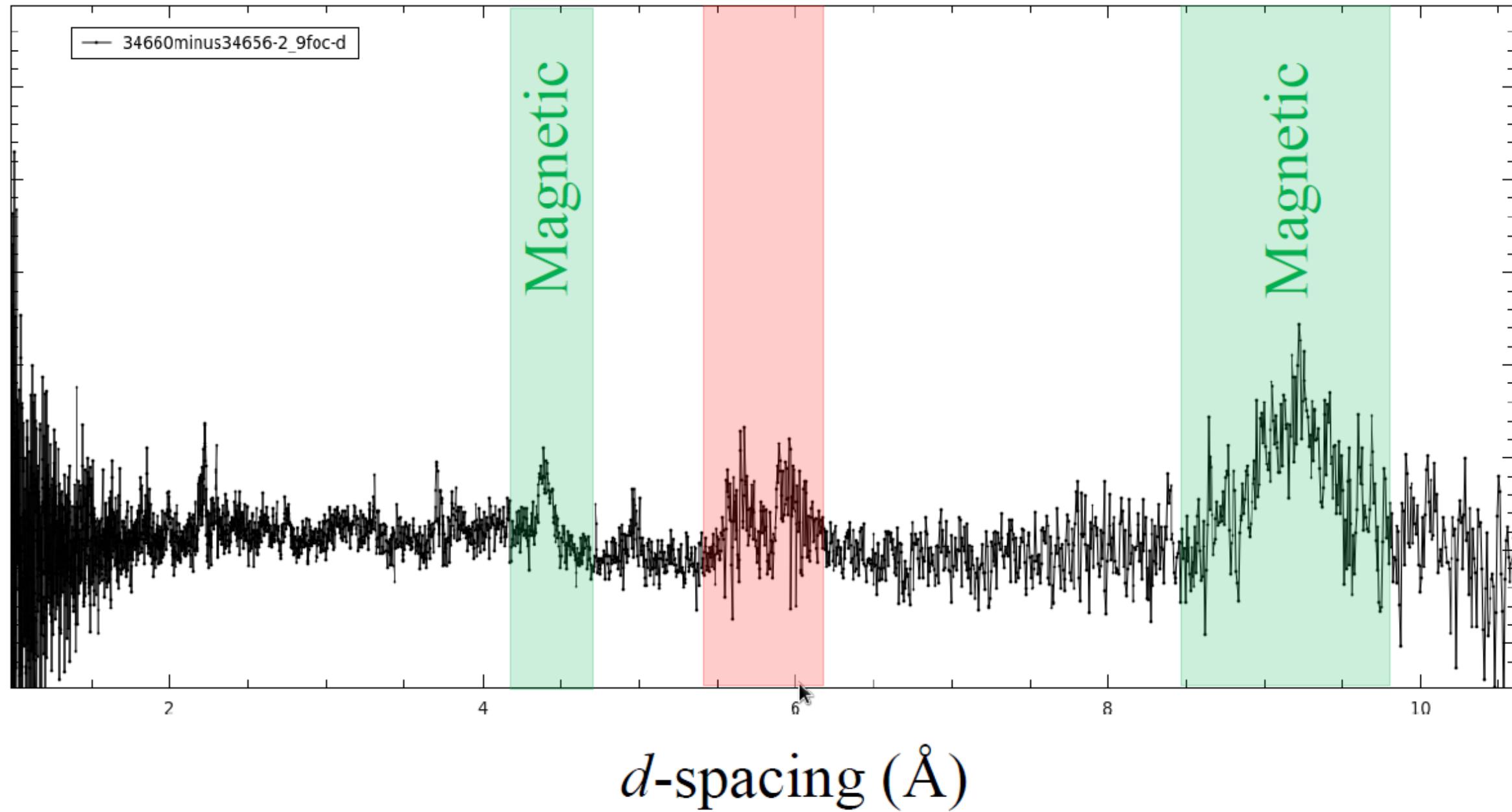
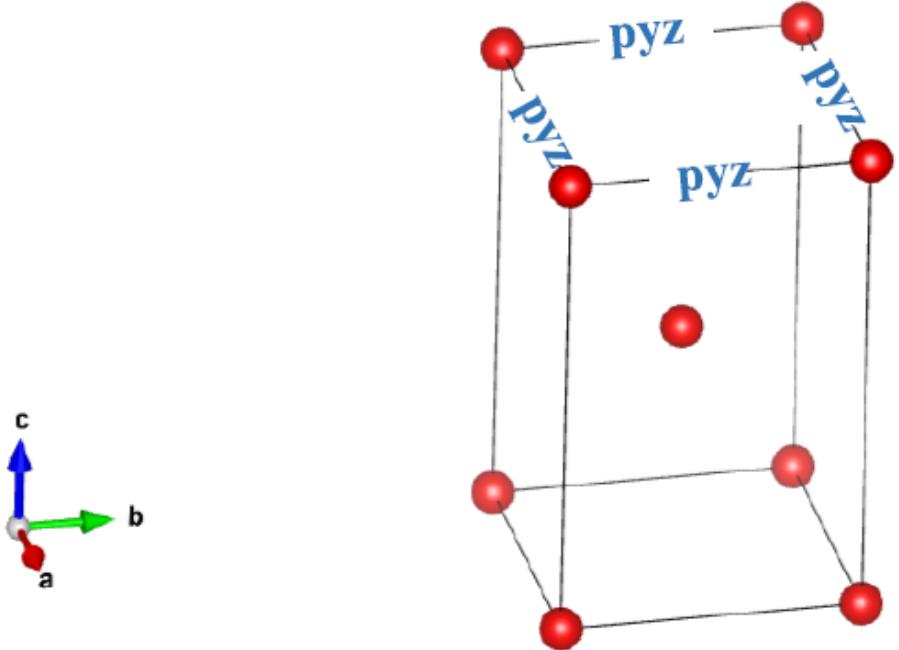


Fig. 3

(a)



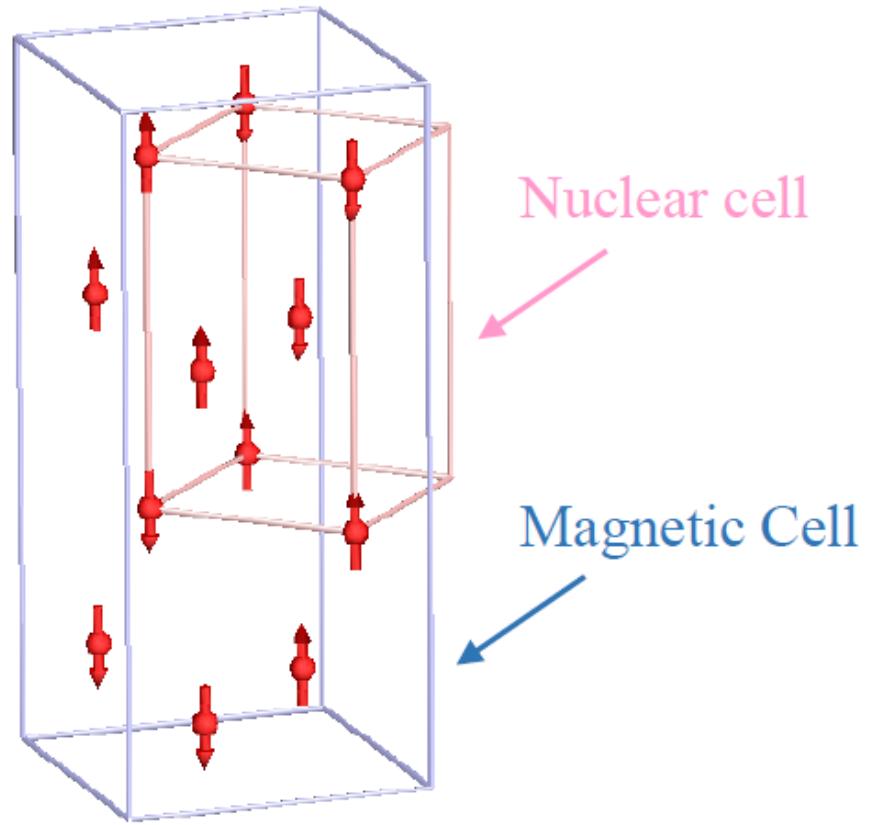
$$a = b = 7.09 \text{ \AA}$$

$$c = 11.3 \text{ \AA}$$

*I*4/*mmm* (tetragonal)

$$T = 298 \text{ K}$$

(b)



$$a = b = 7.04 \text{ \AA}$$

$$c = 11.2 \text{ \AA}$$

*I*4/*mmm* (tetragonal)

$$T = 0.3 \text{ K}$$

$$\boldsymbol{k} = [0.5 \ 0.5 \ 0.5]$$

