Rotational Ligand Dynamics in Mn[N(CN)₂]₂.pyrazine Craig Brown, Nick Butch, Wei Zhou



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Outline

- Mn[N(CN)₂]₂.pyrazine
 - Physical Properties
 - Review data from other techniques
 - Compare behaviour with related compounds
- Categories of experiments performed on DCS
 - Extra slides
 - Quasi-elastic scattering
 - What is it?
 - What it means

Pyrazine



Interactions



Structure and dynamics

- Deuterated sample for coherent Bragg diffraction to obtain structure as a function of temperature
- Protonated to observe both single particle motion (quasielastic) and to weigh the inelastic scattering spectrum in favor of hydrogen (vibrations)
 - Deuteration can help to *assign* particular vibrational modes and provide a '*correction*' to the quasielastic data for the paramagnetic scattering of manganese and coherent quasielastic scattering.

Magnetic Structure





One of the interpenetrating lattices shown. *a* is up, *b* across, *c* into page
Magnetic cell is (½, 0, ½) superstructure
Exchange along Mn-pyz-Mn chain 40x

J. L. Manson *et. al*

- J. Am. Chem. Soc. 2000
- J. Magn. Mag. Mats. 2003

Mn[N(CN)₂]₂.pyrazine







Mn[N(CN)₂]₂.pyrazine

Weight Loss (%)

- **1.3 K** 3-D antiferromagnetic order below ~2.5 K
 - Magnetic moments aligned along a (4.2 $\mu_{\rm B}$)
 - Monoclinic lattice (*a*=7.3 Å, *b*=16.7 Å, *c*=8.8 Å)

~200 K

- Phase transition to orthorhombic structure
- Large Debye-Waller factor on dicyanamide ligand
- Diffuse scattering

408 K

- Phase transition
 - Large Debye-Waller factors on pyrazine



~435 K• Decomposes and loses pyrazine.

As a function of Temperature



The other compounds



AIMS

- **Experience** Practical QENS
 - sample choice

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- geometry consideration
- Learn something about the instrument
 - Wavelength / Resolution / Intensity
- Data Reduction
 - Data Analysis and Interpretation
 - instrument resolution function and fitting
 - extract EISF and linewidth
 - spatial and temporal information

The Measured Scattering





Quasielastic Scattering

- The intensity of the scattered neutron is broadly distributed about zero energy transfer to the sample
- Lineshape is often Lorentzian-like
- Arises from atomic motion that is
 - Diffusive
 - Reorientational
- The instrumental resolution determines the timescales observable
- The *Q*-range determines the spatial properties that are observable
- (The complexity of the motion(s) can make interpretation difficult)

Types of Experiments

- Translational and rotational diffusion processes, where scattering experiments provide information about time scales, length scales and geometrical constraints; the ability to access a wide range of wave vector transfers, with good energy resolution, is key to the success of such investigations
- Low energy vibrational and magnetic excitations and densities of states
- Tunneling phenomena
- **Chemistry** --- e.g. clathrates, molecular crystals, fullerenes, MOFs
- **Polymers** --- bound polymers, glass phenomenon, confinement effects
- **Biological systems** --- protein folding, protein preservation, water dynamics in membranes
- **Physics** adsorbate dynamics in mesoporous systems (zeolites and clays) and in confined geometries, metal-hydrogen systems, glasses, magnetic systems
- **Materials** --- negative thermal expansion materials, low conductivity materials, hydration of cement, carbon nanotubes, proton conductors, metal hydrides, hydrogen diffusion, CH4 dynamics....



Data courtesy of M. Lumsden, ORNL

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Magnetism



x=0.100003



Para has a nuclear spin I=0. This constrains J to be even.

Ortho has a nuclear spin I=1. This constrains J to be odd.

Transition between ortho and para species can occur through flipping the nuclear spin.











Monitor hydrogen diffusion over isotherm







Extract diffusion constants, residency times, activation energies 2 types of hydrogen facilitate modeling, degrees-of-freedom ...



Surface diffusion reduced and shorter hops with loading At 77 K hydrogen behaves like it is at 35 K on carbons

Some notes on data meaning

The Measured Scattering

 $S(Q,\omega) = S(Q,\omega)^{Reorient} \otimes S(Q,\omega)^{Lattice} \otimes S(Q,\omega)^{VIB} \otimes R(Q,\omega)$

The reorientational and/or Lattice parts.

The Lattice part has little effect in the QE region- a flat background (see Bée, pp. 66)

Debye-WallerInstrumentalfactorresolutionFar away from the
QE regionfunction

Quasielastic Neutron Scattering Principles and applications in Solid State Chemistry, Biology and Materials Science M. Bee (Adam Hilger 1988)

Quasielastic Scattering

 $G_s(r,t)$ is the probability that a particle be at r at time t, given that it was at the origin at time t=0(self-pair correlation function)

 $I_{inc}(Q,t)$ is the space Fourier transform of $G_s(r,t)$

(incoherent intermediate scattering function)

$$I_{inc}(Q,t) = \left\langle e^{iQ.r(t)} e^{-iQ.r(0)} \right\rangle$$

 $S_{inc}(Q, \omega)$ is the time Fourier transform of $I_s(Q, t)$ (incoherent scattering law)

$$S_{inc}(\vec{Q},\omega) = \frac{1}{2\pi} \int I_{inc}(Q,t) e^{-i\omega t} dt$$

Quasielastic Scattering $r_1 r_2$ Jump model between two equivalent sites $\frac{\partial}{\partial t} p(\mathbf{r}_1, t) = -\frac{1}{\tau} p(\mathbf{r}_1, t) + \frac{1}{\tau} p(\mathbf{r}_2, t)$ $\frac{\partial}{\partial t} p(\mathbf{r}_2, t) = \frac{1}{\tau} p(\mathbf{r}_1, t) - \frac{1}{\tau} p(\mathbf{r}_2, t)$ $\frac{\partial}{\partial t} \left[p(\mathbf{r}_1, t) + p(\mathbf{r}_2, t) \right] = 0$ $p(r_1,t) + p(r_2,t) = 1$ $p(\mathbf{r}_1, t; \mathbf{r}_1, 0) = \frac{1}{2} [1 + e^{-2t/\tau}] \qquad p(\mathbf{r}_2, t; \mathbf{r}_1, 0) = \frac{1}{2} [1 - e^{-2t/\tau}]$ $p(r_2, t; r_2, 0) = \frac{1}{2} [1 + e^{-2t/\tau}]$ $p(r_1, t; r_2, 0) = \frac{1}{2} [1 - e^{-\frac{2t}{\tau}}]$ $I(Q,t) = [p(r_1,t;r_1,0) + p(r_2,t;r_1,0)e^{iQ(r_2-r_1)}]p(r_1,0)$ +[$p(r_1,t;r_2,0)e^{iQ(r_1-r_2)} + p(r_2,t;r_2,0)$] $p(r_2,0)$ $I(Q,t) = \frac{1}{2} [1 + \cos Q \cdot (r_2 - r_1)] + \frac{1}{2} [1 - \cos Q \cdot (r_2 - r_1)] e^{-\frac{2t}{\tau}}$

Quasielastic Scattering r_{\bullet} Jump model between two equivalent sites $S(Q,\omega) = \frac{1}{2} [1 + \cos Q \cdot (r_2 - r_1)] \delta(\omega) + \frac{1}{2} [1 - \cos Q \cdot (r_2 - r_1)] \frac{1}{\pi} \frac{2\tau}{4 + \omega^2 \tau^2}$ **Powder (spherical) average** $S(Q,\omega) = \frac{1}{2} \left[1 + \frac{\sin Q.(r_2 - r_1)}{Q.(r_2 - r_1)}\right] \delta(\omega) + \frac{1}{2} \left[1 - \frac{\sin Q.(r_2 - r_1)}{Q.(r_2 - r_1)}\right] \frac{1}{\pi} \frac{2\tau}{4 + \omega^2 \tau^2}$ $S(Q,\omega) = A_0 \delta(\omega) + A_1 \frac{1}{\pi} \frac{2\tau}{4 + \omega^2 \tau^2}$ EISF!!!!!!!! **OISF!!!!!!!**



Quasielastic Scattering

Jump model between two equivalent sites



Quasielastic Scattering Jumps between three equivalent sites $S(Q,\omega) = A_0 \delta(\omega) + (1 - A_0) \frac{1}{\pi} \frac{3\tau}{9 + \omega^2 \tau^2} \qquad A_0 = \frac{1}{3} \left[1 + 2j_0 (Qr\sqrt{3}) \right]$ 1.00.80.6 $A_o(Q)$ 0.4

10

 \mathbf{qr}

15

20

0.2

0.0

0

5



TOF spectroscopy, in practice

(1) The neutron guide



(2) The choppers

(3) The sample area

(4) The flight chamber and the detectors