

Last slide of my presentation at this workshop on February 7, 2006:

## Conclusions

- Once very fashionable, the area of GB diffusion is not **hot** anymore. It is **not** considered to be **cool** enough. It cannot compete with carbon nanotubes and quantum dots
- It is only the Herzig group in Muenster that keeps GB diffusion measurements alive
- Development of GB diffusion theory stopped (June 1, 2005)
- Posterity will not forgive us

# Grain boundary diffusion

## The rest of the story

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### *Outline*

- Quick overview of the field
- Diffusion along grain boundaries and their triple junctions in copper
- Diffusion along low-angle grain boundaries in aluminum

## Importance of GB diffusion

GB diffusion is much faster than lattice diffusion  
(E.g.  $D_{gb}/D_L \approx 10^{10}$  at  $0.5T_m$  in fcc metals)

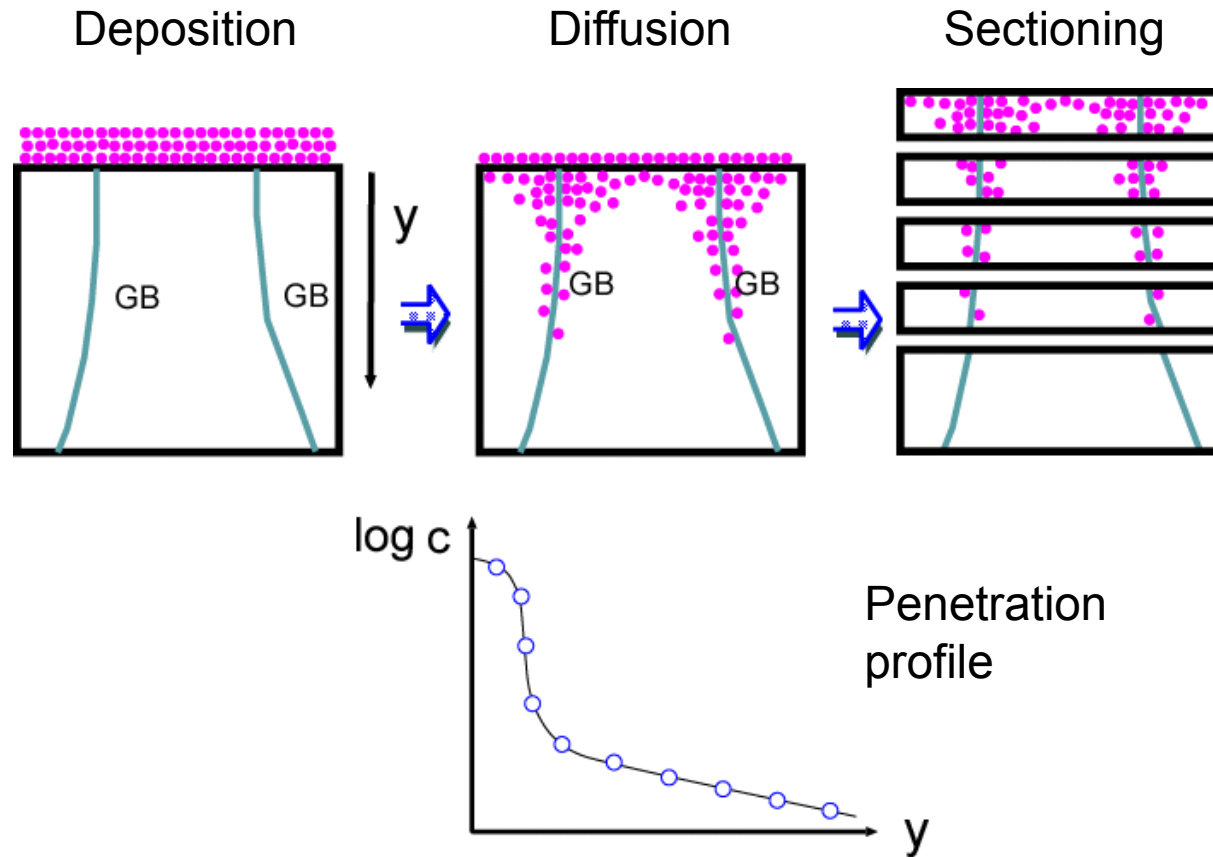


Processes controlled/influenced by GB diffusion:

- Solid state reactions (discontinuous precipitation,...)
- Grain growth
- Deformation and fracture at elevated temperatures
- Coble creep
- GB dislocation climb
- Structural relaxation after fabrication (severe plastic deformation, etc.)

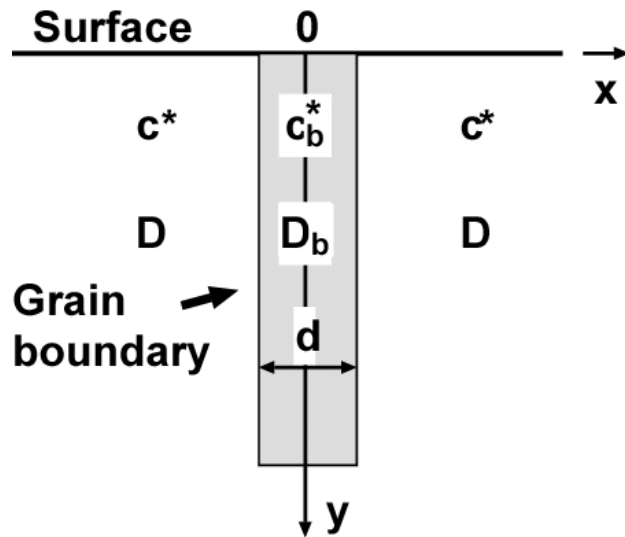
Diffusion is a structure-sensitive property

# Radiotracer sectioning method



Need many GBs to achieve a high accuracy

## Fisher model of GB diffusion



$$\frac{\partial c^*}{\partial t} = D \left( \frac{\partial^2 c^*}{\partial x^2} + \frac{\partial^2 c^*}{\partial y^2} \right)$$

$$\frac{\partial c_b^*}{\partial t} = D_b \frac{\partial^2 c^*}{\partial y^2} + \frac{2D}{\delta} \left( \frac{\partial c^*}{\partial x} \right)_{x=\delta/2}$$

### Coupling conditions:

Self-diffusion  $A^* \rightarrow A$ :  $c_b^* = c^*$

Impurity diffusion  $B^* \rightarrow A$ :  $c_b^* = s c^*$ , where  $s = s_0 \exp(-E_s/kT)$

Self-diffusion in alloy  $B^* \rightarrow A-B$ :  $c_b^* = s c^*$ , where  $s = c_b^B / c^B$

# Solution of the Fisher model

Under typical experimental conditions

$$\log \bar{c} \propto - \left[ \frac{4\pi D}{(s\delta D_b)^2} \right]^{3/10} y^{6/5}$$

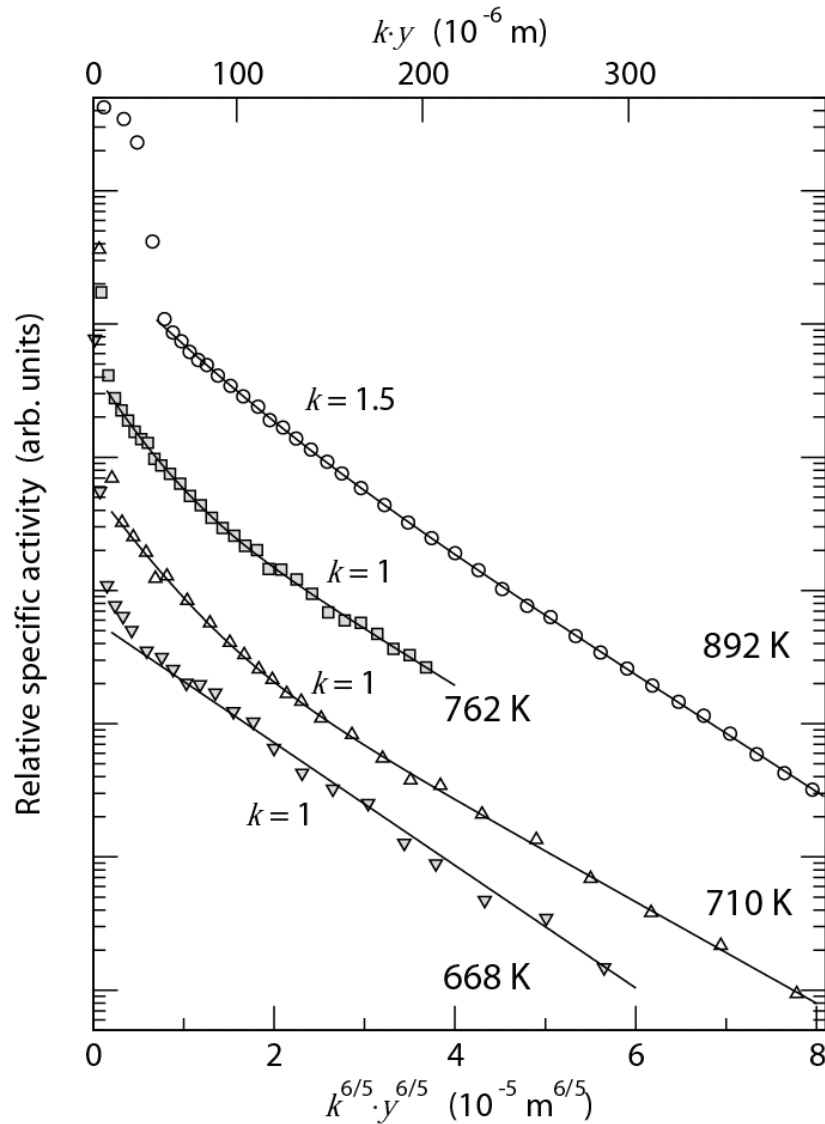


$$s\delta D_b = q \left( \frac{D}{t} \right)^{1/2} \left( - \frac{\partial \ln \bar{c}}{\partial y^{6/5}} \right)^{-5/3}$$

$q$  – numerical factor depending on the surface condition

- Assume  $\delta = 0.5$  nm
- Must know  $D$  from independent measurements

# Examples of GB diffusion profiles



Ag in Cu-0.2at%Ag

Divinski et al., *Interface Science* **11**,  
 21 (2003)

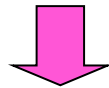
## Combined B and C regime measurements

**Regime B** (high temperatures)

$$(Dt)^{1/2} \gg s\delta \quad \rightarrow \quad \bar{c}(y, s\delta D_b) \quad \rightarrow \quad s\delta D_b$$

**Regime C** (low temperatures – extremely difficult measurements!)

$$(Dt)^{1/2} \ll s\delta \quad \rightarrow \quad \bar{c} \propto \exp\left(-\frac{y^2}{4D_b t}\right) \quad \rightarrow \quad D_b$$

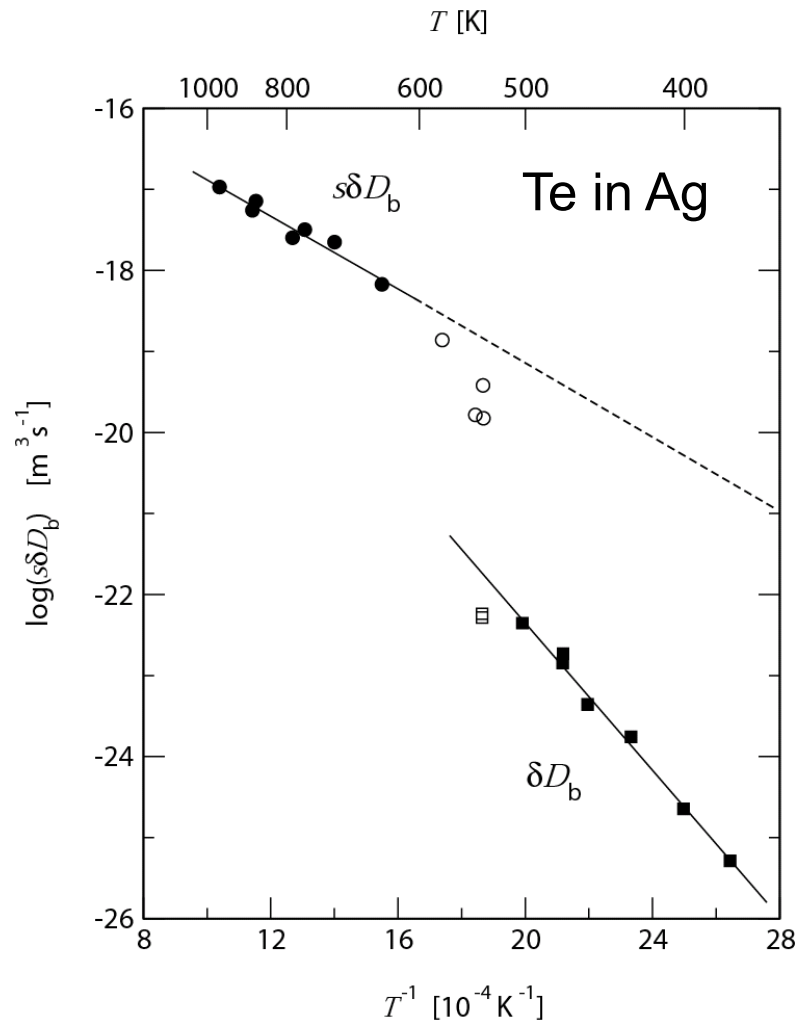


$$s\delta = \frac{(s\delta D_b)_B}{(D_b)_C}$$

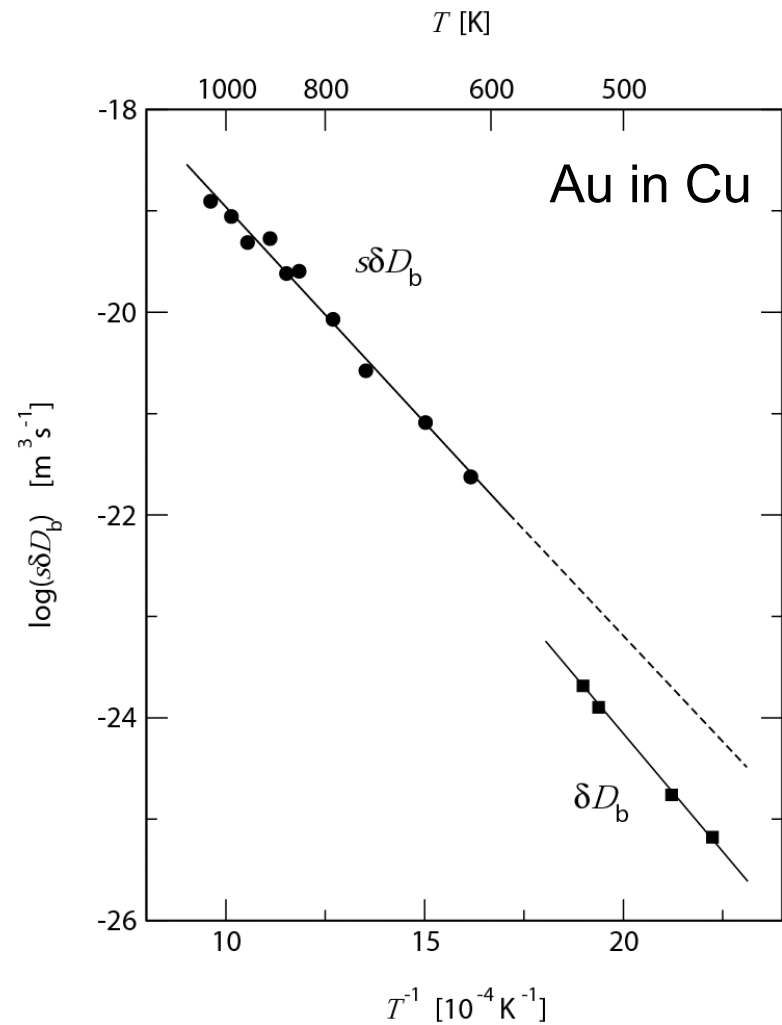
**Self-diffusion:**  $s = 1 \rightarrow \delta \approx 0.5 \text{ nm}$  (Atkinson and Taylor 1981; Sommer and Herzig 1992; Gas, Beke and Bernardini 1992)



## Example of combined B and C regime measurements

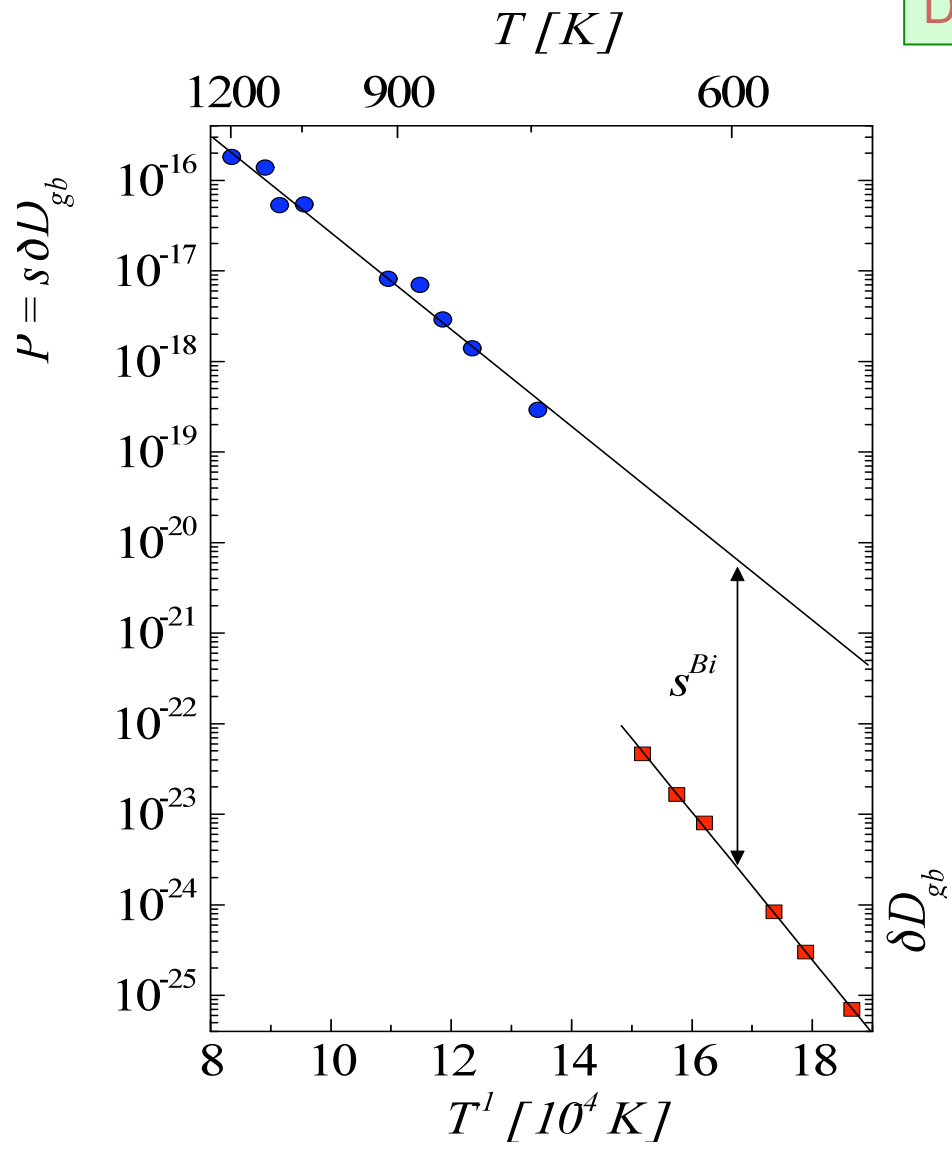


C. Herzig et al., *Acta Mater.* **41**, 1683 (1993)



T. Surholt et al., *Phys. Rev. B.* **50**, 3577 (1994)

Divinski, Herzig, et al Acta Mater (2004)



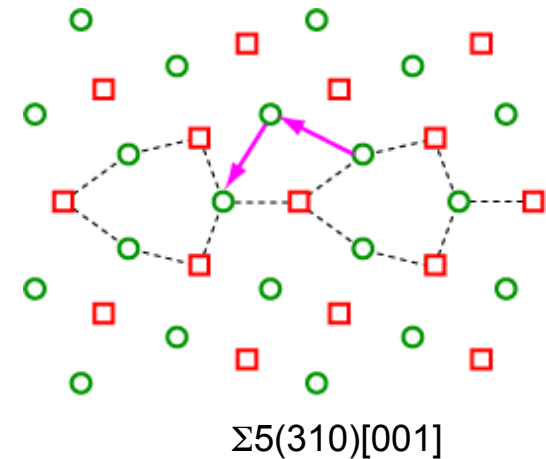
Bi GB diffusion in Cu



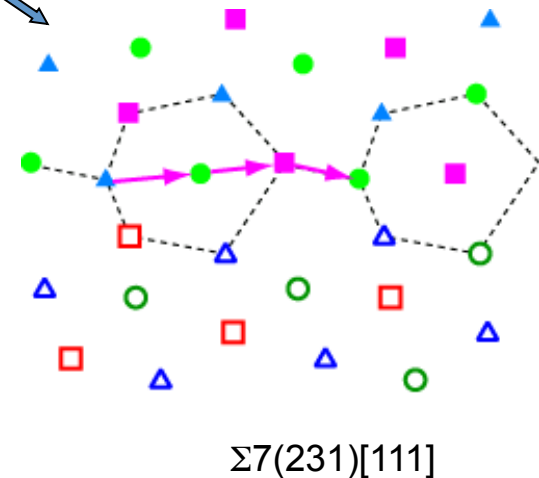
Non-destructive evaluation of GB segregation factors

## Diffusion mechanisms in GBs

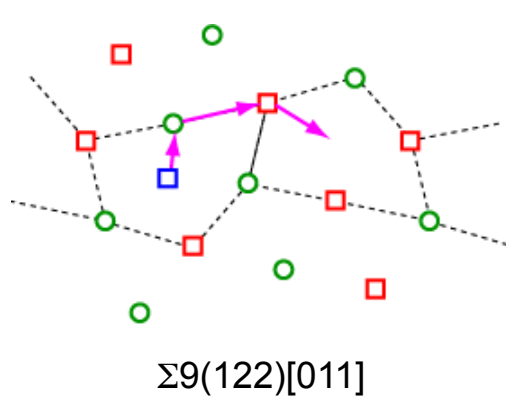
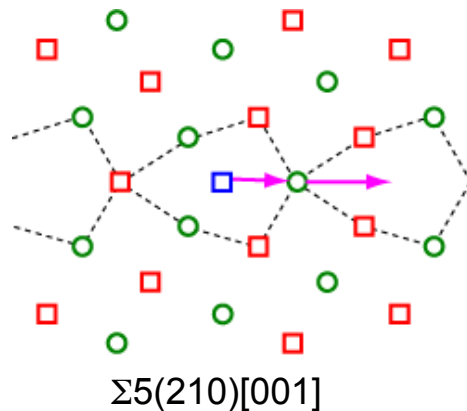
- Vacancy mechanisms
  - Simple vacancy-atom exchanges
  - Long vacancy jumps (2-3 atoms)
- Interstitial mechanisms
  - Direct jumps
  - Collective jumps (2-4 atoms)
- Ring mechanisms (up to 6 atoms)



Vacancies 



Interstitials



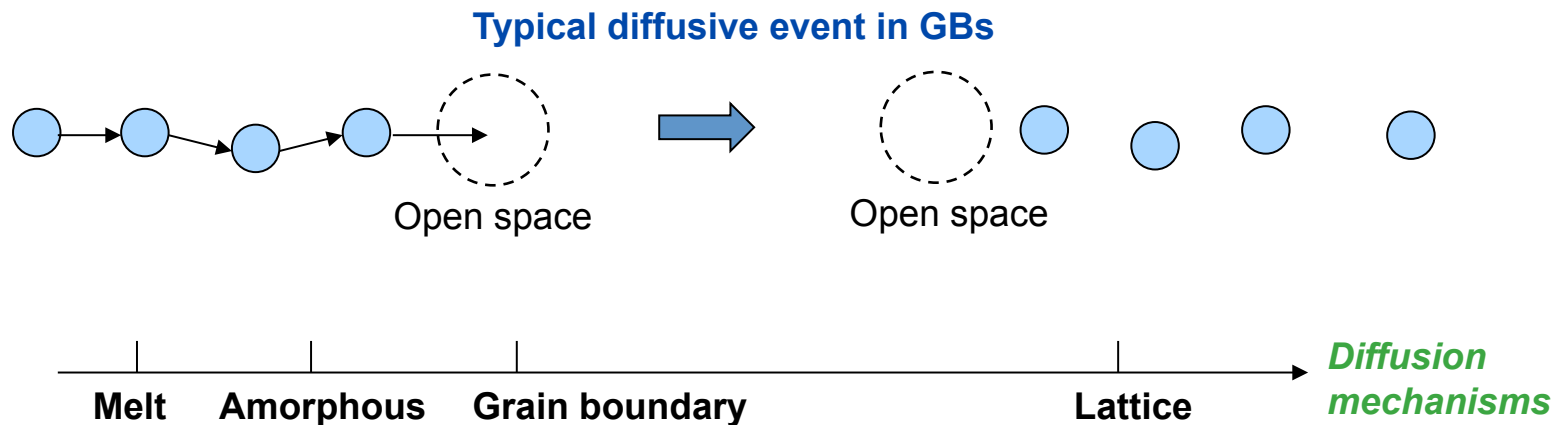
## Comparison of GB and lattice diffusion

### Lattice diffusion

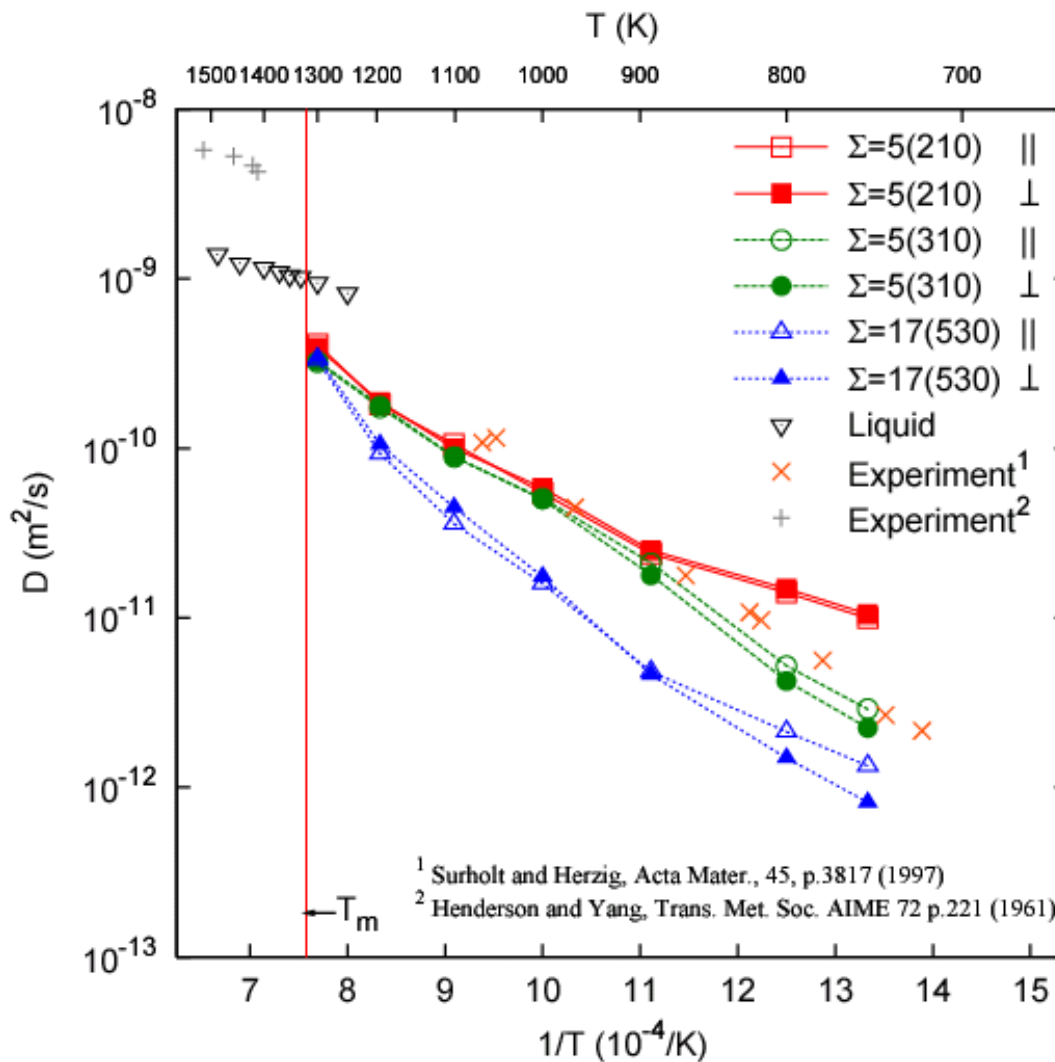
- Vacancies dominate
- Vacancies move by single-atom exchanges

### GB diffusion

- Vacancies and interstitials are equally important
- Variety of point-defect structures
- Variety of diffusion mechanisms
- Most diffusive events are collective



## GB diffusion in Cu: MD calculations



- Agreement with experiment
- Continuous “premelting” ~100K before  $T_m$
- The  $\Sigma 5$ 's merge at high temperatures. Universal diffusion mechanism? “Liquid-like” structure? [Kebllinski et al, 1997, 1999]

High-temperature mechanisms remain unknown !

A. Suzuki and Y. Mishin, *Journal of Materials Science* **40**, 3155 (2005)

# Diffusion along GBs and triple junctions (TJ) in Cu

Tim Frolov

# Triple junction (TJ) is a line where three grain boundaries (GBs) meet together

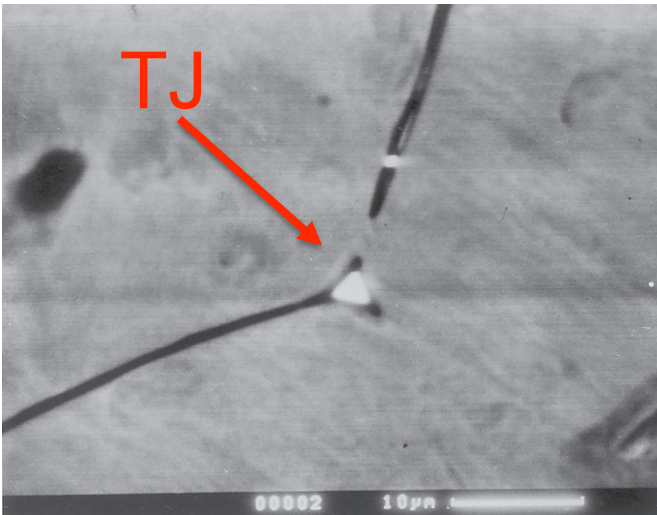
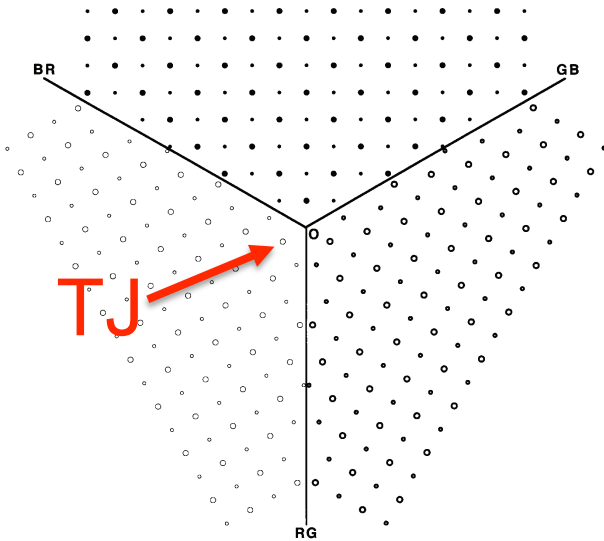


Fig. 1. Electron microphotography (BEI) of sockets for the accelerated penetration Bi in Cu along triple junctions.

- TJs may play a role in behavior of nano-crystalline materials  
V.B. Rabukhin et al (1986), B. Bokstein et al (2001)
- Enhanced diffusivity of nano-crystalline materials
  - Plastic deformation  
L.M. Klinger et al (1997), A.A. Fedorov et al (2002), Ovid'ko et al (2004), H. Wang et al (2005), Y. Chen et al (2007)
  - Diffusional creep
- Sites of enhanced segregation  
K. M. Yin et al (1997)
- Can limit GB mobility during recrystallization
- Preferable sites of nano-cracks and corrosion  
G. Palumbo et al (1989), W.M. Kane et al (2008), L.S. Kumar et al (2003)



# Simulation Methodology

- Copper as a model material

- EAM potential for copper

(lattice parameter, cohesive energy, elastic constants, phonon frequencies, thermal expansion, lattice-defect energies, predicted melting point **1327K**)

- Simulation method

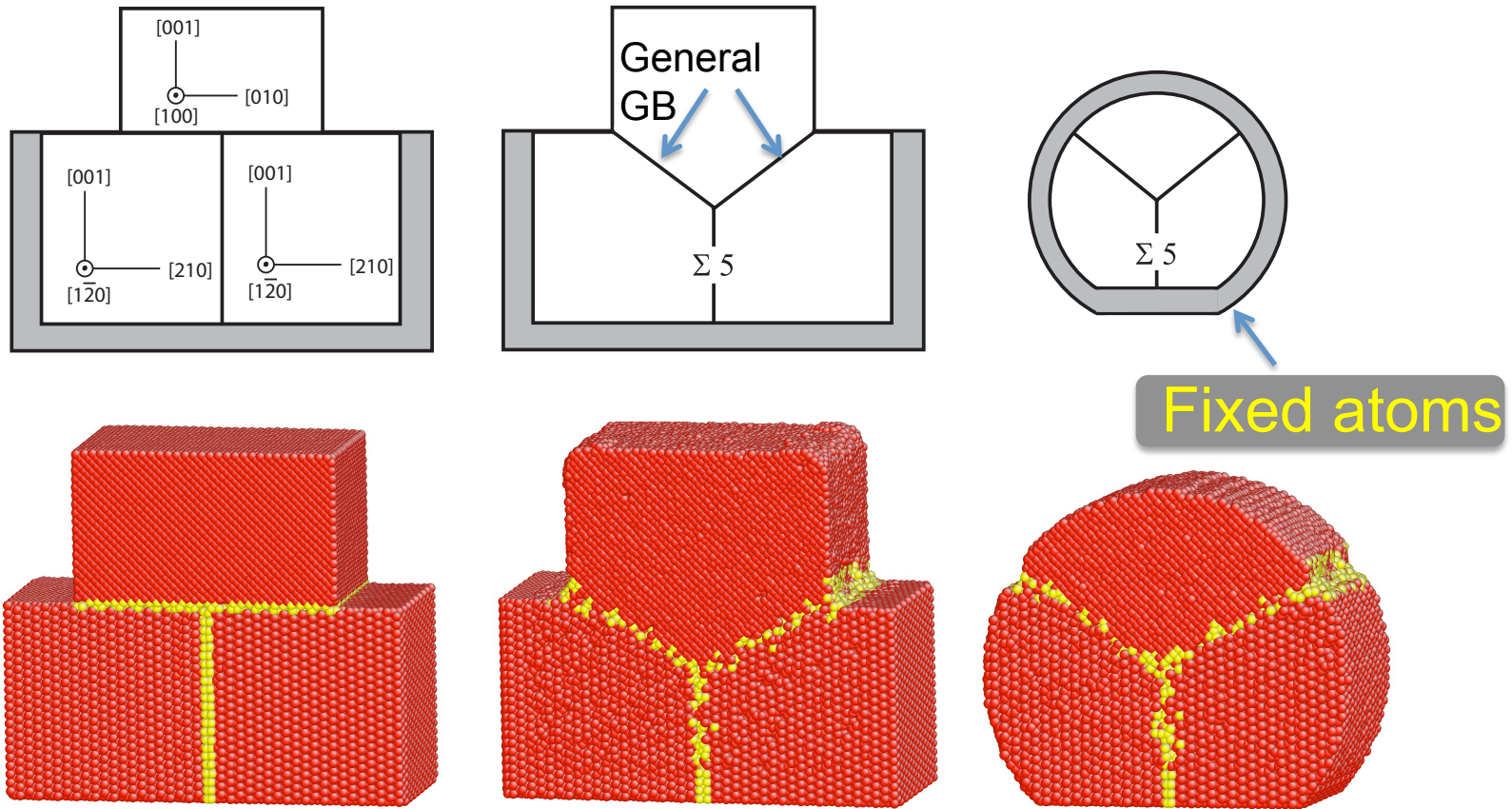
- Molecular Dynamics (MD) ([IMD, Stuttgart, Germany](#))
- Nose-Hoover thermostat
- 81,137 atoms
- Temperature range 700K to 1315K

- Visualization

- Atomeye ([J. Li, Modelling Simul. Mater. Sci. Eng. 11 \(2003\) 173](#))

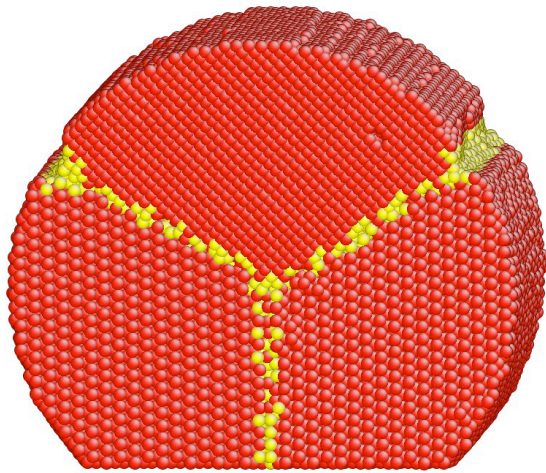


# Creation of a triple junction

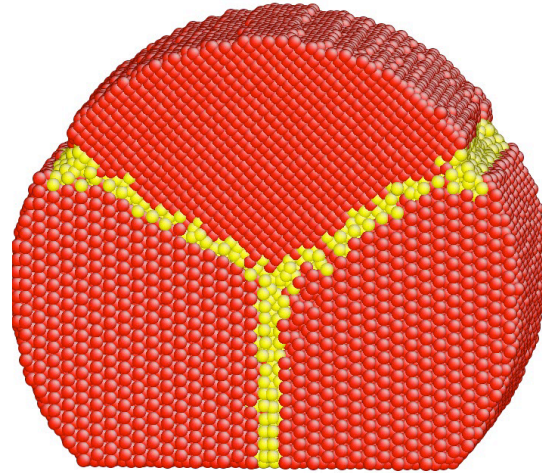


- Three grains with known orientation are brought together
- The structure equilibrates during MD run
- Cylindrical region is cut out from the original block
- Boundary conditions are **periodic** in the direction parallel to TJ

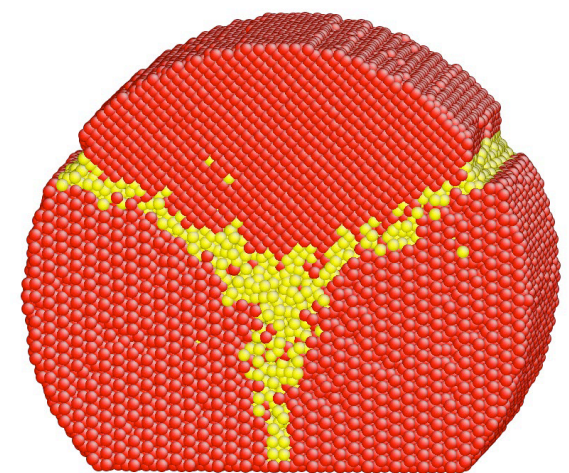
# Simulation block at different temperatures



700 K



1100 K

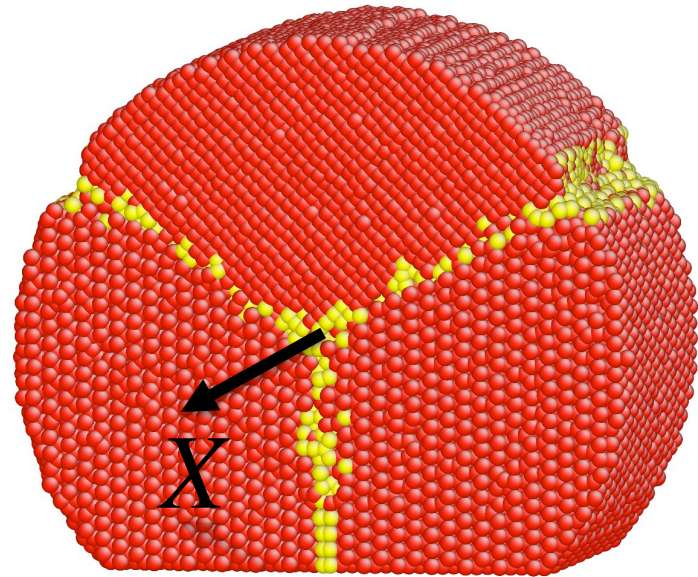


1315 K

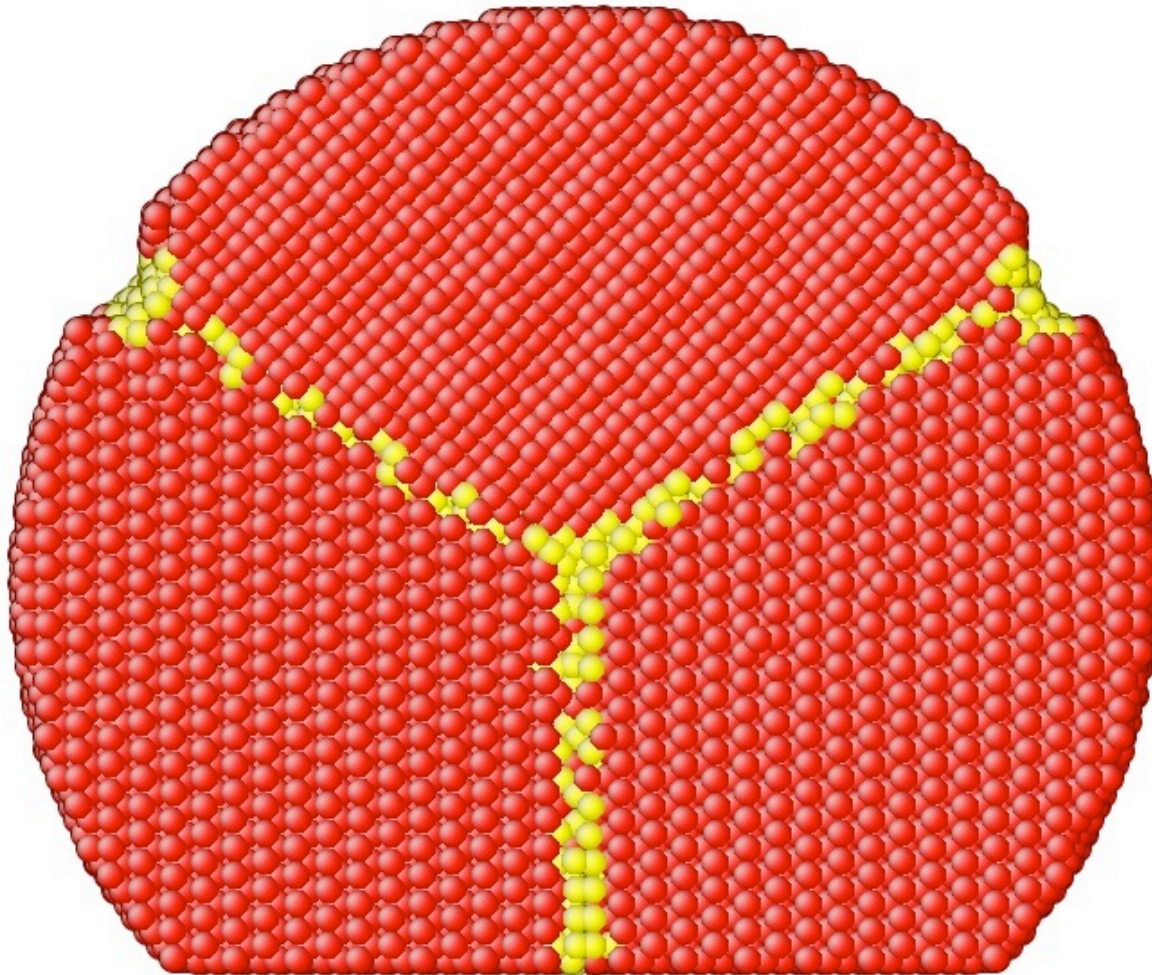
$T_m = 1327$  K

# Diffusion calculations

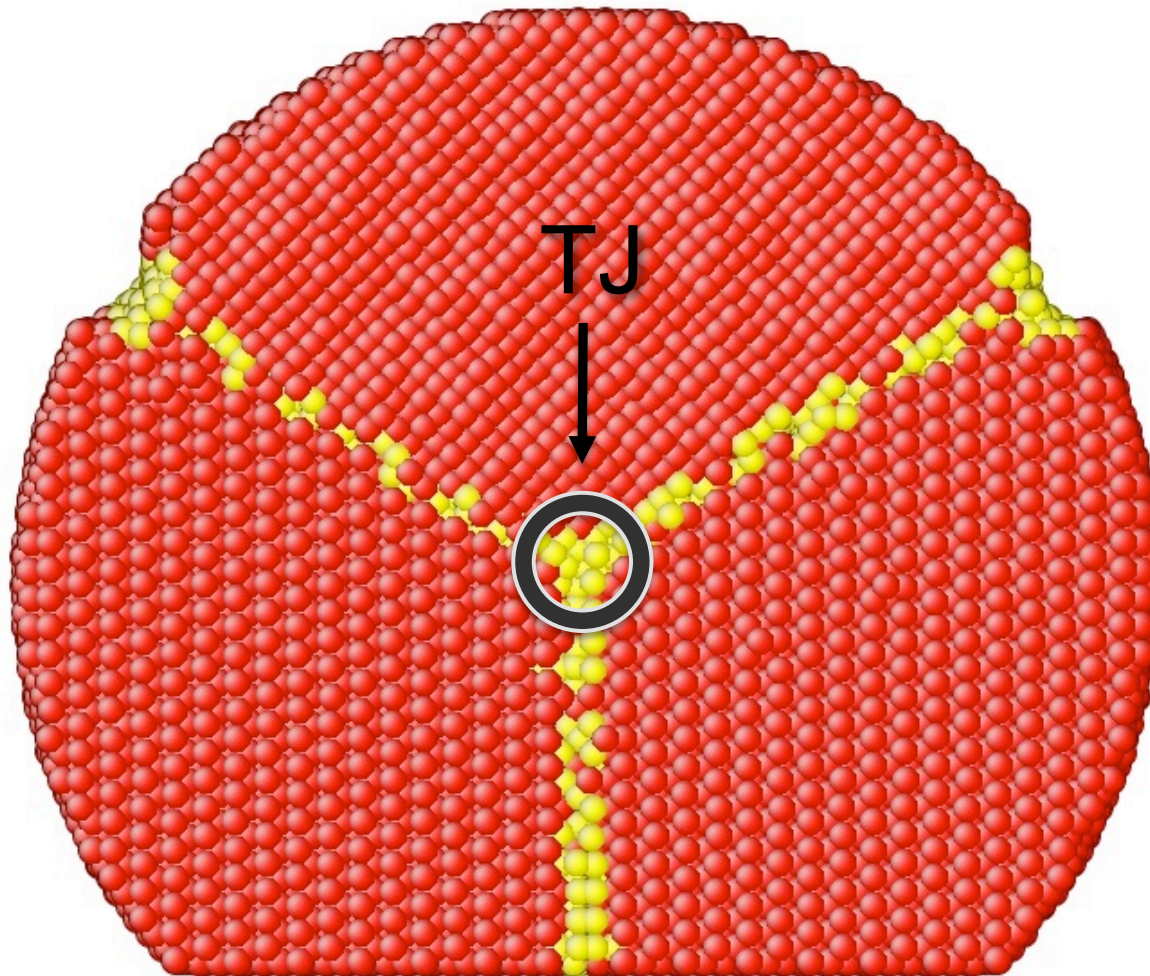
- Einstein relation  $\langle X^2 \rangle = 2Dt$ 
  - $X$  direction parallel to the TJ
  - $D$  diffusion coefficient
  - $t$  time
- Atomic displacements computed for particular regions



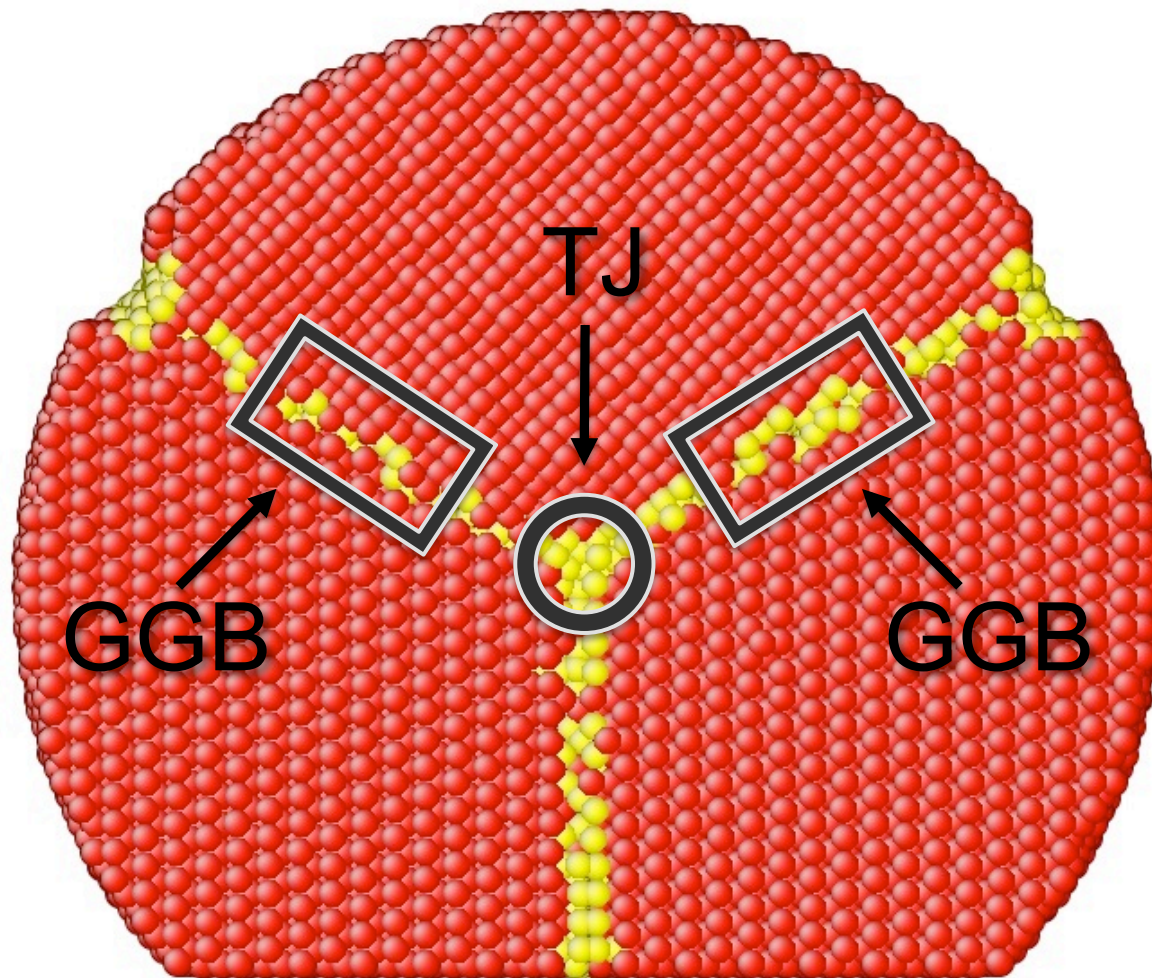
# Regions for TJ and GBs diffusion calculation



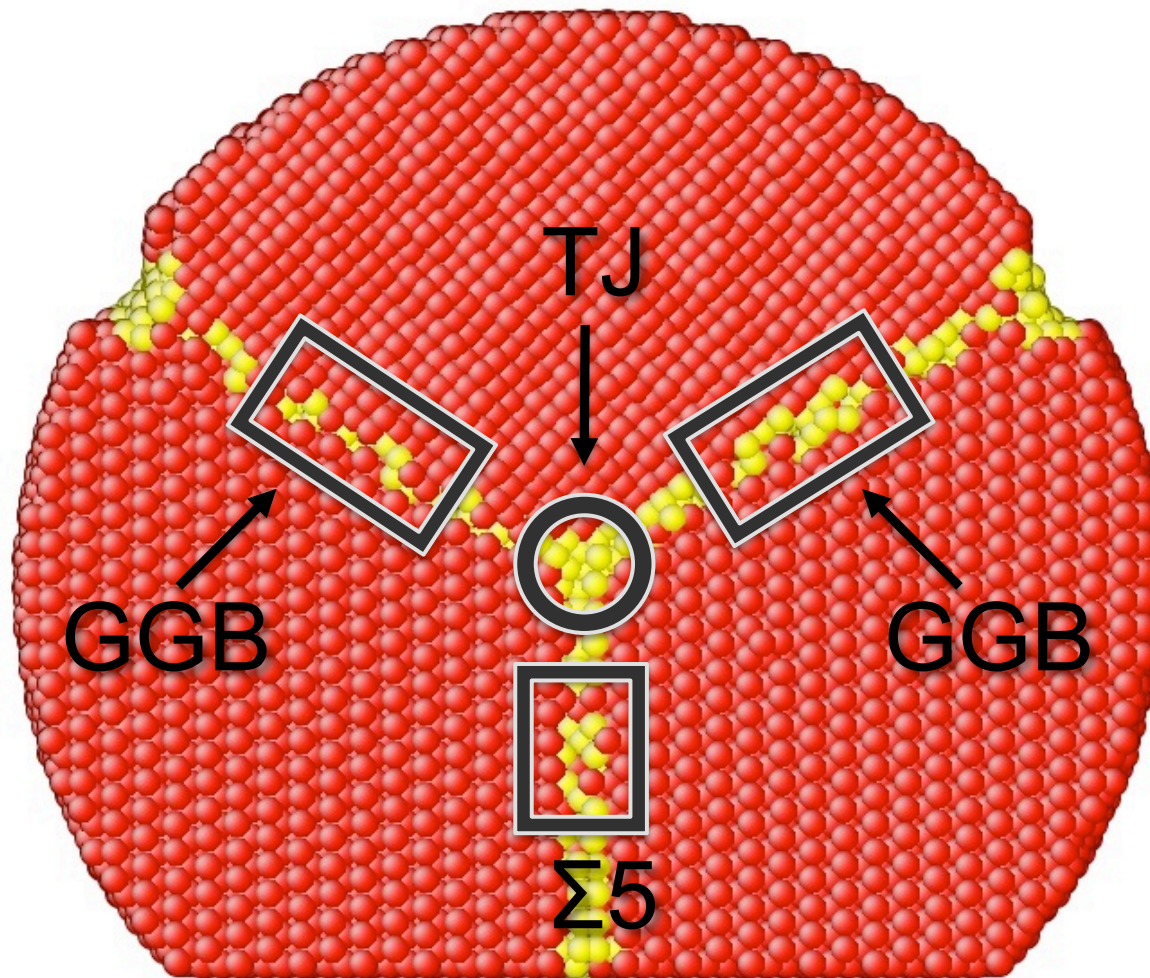
# Regions for TJ and GBs diffusion calculation



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# Regions for TJ and GB diffusion calculation

- TJ

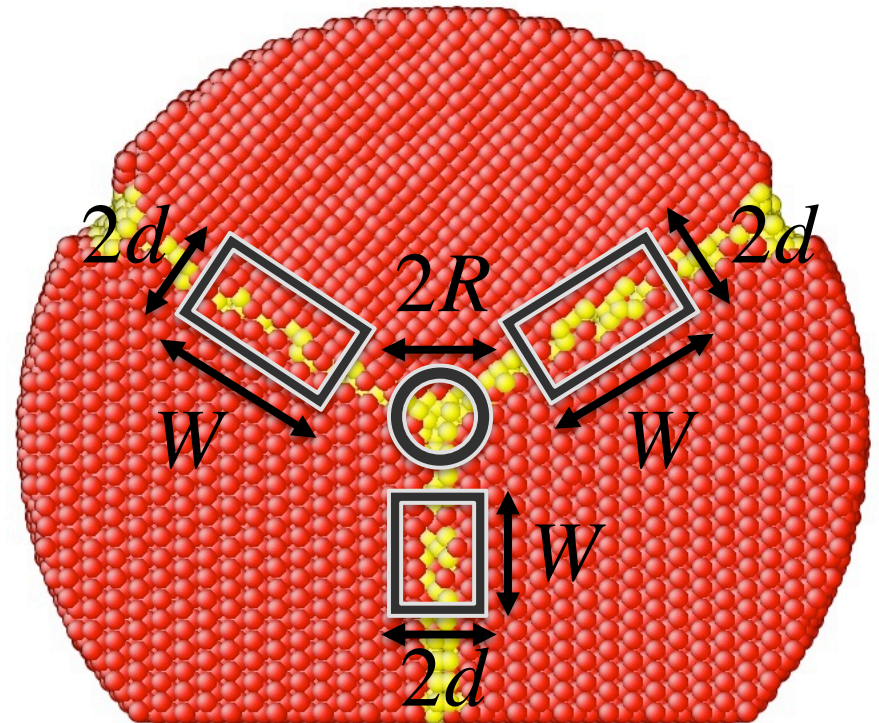
$$2R = 10 \text{ \AA} \quad L = 66 \text{ \AA}$$

- $\Sigma 5$

$$2d = 10 \text{ \AA} \quad W = 25 \text{ \AA} \quad L = 66 \text{ \AA}$$

- GGBs

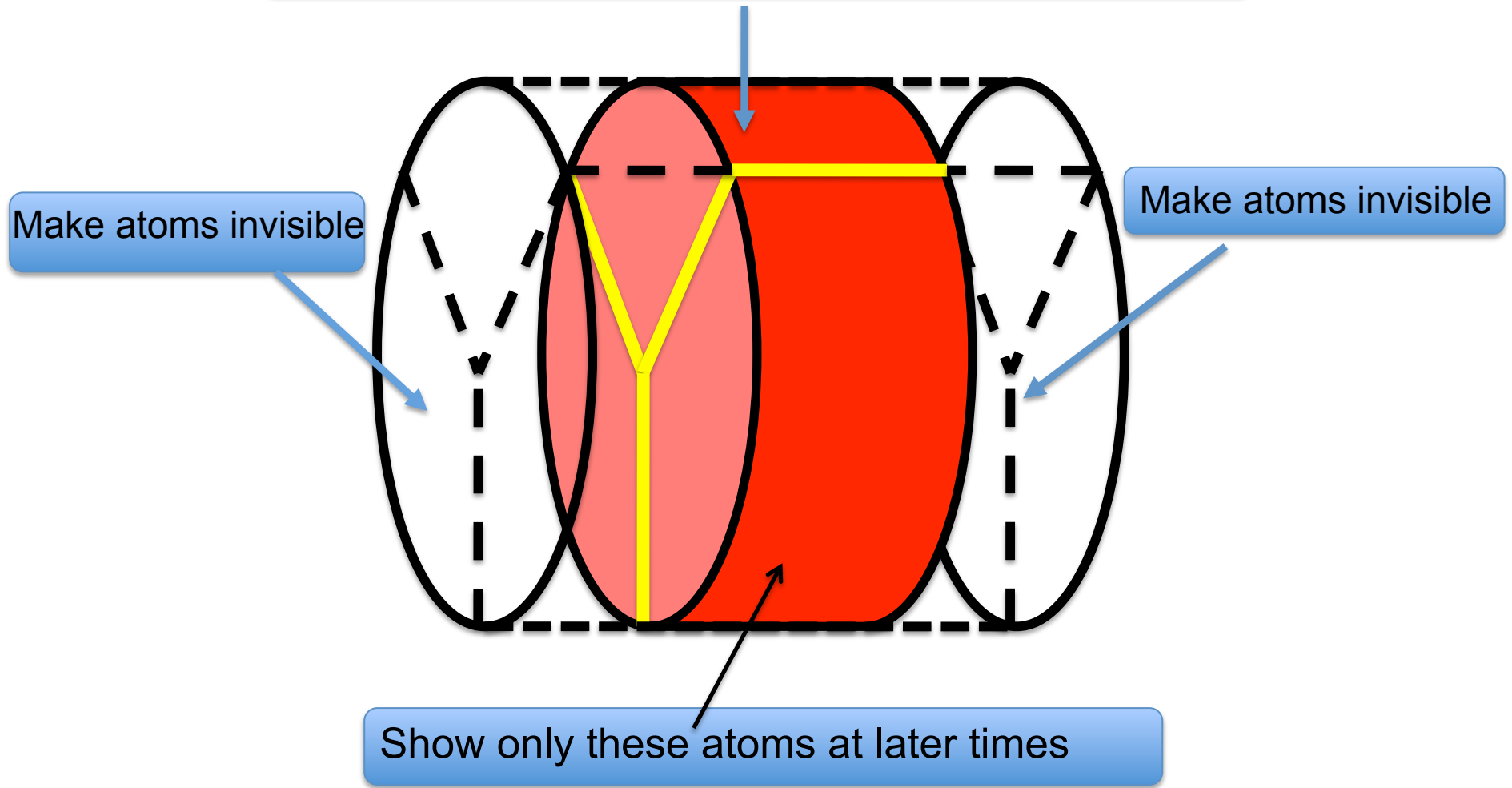
$$2d = 10 \text{ \AA} \quad W = 30 \text{ \AA} \quad L = 66 \text{ \AA}$$



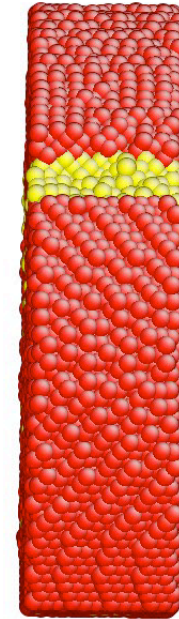
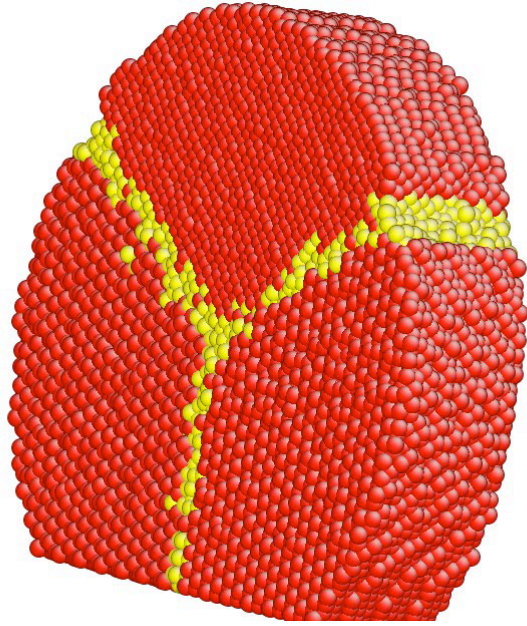


# Animation

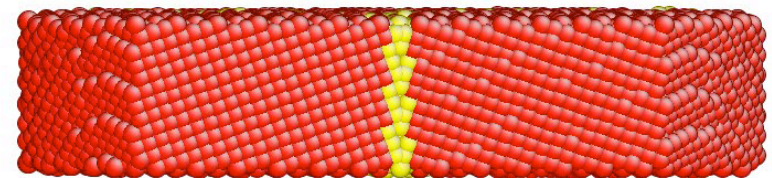
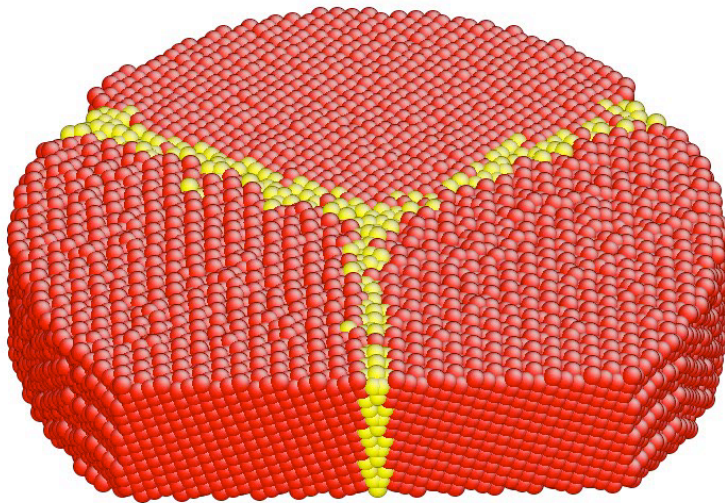
Select a slab of atoms in the middle of the block



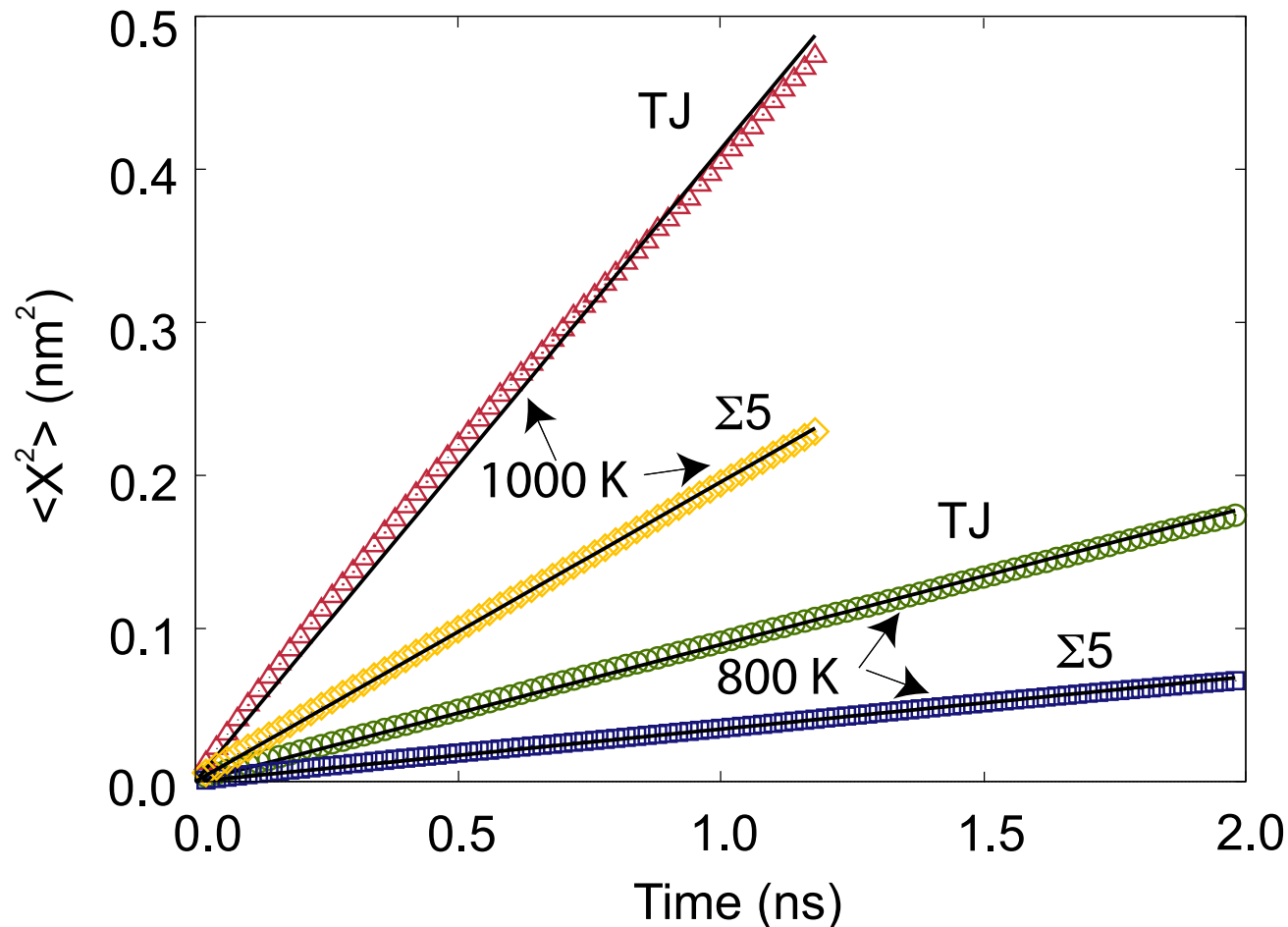
$T=1100\text{ K}$



$T_J$

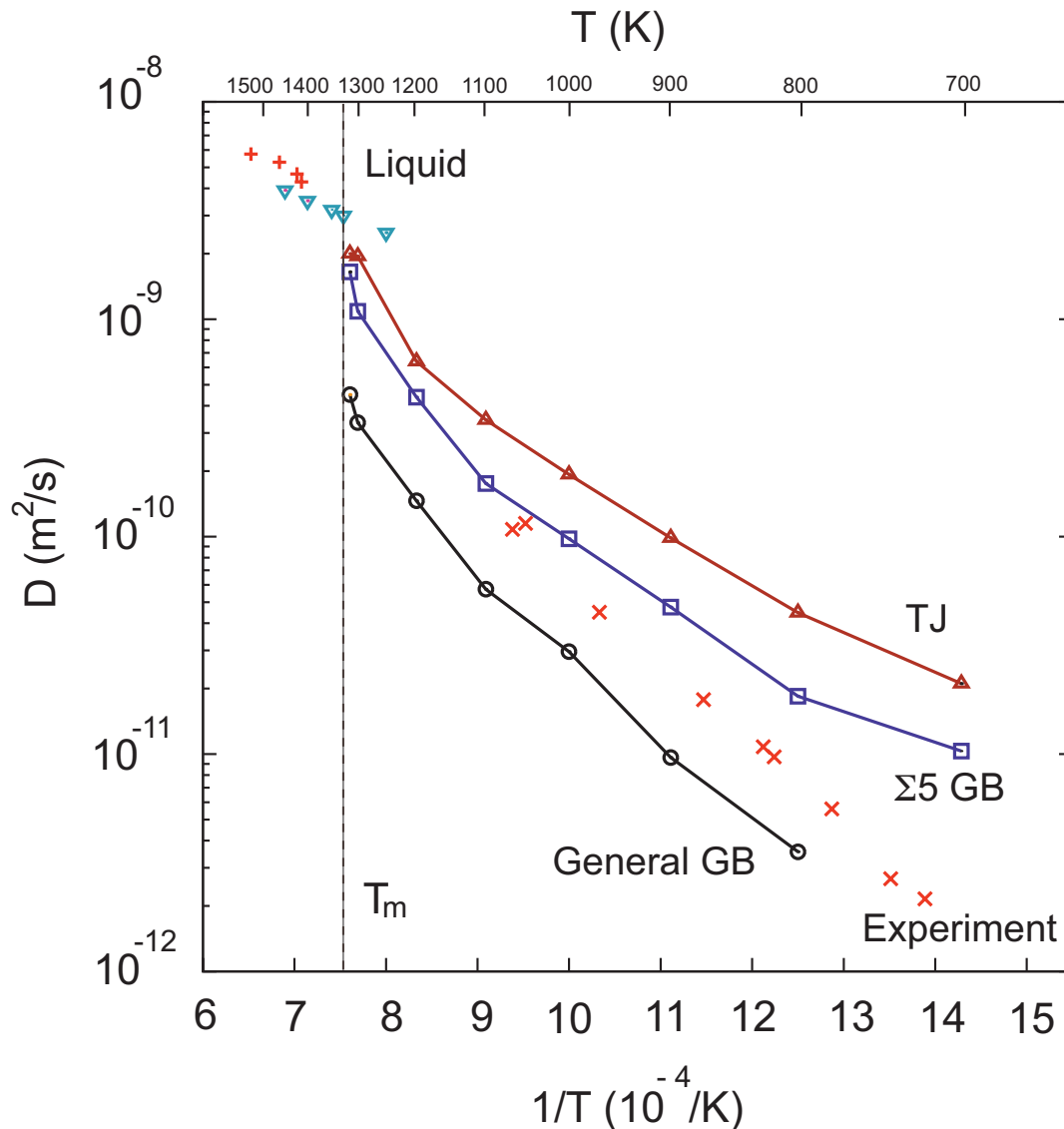


# Typical RMS displacement versus time plots for TJ and GBs



- Displacements of atoms follow the Einstein relation

# Arrhenius diagram of diffusion



- At all T  $D_{TJ} > D_{\Sigma 5} > D_{GGB}$

- Premelting at high T

- Experimental data:  
polycrystalline copper  
(average over different GBs and  
TJs)

T. Surholt et al (1997)

liquid copper

J. Henderson et al (1961)

# Conclusions

- Stable TJs with controlled crystallographic orientations of the grains can be created in computer simulations
- Self-diffusion in the TJ and adjacent GBs was computed over a range of temperatures
- TJ diffusivity is higher than GB diffusivity at all simulated temperatures
- TJ diffusivity **only twice** large than the diffusivity in  $\Sigma 5$  GB
- Contribution of TJs to diffusivity in nano-crystalline materials is probably overestimated

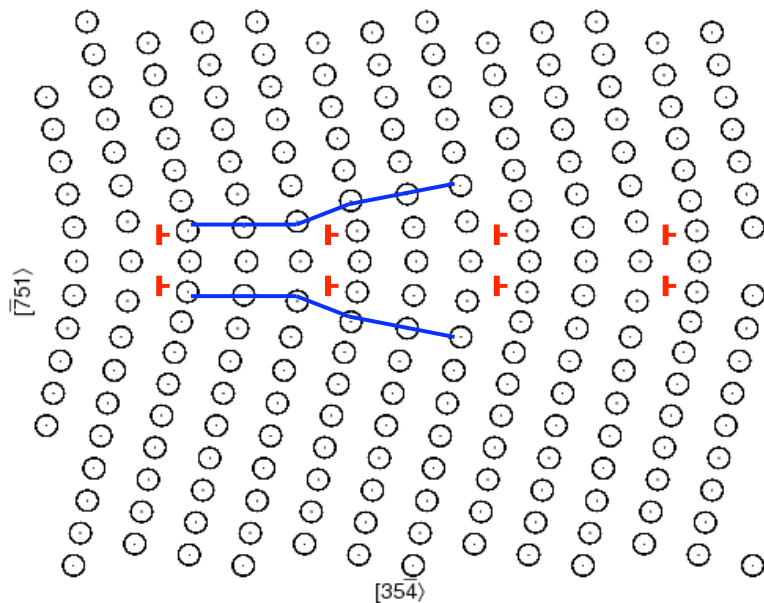
# Diffusion along a dislocation grain boundary in Al

Ganga Pun

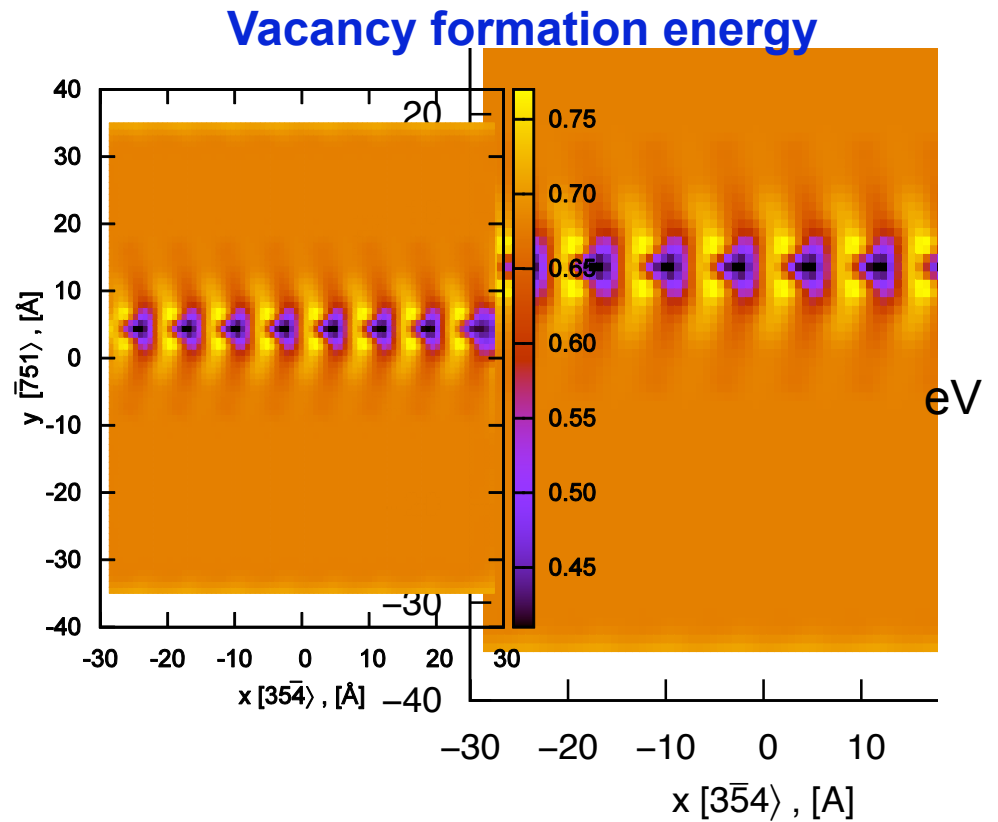
# Symmetrical tilt grain boundary in Al

$\Sigma 75$  (751) [112], 23.07°

$N_{\text{free}} = 12,000$  atoms  
 $L_x = 58.3 \text{ \AA}$   
 $L_y = 107.4 \text{ \AA}$   
 $L_z = 50.4 \text{ \AA}$



XY view



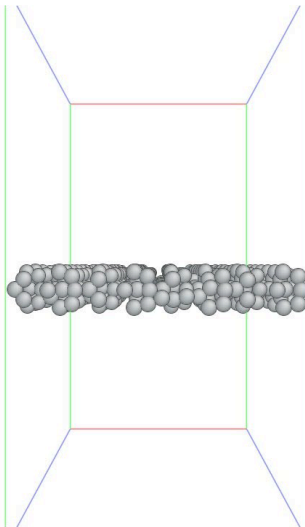
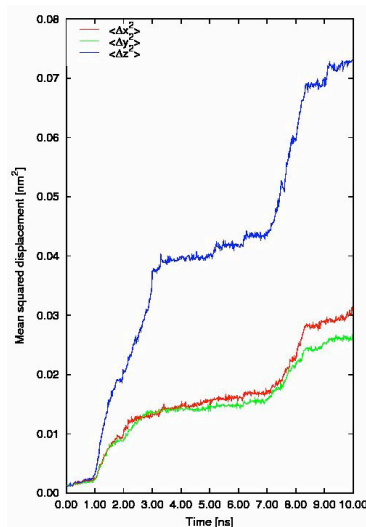
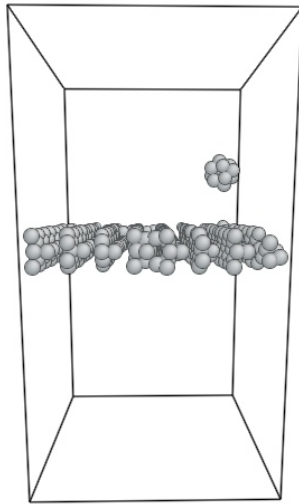
# Three types of simulations

- Pre-existing vacancies
- Pre-existing interstitials
- No pre-existing defect (**intrinsic** mechanism)

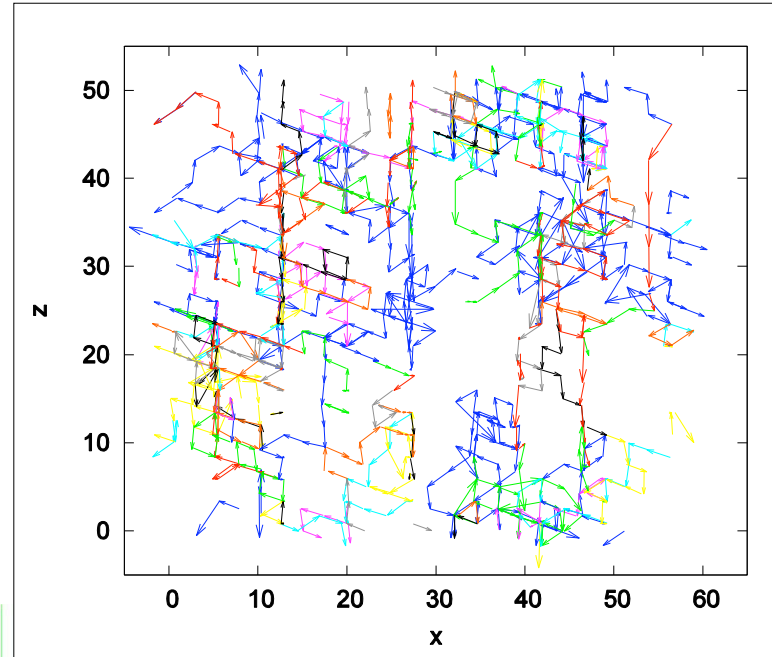


# Diffusion along edge dislocations in $\Sigma 75$ (751) [112]

## 800K (w/vacancy)



## Atomic jumps: 900K (with 8 vacancies)



## 850K (intrinsic)

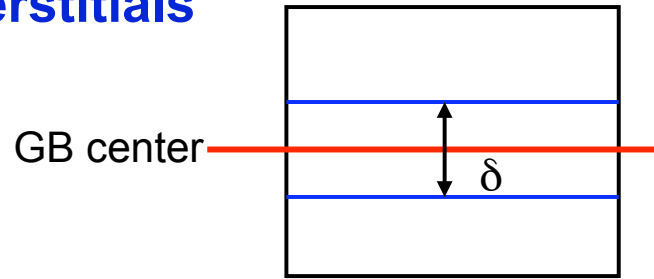
### Intrinsic mechanism:

- Frenkel pair forms in the GB
- Vacancy escapes to the lattice
- Interstitial mediates fast diffusion
- The Frenkel pair recombines

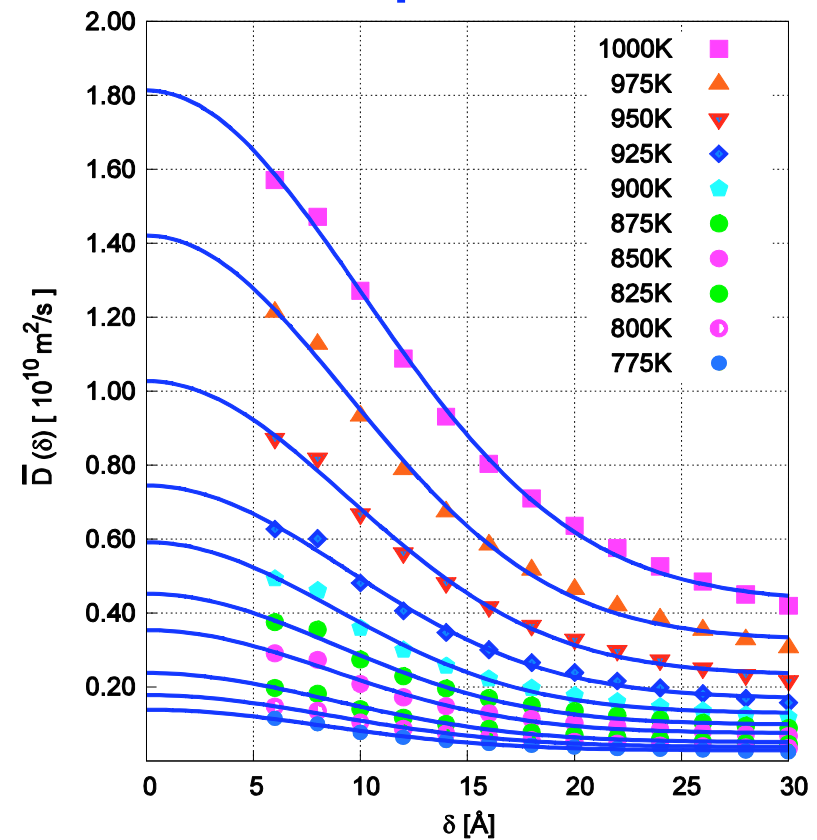
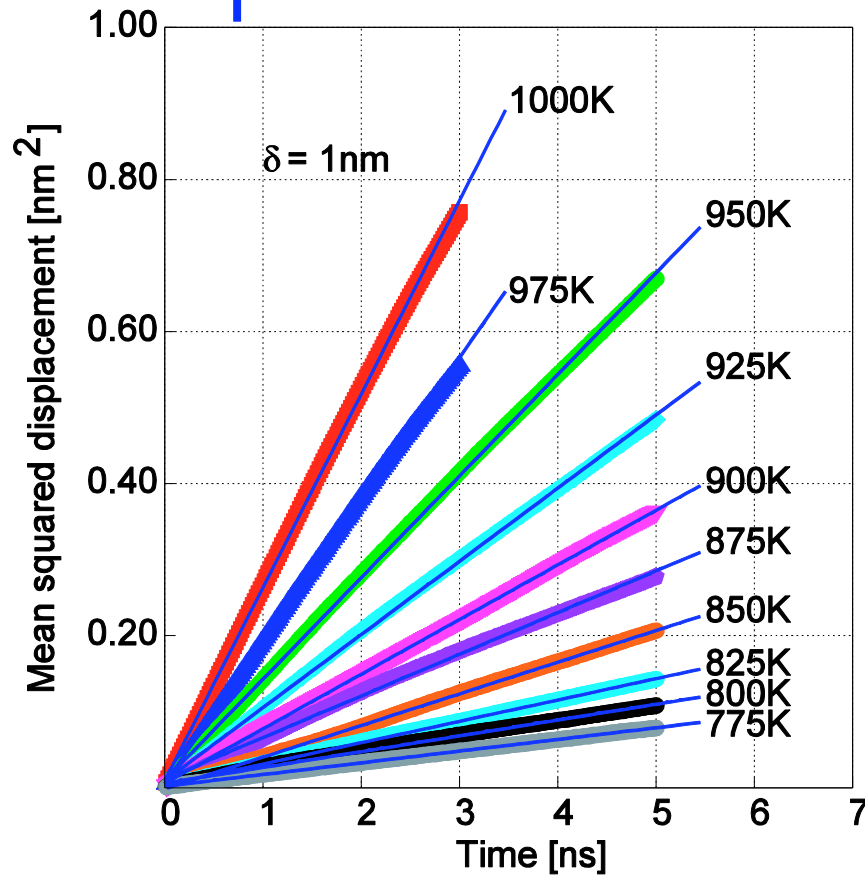
# Calculation of “raw” diffusivity

## Case with 8 interstitials

$$\bar{D}^{raw}(\delta) = \frac{\langle z^2 \rangle}{2\tau}$$



$$\bar{D}^{raw}(\delta) = A \exp\left(-\frac{\delta^2}{\delta_0^2}\right) + B$$



# Calculation of actual diffusivity

$$\bar{D}^{raw}(R) = \frac{\langle z^2 \rangle}{2\tau}$$



$$\bar{D}^{raw}(\delta) = A \exp\left(-\frac{\delta^2}{\delta_0^2}\right) + B$$



$$\bar{D}_{GB}^{raw}(\delta = \delta_0) \approx A/e$$

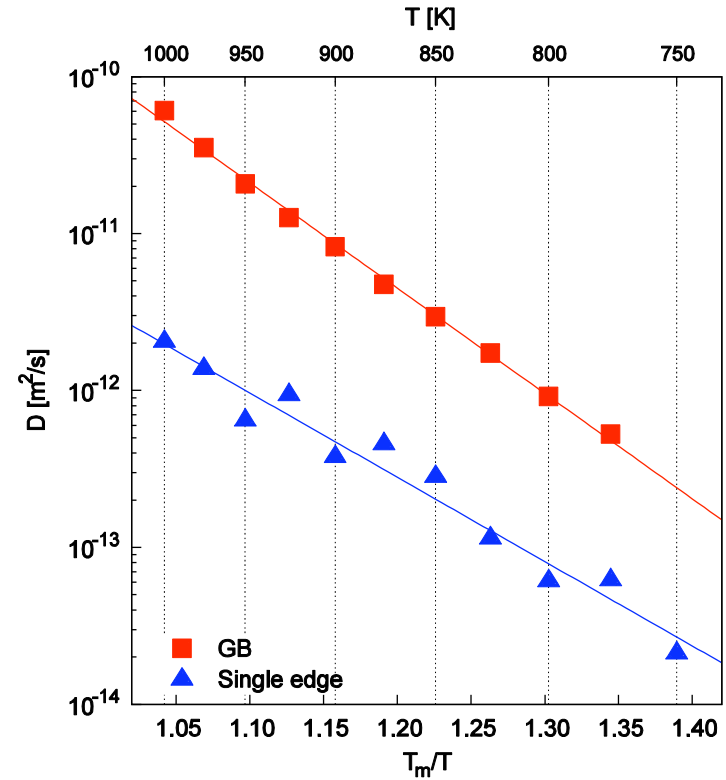
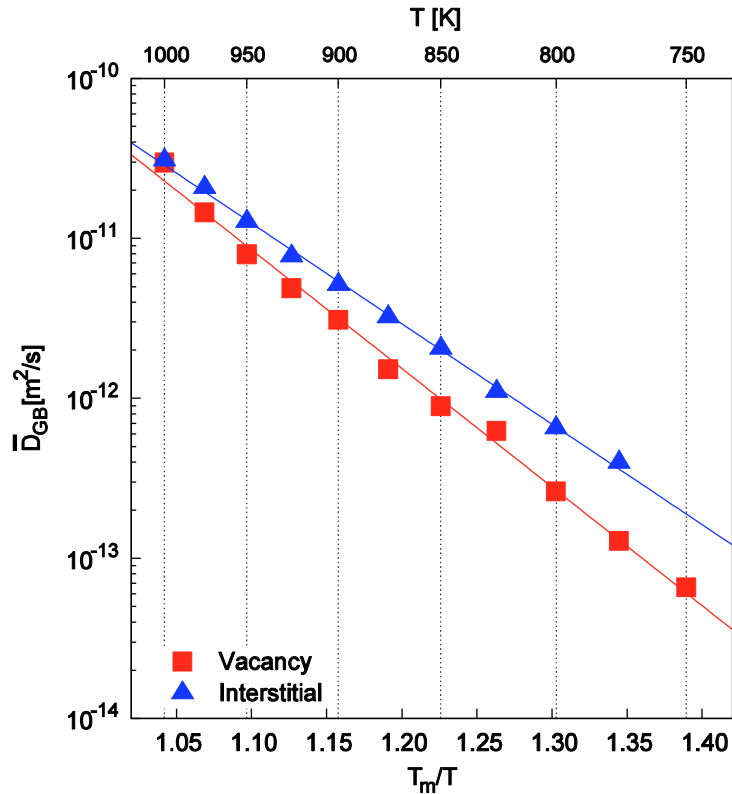


$$\bar{D}_{GB} = \bar{D}_{GB}^I + N^{eq}_v (\bar{D}_{GB}^{raw} - \bar{D}_{GB}^I) / N_v$$

$\bar{D}_{GB}^I$  : average intrinsic diffusivity

$$N^{eq}_v = \sum_{i=1}^N \exp\left(-\frac{E_{vi}}{kT}\right)$$

# Arrhenius plots



$\delta_0 = 0.35-0.40$  nm

	Vacancy	Interstitial
$E_a$ (eV)	$1.53 \pm 0.03$	$1.30 \pm 0.02$
$\text{Log}(D_0)$ (m <sup>2</sup> /s)	$-2.91 \pm 0.20$	$-4.00 \pm 0.11$

	Grain boundary	Single edge
$E_a$ (eV)	$1.39 \pm 0.02$	$1.11 \pm 0.07$
$\text{Log}(D_0)$ (m <sup>2</sup> /s)	$-3.29 \pm 0.15$	$-6.10 \pm 0.43$

# Conclusions

- $\Sigma 75$  (751) [112], 23.07° symmetrical tilt grain boundary:
  - Interstitial diffusion dominates over vacancies
  - High diffusivity compared to single edge core
- Working on other low angle symmetrical tilt and twist grain boundaries

Last slide of my presentation at this workshop on February 7, 2006:

## Conclusions

- Once very fashionable, the area of GB diffusion is not **hot** anymore. It is **not** considered to be **cool** enough. It cannot compete with carbon nanotubes and quantum dots
- It is only the Herzig group in Muenster that keeps GB diffusion measurements alive
- Development of GB diffusion theory stopped (June 1, 2005)
- Posterity will not forgive us

**Never say never!**

