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Event identification in ³He proportional counters using risetime discrimination

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ABSTRACT

We present a straightforward method for particle identification and background rejection in ³He proportional counters for use in neutron detection. By measuring the risetime and pulse height of the preamplifier signals, one may define a region in the risetime versus pulse height space where the events are predominately from neutron interactions. For six proportional counters surveyed in a low-background environment, we demonstrate the ability to reject alpha-particle events with an efficiency of 99%. By applying the same method, we also show an effective rejection of microdischarge noise events that, when passed through a shaping amplifier, are indistinguishable from physical events in the counters. The primary application of this method is in measurements where the signal-to-background for counting neutrons is very low, such as in underground laboratories.

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1. Introduction

³He proportional counters are used as neutron detectors in a variety of fields in physics such as neutron scattering and nuclear and particle physics, particularly those experiments that operate underground in a low-background environment. They also have a commercial application as neutron detectors in health physics, medical physics, and nuclear safeguards. ³He has the advantages of a large thermal neutron capture cross-section, easily detected charged reaction products, and low gamma-ray sensitivity. Although the basic properties of ³He proportional counters have been known for decades [1], more recent applications in counting low rates of neutrons have spurred work in better understanding their backgrounds. Methods for improving background rejection in the counters [2,3] can be a significant benefit for those experiments or applications where the neutron signal-to-background ratio is particularly poor.

Specifically, alpha-particle decays are a pernicious background for ³He detectors used in low counting rate applications. The rate of alpha-particle decays originating from the walls of the counter which can be comparable to, or even exceed, the rate of neutron captures. Alpha particles arise primarily from daughters in the uranium and thorium decay chains. The amount of energy deposited in the detector can be from essentially zero to several MeV, spanning the region of the neutron capture products. Thus, energy discrimination alone is not sufficient to reject background events from alpha particles. The method of risetime discrimination provides a straightforward method to reject a large fraction of these events.

For experiments operating in deep underground environments, cosmic-ray induced neutron events can be a particularly difficult background because one cannot turn off the source of the background [4,5]. The experiments can only reduce the ambient neutron background through shielding and precisely quantifying the remaining neutron flux [6]. Solar neutrino experiments have used ³He proportional counters as triggers for neutral current interactions and monitors of muon-induced neutron backgrounds [7]. Dark matter experiments are particularly susceptible to neutron backgrounds, due to neutron recoil events that can mimic WIMP dark matter interactions. To this end, many experiments use ³He proportional counters to monitor and characterize the neutron flux at their experimental site [8,9].

In one notable example, the Sudbury Neutrino Observatory (SNO) collaboration took considerable effort in constructing ultralow background ³He proportional counters containing low concentration of alpha emitters. That effort involved manufacturing proportional counters from materials with low concentrations of uranium and thorium, as well as working in a radon-free

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environment to prevent surface contamination [10,11]. During the Phase III run, 3.0 n/day were observed compared to 14.4 alphas/ day in the signal region of their helium proportional counters [12]. For experiments that rely upon neutron counters as their triggers, these alpha-particle interactions in the proportional counter can be a leading source of background. Therefore, experiments such as SNO have distinguished alpha-particle background from neutron signals either through risetime analysis [13,14] or chi-squared fitting of the entire pulse waveforms [15].

More generally, fast neutron spectroscopy can be done directly using a ³He proportional counter [16,17], and it is also routinely done using Bonner spheres, where fast neutrons are moderated in a hydrogenous material and a fraction of those thermalized neutrons may be captured by a central ³He proportional counter [18,19]. Another technique for fast neutron detection using ³He proportional counters is capture-gated spectroscopy [20]. The fast neutrons recoil from protons in an active hydrogenous moderator; the moderated neutron diffuses in the medium and may be subsequently captured on ³He. The time-delayed capture signifies that the preceding interaction was a fully thermalized neutron. This allows for an accurate measurement of the incident particles energy, as well as a highly efficient method for background reduction.

Improving the particle identification and rejection of nonneutron events in ³He proportional counters have clear relevance to these applications. More specifically, this collaboration is constructing a fast neutron spectrometer [21] that employs a large number of ³He counters, and hence it is critical to understand all the background contributions in each counter that may be part of the spectrometer. Approximately 75 ³He proportional counters of the same model were surveyed as part of this study. Six proportional counters were surveyed at the Kimballton Underground Research Facility (KURF) in Ripplemead, VA. An underground lab with approximately 1450 mew overburden, KURF houses a variety of low background experiments [22–25].

In this paper, we discuss how a combination of risetime and energy analysis can significantly improve the identification of neutron capture events in a ³He proportional counter. Section 2 briefly reviews the operation of ³He proportional counters. In Section 3 we describe the experimental setup and analysis method for testing ³He proportional counters. Section 4 gives the approach for identifying the origin of various waveforms that may occur in a ³He proportional counters. We apply that method to quantify the rejection of non-neutron events for a specific counter in Section 5. Section 6 describes an auxiliary measurement to quantify alphaparticle rates from the proportional counter body, and Section 7 summarizes the results.

2. ³He proportional counters

2.1. Principles of operation

³He provides an excellent medium for detecting thermal neutrons due to the high capture cross-section (5330×10^{-24} cm²) and a final state that consists of charged particles, as opposed to gamma rays. A neutron is captured by a ³He nucleus, resulting in a proton and a triton, which share 764 keV of kinetic energy

$$n+{}^{3}\text{He} \rightarrow p+t+764 \text{ keV}. \tag{1}$$

These heavy charged particles ionize the gas in the proportional counter, creating electron-ion pairs that are proportional to the energy deposited. For cylindrical proportional counters, the electrons drift towards a central anode wire, and when the electric field gradient reaches a critical level, begin to avalanche. Thus, the detected current is created from the drifted charge.

Because the neutron capture on ³He has a two body final state, the charged particles will be emitted in opposite directions. If the ionization tracks are parallel to the central anode wire, all of the charge will be collected at approximately the same time, leading to a short risetime of the detected signal. However, if the particles are emitted perpendicular to the anode wire, one particle will be moving towards the wire while the other is moving away. The collected charge will be spread out in time due to the radial variation caused by the track geometry [26].

Typically, a ³He proportional counter is biased through a preamplifier, and the resulting signals are sent through a shaping amplifier to be analyzed with a multi-channel analyzer or peaksensing analog to digital converter. However, shaping the preamplifier signal smoothes out the slight deviations in signal shape caused by the track geometry. With high-speed waveform digitizers, it is possible to study the shape of preamplifier signals directly and apply digital signal processing to each event to identify the type of incident radiation.

2.2. Specific energy loss and particle range

The specific energy loss of charged particles with charge Z, mass m, and kinetic energy E_{kin} , is approximately proportional to

$$-\frac{dE}{dx} \propto \frac{Z^2 m}{E_{kin}}.$$
 (2)

Thus, the track length for a given E_{kin} decreases with increasing mass and charge of the projectile. For neutron capture on ³He, the proton is emitted with 573 keV and the triton with 191 keV. The stopping power and particle ranges can be accurately determined using the TRIM software package [27]. At these energies, the proton has a range of 5.76 mm and the triton has a range of 1.6 mm for the gas mixture of our proportional counters, described in Section 3. Because the proton and triton are emitted back-to-back, this yields a total track length of 7.36 mm for a neutron capture. An alpha particle with the same energy as the neutron capture, 764 keV, will have a track 2.54 mm long.

Beta emitters, electrons from γ -interactions, and cosmic rays leave long tracks and deposit little energy resulting in small signals with a wider range of risetimes. Alpha particles with much higher specific energy loss leave shorter tracks and deposit more of their energy, which yield large signals with a relatively fast risetime. These features can be exploited in analysis to help identify the original radiation. Details about the energy loss of charged particles and the resulting pulse shapes in proportional counters can be found elsewhere [26,28–31].

3. Experimental setup

The experimental setup is shown in Fig. 1. We used 2.54 cm outer diameter, 46.3 cm active length, aluminum-bodied, cylindrical ³He proportional counters manufactured by GE-Reuter



Fig. 1. A schematic showing the setup for a single channel of our measurements. Particles in the ³He proportional counters generate current signals that go to the preamplifier. The waveform digitizer captures the resulting long-tailed signals and records them for off-line analysis.



Fig. 2. Typical energy spectrum of neutron capture events in one of the ³He proportional counters used in this study. The dominant feature is the full-deposition peak at 764 keV. The two edges in the spectrum near 200 keV and 600 keV are related to partial energy deposition in the gas when either the proton or the triton interacts in the wall.

Stokes, model number RS-P4-0819-103.² The ³He partial pressure in the counters was 404 kPa (4 atm) with a buffer gas consisting of 111 kPa (1.1 atm) of krypton. Helium does not have a high stopping power for heavy charged particles. Therefore, manufacturers tend to add in a buffer gas with higher stopping power to ensure full energy deposition from the capture products. The krypton increases the stopping power of the gas for charged particles but has little effect on neutron capture.

The proportional counters were biased through a charge sensitive preamplifier, and the traces were recorded for offline analysis by a 125 MSamples/s 12-bit GaGe waveform digitizer or a 250 MSamples/s 12-bit CAEN waveform digitizer. There was no discernible effect on the analyzed data using the faster digitization rate. Data were taken at different voltage levels to test the variability of the counters' response. For these specific counters, 2000 V was a reasonable operating voltage. An exploration of the effect of different counter bias voltage indicated that there was little variation of the risetime distribution.

Fig. 2 shows a typical energy spectrum measured by one of the proportional counters used in this work. Note that although the neutron capture reaction is monoenergetic, there are still features in the energy spectrum related to when one of the particles interacts in the counter wall, leading to reduced energy deposition. This well-known wall effect is clearly seen in the two edges around 200 keV and 600 keV in Fig. 2. Events below 200 keV are largely betas liberated from the counter body by gammas from the ²⁵²Cf source.

3.1. Analysis method

The digitized waveforms were analyzed, and both the amplitude and risetime of the signals were extracted. Fig. 3 shows sample traces corresponding to three different sources: a neutron event, an alpha-particle event, and a microdischarge from the proportional counter. Microdischarges are electronic artifacts originating from current leakage in the bias voltage breakdown applied to the counter [32–34]. To determine the pulse height of the waveform, one first removes the baseline offset by subtracting off the average of the first few hundred samples of the trace, well before the pulse onset. To remove high-frequency noise,



Fig. 3. Example of the raw preamplifier traces showing the risetimes from three different sources. Shown here, in sequence of increasing risetime (left to right), are a microdischarge, an alpha-particle trace, and a neutron trace.

a Gaussian smoothing routine was applied to the waveform, and the amplitude was determined from the resulting wave. Because the signals were sent though a charge-sensitive preamplifier, the energy of the particle is proportional to the amplitude of the recorded signal. The fall time of the signals was a property of the AC coupling and decay time of the preamplifier. The detectors were calibrated using the thermal neutron capture peak to set the energy scale.

For this work, the risetime was defined as the time difference between the signal reaching 10% and 50% of its maximum amplitude. This is the same interval used in Ref. [13]. This choice was based upon a reasonable separation of the three signal types. If the interval was too small (less than 20%), there was not adequate separation between the microdischarges and the real signals. However, the interval could not be made too large, otherwise the slow component of the risetime, due to the drift of positive ions [30], decreases the separation of alpha-particle and neutron events. For both the alpha particles and neutron traces shown in Fig. 3, this begins at about 75% of the maximum signal amplitude.

4. Event identification through risetime analysis

4.1. Physical signals from source radiation

Events in the counters arise from radiation sources or bias voltage breakdown in the counter. Data were collected under a few different conditions to characterize the ³He proportional counters' response to different sources of radiation. By displaying the risetime versus the pulse height of each signal, it is possible to identify regions of this two-dimensional space where events from different source radiation will occur. Fig. 4 shows the proportional counter response to thermal neutron, gamma, beta, alpha, and microdischarge events.

Thermal neutrons were generated from neutrons from ²⁵²Cf that were moderated by a 5 cm thick piece of polyethylene placed between the source and proportional counter. Fig. 4a illustrates the region in both risetime and pulse height where the counter detects the thermal neutron capture reaction products. The large concentration of events around 764 keV corresponds to full energy deposition of both the proton and triton. The risetime varies depending upon the orientation of the decay products in the counter with respect to the central anode wire.

To test the counter's response to beta and gamma radiation, ⁹⁰Sr and ⁶⁰Co sources were placed directly above the counter in separate measurements. Fig. 4b and c shows that these interactions are comparatively low energy that can be easily

² Certain trade names and company products are mentioned in the text or identified in an illustration in order to adequately specify the experimental procedure and equipment used. In no case does such identification imply recommendation or endorsement by the National Institute of Standards and Technology, nor does it imply that the products are necessarily the best available for the purpose.



Fig. 4. Scatter plots showing the risetime versus pulse height from different sources of radiation and microdischarges in ³He proportional counters: (a) neutron source data; (b) gamma source data; (c) beta source data; (d) alpha-particle background data; and (e) microdischarges. The solid black lines illustrate regions where the indicated events occur. Further description is found in the text of Section 4.

discriminated from neutron events. Note that there are still neutron events from the ambient background in the laboratory where the measurements were performed.

Fig. 4d illustrates why alpha particles are the most difficult background to address. Many events overlap with neutrons in both energy and risetime. For this measurement, the source of the alpha particles was the trace amount of uranium and thorium contamination in the wall of the proportional counter body. To emphasize the alpha-particle events, the counter was wrapped in boronloaded silicone rubber to absorb most of the external ambient background of thermal neutrons. As the typical alpha-particle detection rates for the counters in this work are quite low, data were taken over the course of a few days to acquire sufficient statistics. Note that with such a long run, there is still some leakage of thermal neutron events into the spectrum.

We characterized the events from about 75 3 He proportional counters for this risetime study. We note that they had a wide

range of alpha activity as well as variations in spectral shape, even though they all came from the same manufacturer. Some counters had rates consistent with or higher than the raw aluminum expectation, discussed in Section 6, while some were considerably lower. The counter with the lowest alpha-particle rate had approximately $6 \times 10^{-4} \text{ s}^{-1}$, and the highest alpha-particle rate counter was approximately $3 \times 10^{-2} \text{ s}^{-1}$. Certain counters exhibited a broad spectrum of alphas with no peaks while others displayed peaks at energies consistent with isotopes in the decay chains of uranium and thorium. During construction of the counters, a thin layer of nickel was deposited onto the inner surface of the cylinders. This was done by the manufacturer to minimize the contribution of alpha particles to the background [35]. We speculate that deviations in the thickness of this coating could cause the observed variation.

In counters with high alpha-particle rates, events were observed with pileup of a large alpha-like signal on top of a small



Fig. 5. A sample trace attributed to an alpha-particle and gamma-ray event in the proportional counter. These are possibly from a correlated decay, where a nucleus alpha decays to an excited state, which then de-excites by emitting a gamma.

gamma-like signal. These events are possibly from correlated decays, where a nucleus alpha-decays to an excited state, which subsequently emits a gamma. An example trace of this process is shown in Fig. 5. The low-amplitude gamma signal will appear first because its track is extended and closer to the anode. The alpha particles predominately originate from the counter walls, the furthest distance from the anode, and therefore arrive after the gamma event. As a result, these signals have longer overall risetimes, which could lead to misidentification as neutron-like signals with this analysis. As a check of this type of event, helium tubes were operated with plastic scintillator detectors surrounding them. There was a significant fraction of events with prompt coincidence between an alpha signal in the helium proportional counter and a gamma signal in the plastic scintillator.

4.2. Microdischarges

In the course of the measurements, a few counters were found to have spurious signals with very rapid risetimes. After studying the counters in question, the source was determined to be microdischarges from the high voltage feedthrough to the grounded case of the counter. An example of such a signal is shown in Fig. 3 along with neutron-capture and alpha-particle signals of a similar amplitude. This is a known effect in ³He proportional counters, and a thorough treatment of the origin of these signals can be found in Ref. [36]. The microdischarge is seen by the preamp as a current pulse and is treated as a normal signal. When sent through a shaping amplifier, there is no way to distinguish a discharge from a signal generated by an incident particle. By using the preamplifier signal, it is possible to measure the fast risetime of these spurious signals and discriminate against them. Fig. 4e shows an example of a detector with a clear band of fast (< 100 ns) risetime events that span a wide range of pulse height. The rate of such events may vary significantly among proportional counters, even for counters made by the same manufacturer.

5. Non-neutron event rejection using energy and risetime discrimination

Utilizing the energy and risetime of these signals provides a straightforward and fast method to characterize the shape and origin of events in a counter, as well as a rationale to define cuts that minimize backgrounds specific to a given application. Fig. 6 is a risetime versus pulse height plot from one of the ³He proportional counters placed in a low counting facility. Based on the regions and limits shown in Fig. 4, one may apply analysis cuts to



Fig. 6. Risetime versus pulse height data from a ³He proportional counter placed in a low background underground lab at the Kimballton Underground Research Facility. The black lines and labels serve to indicate the predominate type of event in the region. The low neutron flux facilitates the study of the internal alphaparticle contamination of the counters. This counter has a broad alpha spectrum and a low microdischarge rate, as shown in the labeled regions.

isolate different event types. Two methods of performing cuts on the data proved useful in this work and are described in this section. First, one may place restrictive cuts that define an essentially neutron-only region. This is particularly important for the applications where the rate from the neutron signal is expected to be very low. Second, one may place less restrictive cuts that eliminate microdischarges and gamma and beta backgrounds but accept the contribution from alpha-particle contamination.

5.1. Neutron-only cut region

In the risetime versus pulse height space, one can define a region essentially free of alpha and beta/gamma backgrounds, as well as microdischarge noise. This region is labeled "Neutrons" in Fig. 6. Three cuts are required to isolate thermal neutron capture events: two diagonal cuts and one in pulse height only. The first diagonal cut is defined by events above the upper edge of the alpha region. The second diagonal cut is defined as events below the upper edge of the distribution from neutron source data. Finally, an upper energy threshold is made to reject events with energy above 0.8 MeV, which is the approximate end of the thermal neutron capture peak.

To test leakage of alpha-particle events into the neutron-only region, we operated several of the proportional counters in a lowneutron environment at the Kimballton Underground Research Facility (KURF) in Ripplemead, VA. The facility is located in an active limestone mine at a depth of 1450 meters water equivalent (mwe) and provides a good low-radioactivity counting environment [37]. Six helium detectors were counted for approximately 16 h. These detectors showed a variety of alpha and microdischarge rates; the risetime versus pulse height spectra from one of the counters is shown in Fig. 6.

It is possible to estimate the leakage of alpha events into the signal region by looking at events above the neutron capture peak in Fig. 6. There should be no neutrons in this region, but the alpha events should have the same probability to spill into the cutregion. For the six detectors tested at KURF, we see an average of 1% of alpha events which are not cut by this method. When these cuts are applied to neutron source data, we find that approximately 55% of the neutrons are rejected. The location of the cuts can be refined to match distributions present in specific counters, as well as the desired signal-to-background ratio.

5.2. Elimination of microdischarge noise only

After surveying the 75 proportional counters using the risetime method, we selected counters that contained the lowest alpha-particle backgrounds to minimize the contamination of the neutron signal region. This permits one to focus on the microdischarge backgrounds. Because the microdischarge region is well-separated from the neutron region, we can use this cut to eliminate the microdischarge events while preserving the entire neutron region.

Without utilizing the information in the pulse risetime, there is nothing to distinguish microdischarge events from alpha-particle events. Thus, more proportional counters would have been categorized as high-background that are low in alpha contamination but high in microdischarges. Counters that would otherwise have been unsuitable for low background experiments can still be used, as long as the microdischarge events are rejected.

6. Direct measurement of alpha-particle rates from ³He counter walls

To verify that the background attributed to alpha particles was being correctly identified, we performed a separate measurement of the alpha activity at the surface of a ³He proportional counter using a silicon detector. These measured rates were compared with existing measurements [38,39] and the internal alpha rates from the counters themselves. Differences between the internal alpha rate and the external rate are expected due to some parameters among the counters cannot be controlled, such as the manufacturing process and knowledge of an absorbing layer deposited on the inner surface.

To measure the alpha-particle rate from the aluminum surface, a ³He proportional counter was placed inside a vacuum chamber, and a silicon surface barrier detector was placed on the surface. The chamber was lined with thin plastic and evacuated to minimize events originating from the chamber and radon in the atmosphere. The silicon detector signal was processed through a preamp and shaping amplifier, and the peak heights were registered on a multichannel analyzer. Data were collected under a number of different conditions in order to identify sources of background. It is important to differentiate between alphas that originate from the metallic surface of the ³He counter and those that may come from the chamber, detector mounting hardware, or the silicon detector itself. Contamination on the surface the ³He counter body is a particular concern, and care was taken to clean the surface with a mild abrasive and detergent before measurements were performed.

Fig. 7 shows the energy spectrum of particles from a ³He counter along with a background spectrum when the counter was removed. To compare with values in the literature, we summed the counts between 2 MeV and 10 MeV after background subtraction when the counter was removed. Taking into account the detector solid angle and scaling up to the full size of the proportional counter, a typical alpha rate was 3.9×10^{-2} s⁻¹. Using the internal surface area of 370 cm² for a counter, one obtains a count rate per area of 1.1×10^{-4} cm⁻² s⁻¹. In Ref. [39], researchers tabulated alpha count rates per area ranging from 1.4×10^{-5} to 1.2×10^{-4} cm⁻² s⁻¹ from commercial aluminum. Rates will vary widely due to several factors, such as the supplier and surface contamination [40], but the values obtained with this proportional counter are consistent with both the literature and our own internal measurements. This gives additional confidence that alpha particles in the proportional counters is being correctly identified.



Fig. 7. The two energy spectra recorded with a silicon surface barrier detector are shown. First is the spectrum from the outside surface of a ³He proportional counter that had been scrubbed clean to remove surface contamination. Also shown is a background spectrum obtained when the ³He counter was removed. The broad, peaked, spectrum of alpha-particle energies is consistent with bulk contamination of the counter body material.

7. Summary

We have detailed a straightforward method to identify neutron, alpha, gamma and beta, and microdischarge events from ³He proportional counter preamplifier signals collected with waveform digitizers. By using the risetime and energy information of each signal, one can isolate neutron capture events and eliminate contamination of the neutron signal from the gamma and beta backgrounds and microdischarges. We have verified that the dominant background in the neutron risetime and energy window arises from alpha decays from the counter body itself. We have identified a set of selection criteria to improve significantly the neutron signal-to-background ratio in that region. If counting statistics is sufficient for a given experiment, or if one desires the highest signal-to-background ratio, one may choose to sacrifice approximately half of the neutron events in order to eliminate nearly all of the alpha-particle events. In many low event rate applications, that may not be the preferred option; in which case, this method will provide an important tool to identify and mitigate the contribution of alpha backgrounds in ³He proportional counters.

A fast neutron spectrometer utilizing six ³He proportional counters [21] was placed underground at KURF. Using the risetime method to identify and reduce backgrounds has led to a significant improvement of the experiment. In future work, we will investigate the improvement in the neutron signal-to-background from data acquired in a low-background environment.

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References

- W.R. Mills, R.I. Caldwell, I.L. Morgan, Review of Scientific Instruments 33 (1962) 866.
- [2] A. Sayres, M. Coppola, Review of Scientific Instruments 35 (1964) 431.
- [3] A. Izumi, Y. Murata, Nuclear Instruments and Methods in Physics Research Section A 94 (1971) 141.
- [4] J. Delorme, M. Ericson, T. Ericson, P. Vogel, Physical Review D 52 (1995) 2222.
 [5] J. Formaggio, C. Martoff, Annual Review of Nuclear and Particle Science 54 (2004) 361.
- [6] D.-M. Mei, A. Hime, Physical Review D 73 (5) (2006) 053004, http://dx.doi.org/ 10.1103/PhysRevD.73.053004.
- [7] R.G.H. Robertson, Progress in Particle and Nuclear Physics 40 (1998) 113.
- [8] V. Chazal, R. Brissot, J. Cavaignac, B. Chambon, M. de Jesus, D. Drain, Astroparticle Physics 9 (1998) 163.
- [9] S.V. Rozov, V.B. Brudanin, A.V. Lubashevskiy, S.S. Semikh, D.V. Filosofov, E. A. Yakushev, Bulletin of the Russian Academy of Science: Physics 74 (2010) 464.
- [10] G. Cox, C. Duba, M. Howe, S. McGee, A. Myers, K. Rielage, R. Robertson, L. Stonehill, B. Wall, J. Wilkerson, T. Van Wechel, IEEE Transactions on Nuclear Science NS-51 (2004) 2227.
- [11] J.F. Amsbaugh, J.M. Ánaya, J. Banar, T.J. Bowles, M.C. Browne, T.V. Bullard, T. H. Burritt, G.A. Cox-Mobrand, X. Dai, H. Deng, M. di Marco, P.J. Doe, M. R. Dragowsky, C.A. Duba, F.A. Duncan, E.D. Earle, S.R. Elliott, E. Esch, H. Fergani, J.A. Formaggio, M.M. Fowler, J.E. Franklin, P. Geissbühler, J.V. Germani, A. Goldschmidt, E. Guillian, A.L. Hallin, G. Harper, P.J. Harvey, R. Hazama, K. M. Heeger, J. Heise, A. Hime, M.A. Howe, M. Huang, L.L. Kormos, C. Kraus, C. B. Krauss, J. Law, I.T. Lawson, K.T. Lesko, J.C. Loach, S. Majerus, J. Manor, S. McGee, K.K.S. Miknaitis, G.G. Miller, B. Morissette, A. Myers, N.S. Oblath, H. M. O'Keeffe, R.W. Ollerhead, S.J.M. Peeters, A.W.P. Poon, G. Prior, S.D. Reitzner, K. Rielage, R.G.H. Robertson, P. Skensved, A.R. Smith, M.W.E. Smith, T. D. Steiger, L.C. Stonehill, P.M. Thornewell, N. Tolich, B.A. Vandevender, T. D. van Wechel, B.L. Wall, H. Wan Chan Tseung, J. Wendland, N. West, J. B. Wilhelmy, J.F. Wilkerson, J.M. Vouters, Nuclear Instruments and Methods in Physics Research Section A 579 (2007) 1054.
- [12] B. Aharmim, et al., Combined Analysis of all Three Phases of Solar Neutrino Data from the Sudbury Neutrino Observatory, 2011. arXiv:1109.0763.
- [13] M. Browne, Preparation for Deployment of the Neutral Current Detectors (NCDs) for the Sudbury Neutrino Observatory, Ph.D. Thesis, University of Washington, 1999.
- [14] L.C. Stonehill, Deployment and Background Characterization of the Sudbury Neutrino Observatory Neutral Current Detectors, Ph.D. Thesis, University of Washington, 2005.
- [15] R. Martin, An Analysis of the ³He Proportional Counter Data from the Sudbury Neutrino Observatory Using Pulse Shape Discrimination, Ph.D. Thesis, Queen's University, 2009.
- [16] R. Batchelor, R. Aves, T.H.R. Skryme, Review of Scientific Instruments 26 (1955) 1037.
- [17] D.L. Chichester, J.T. Johnson, E.H. Seabury, Applied Radiation and Isotopes 70 (2012) 1457.
- [18] D.J. Thomas, A.V. Alevra, Nuclear Instruments and Methods in Physics Research Section A 476 (2002) 12.
- [19] B. Wiegel, A. Alevra, Nuclear Instruments and Methods in Physics Research Section A 476 (2002) 36.
- [20] J. Abdurashitov, V. Gavrin, A. Kalikhov, V. Matushko, A. Shikhin, V. Yants, O. Zaborskaia, Nuclear Instruments and Methods in Physics Research Section A 476 (2002) 322.
- [21] T. Langford, E. Beise, H. Breuer, D. Erwin, C. Bass, C. Heimbach, J. Nico, Bulletin of the American Physical Society 55 (2010) 1.
- [22] C. Aalseth, M. Amman, F.A. III, H. Back, A. Barabash, P. Barbeau, M. Bergevin, F. Bertrand, M. Boswell, V. Brudanin, W. Bugg, T. Burritt, M. Busch, G. Capps, Y.-D. Chan, J. Collar, R. Cooper, R. Creswick, J. Detwiler, J. Diaz, P. Doe,

Y. Efremenko, V. Egorov, H. Ejiri, S. Elliott, J. Ely, J. Esterline, H. Farach, J. Fast, N. Fields, P. Finnerty, B. Fujikawa, E. Fuller, V. Gehman, G. Giovanetti, V. Guiseppe, K. Gusey, A. Hallin, G. Harper, R. Hazama, R. Henning, A. Hime, E. Hoppe, T. Hossbach, M. Howe, R. Johnson, K. Keeter, M. Keillor, C. Keller, J. Kephart, M. Kidd, A. Knecht, O. Kochetov, S. Konovalov, R. Kouzes, L. Leviner, I. Loach, P. Luke, S. MacMullin, M. Marino, R. Martin, D.-M. Mei, H. Milev, M. Miller, L. Mizouni, A. Meyers, M. Nomachi, J. Orrell, D. Peterson, D.P. II, A. Poon, G. Prior, J. Qian, D. Radford, K. Rielage, R. Robertson, L. Rodriguez, K. Rykaczewski, H. Salazar, A. Schubert, T. Shima, M. Shirchenko, D. Steele, J. Strain, G. Swift, K. Thomas, V. Timkin, W. Tornow, T.V. Wechel, I. Vanyushin, R. Varner, K. Vetter, J. Wilkerson, B. Wolfe, W. Xiang, E. Yakushev, H. Yaver, A. Young, C.-H. Yu, V. Yumatov, C. Zhang, S. Zimmerman, Nuclear Instruments and Methods in Physics Research Section A 652 (1) (2011) 692 http://dx.doi org/10.1016/j.nima.2010.08.100. (Symposium on Radiation Measurements and Applications (SORMA) XII 2010. (http://www.sciencedirect.com/science/arti cle/pii/S0168900210018875>).

- [23] R.S. Raghavan, the LENS Collaboration, Journal of Physics: Conference Series 120 (5) (2008) 052014. (http://stacks.iop.org/1742-6596/120/i=5/a=052014).
- [24] S. MacMullin, G. Giovanetti, M. Green, R. Henning, R. Holmes, K. Vorren, J. Wilkerson, Journal of Environmental Radioactivity 112 (2012) 165, http://dx. doi.org/10.1016/j.jenvrad.2012.01.026. science/article/pii/S0265931X12000446>.
- [25] S.D. Rountree, R.B. Vogelaar, Kimballton Underground Research Facility, in: APS Meeting Abstracts, 2012, p. 11001.
- [26] N. Takeda, K. Kudo, IEEE Transactions on Nuclear Science NS-42 (1995) 548.
- [27] J.F. Ziegler, M. Ziegler, J. Biersack, Nuclear Instruments and Methods in Physics Research Section A 268 (11–12) (2010) 1818, http://dx.doi.org/10.1016/j. nimb.2010.02.091. (19th International Conference on Ion Beam Analysis. (http://www.sciencedirect.com/science/article/pii/S0168583X10001862)).
- [28] H. Andersen, J. Ziegler, J. Biersack, U. Littmark, The Stopping and Ranges of Ions in Matter, Pergamon Press, 1977.
- [29] W. Leo, Techniques for Nuclear and Particle Physics Experiments: A How-To Approach, 2nd ed., Springer-Verlag, 1994.
- [30] G.F. Knoll, Radiation Detection and Measurement, 2nd Ed., Wiley, 1989.
- [31] B. Beltran, H. Bichsel, B. Cai, G.A. Cox, H. Deng, J. Detwiler, J.A. Formaggio, S. Habib, A.L. Hallin, A. Hime, M. Huang, C. Kraus, H.R. Leslie, J.C. Loach, R. Martin, S. McGee, M.L. Miller, B. Monreal, J. Monroe, N.S. Oblath, S.J. M. Peeters, A.W.P. Poon, G. Prior, K. Rielage, R.G.H. Robertson, M.W.E. Smith, L.C. Stonehill, N. Tolich, T.V. Wechel, H.W.C. Tseung, J. Wendland, J. F. Wilkerson, A. Wright, New Journal of Physics 13 (2011) 073006.
- [32] I. Kitani, K. Yamauchi, K. Arii, The charge spreading due to microdischarges on insulating films from an evaporated metal electrode subjected to a DC ramp voltage, in: IEEE Proceedings of the 4th International Conference on Properties and Applications of Dielectric Materials, vol. 1, 1994, p. 282
- [33] M.G. Danikas, IEEE Transactions on Dielectrics and Electrical Insulation 4 (1997) 863.
- [34] M.G. Danikas, G. Adamidis, Electrical Engineering 80 (1997) 105.
- [35] N. Johnson, private communication, 2011.
- [36] K. Heeger, S. Elliott, R. Robertson, M. Smith, T. Steiger, J. Wilkerson, IEEE Transactions on Nuclear Science NS-47 (2000) 1829.
- [37] P. Finnerty, S. MacMullin, H.O. Back, R. Henning, A. Long, K.T. Macon, J. Strain, R.M. Lindstrom, R.B. Vogelaar, Nuclear Instruments and Methods in Physics Research Section A 642 (2010) 65.
- [38] S.R. Hashemi-Nezhad, L.S. Peak, Nuclear Instruments and Methods in Physics Research Section A 416 (1998) 100.
- [39] B. Al-Bataina, J. Janecke, Nuclear Instruments and Methods in Physics Research Section A 255 (1987) 512.
- [40] Z. Dębicki, K. Jędrzejczak, J. Karczmarczyk, M. Kasztelan, R. Lewandowski, J. Orzechowski, B. Szabelska, J. Szabelski, P. Tokarski, T. Wibig, Astrophysics and Space Sciences Transactions 7 (2011) 511.